# GENERIC REGULARITY OF FREE BOUNDARIES FOR THE OBSTACLE PROBLEM 

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#### Abstract

The goal of this paper is to establish generic regularity of free boundaries for the obstacle problem in $\mathbf{R}^{n}$. By classical results of Caffarelli, the free boundary is $\mathrm{C}^{\infty}$ outside a set of singular points. Explicit examples show that the singular set could be in general $(n-1)$-dimensional - that is, as large as the regular set. Our main result establishes that, generically, the singular set has zero $\mathcal{H}^{n-4}$ measure (in particular, it has codimension 3 inside the free boundary). Thus, for $n \leq 4$, the free boundary is generically a $\mathrm{C}^{\infty}$ manifold. This solves a conjecture of Schaeffer (dating back to 1974) on the generic regularity of free boundaries in dimensions $n \leq 4$.


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## 1. Introduction

Several fundamental problems in science (physics, biology, finance, geometry, etc.) can be described by PDEs that exhibit a-priori unknown interfaces or boundaries. They are called free boundary problems, and have been a major line of research in the PDE community in the last 60 years; see for instance [LS67, LS69, Kin73, BK74, KN77, Caf77, CR77, Sak91, Caf98, W99, CKS00, Mon03, SU03, ACS08, GP09, ALS13, FS19].

The obstacle problem
(1.1)

$$
\Delta u=\chi_{\{u>0\}} \quad \text { in } \quad \Omega \subset \mathbf{R}^{n}
$$

$$
u \geq 0
$$

[^0]is the most classical and among the most important elliptic free boundary problems, and it arises in a variety of situations; see e.g. [DL76, Fri82, Rod87, PSU12, Ser 15].

From the mathematical point of view, the most challenging question in this context is to understand the regularity of free boundaries. The modern development of the regularity theory for free boundaries started in the late 1970's with the seminal paper of Caffarelli [Caf77], and since then it has been a very active area of research.

The main result in [Caf77] establishes that, for any solution of (1.1), the free boundary $\partial\{u>0\}$ is $\mathrm{C}^{\infty}$ outside a closed set of singular points. Singular points arise for example when the free boundary creates cusps, and they may appear in any dimension $n \geq 2$. By [CR76, Caf98, Mon03], these points are locally contained in a $\mathrm{C}^{1}$ manifold of dimension $n-1$. More recently, finer estimates at singular points were established in [CSV18, FS19].
1.1. Generic regularity for the obstacle problem. - A major question in the understanding of singularities in PDE theory is the development of methods to prove generic regularity results. In the context of the obstacle problem (1.1), the key question is to understand the generic regularity of free boundaries. Explicit examples [Sch76] show that singular points in the obstacle problem can form a set of dimension $n-1$ (thus, as large as the whole free boundary). Still, singular points are expected to be rare [Sch74]:

Conjecture (Schaeffer, 1974). - Generically, free boundaries in the obstacle problem have no singular points.

The conjecture is only known to hold in the plane $\mathbf{R}^{2}$ [Mon03], and up to now nothing was known in the physical space $\mathbf{R}^{3}$ or in higher dimensions.

Notice that, in the obstacle problem, the question of generic regularity is particularly relevant, since in such context the singular set can be as large as the regular setwhile in other problems the singular set has lower Hausdorff dimension [Giu84]. Also, from the point of view of applications (see [Bai74, DL76, Rod87, Ser15]), it is particularly important to understand the problem in the physical space $\mathbf{R}^{3}$.

A main goal of this paper is to prove Schaeffer's conjecture in $\mathbf{R}^{3}$ and $\mathbf{R}^{4}$. To this aim, we consider any monotone family of solutions $\left\{u^{t}\right\}_{t \in(-1,1)}$ of (1.1) in $\mathrm{B}_{1}$ satisfying the following "uniform monotonicity" condition:

For every $t \in(-1,1)$ and any compact set $\mathrm{K}_{t} \subset \partial \mathrm{~B}_{1} \cap\left\{u^{t}>0\right\}$ there exists $c_{\mathrm{K}_{t}}>0$ such that

$$
\begin{equation*}
\inf _{x \in \mathrm{~K}_{t}}\left(u^{t^{\prime}}(x)-u^{t}(x)\right) \geq c_{\mathrm{K}_{t}}\left(t^{\prime}-t\right), \quad \text { for all }-1<t<t^{\prime}<1 . \tag{1.2}
\end{equation*}
$$

This condition rules out the existence of regions that remain stationary as we increase the parameter $t$. In case that $u^{t}$ is continuously differentiable with respect to $t$, then such condition is equivalent to saying that $\partial_{t} u^{t}>0$ inside $\left\{u^{t}>0\right\}$.

We shall also assume that $\left.(-1,1) \ni t \mapsto u^{t}\right|_{\partial \mathrm{B}_{1}} \in \mathrm{~L}^{\infty}\left(\partial \mathrm{B}_{1}\right)$ is continuous with respect to $t$. Note that, by the maximum principle, this implies that $(-1,1) \ni t \mapsto u^{t} \in$ $L^{\infty}\left(B_{1}\right)$ is continuous. Under this assumption, we prove the following:

Theorem 1.1.- Let $\left\{u^{t}\right\}_{t \in(-1,1)}$ be a monotone and continuous family of solutions to (1.1) in $\mathrm{B}_{1} \subset \mathbf{R}^{n}$ satisfying (1.2), and let $\Sigma^{t} \subset \partial\left\{u^{t}>0\right\} \cap \mathrm{B}_{1}$ be the set of singular points for $u^{t}$. Then

$$
\mathcal{H}^{n-4}\left(\Sigma^{t}\right)=0 \quad \text { for a.e. } t \in(-1,1)
$$

In particular, Schaeffer's conjecture holds for $n \leq 4$.
We remark that very few results are known in this direction for elliptic PDE, and most of them deal only with simpler situations (for instance the obstacle problem in $\mathbf{R}^{2}$ [Mon03]), or when the singular set is known to be very small (as in the case of areaminimizing hypersurfaces in $\mathbf{R}^{8}$ [Sma93]).

As a particular family of solutions to which our Theorem 1.1 applies, one can consider the solution $u^{t}$ to the obstacle with boundary data $\left.u^{t}\right|_{\partial \mathrm{B}_{1}}=g+t$ (similarly to what was done in [Mon03]), but many other choices are possible.

In particular, due to the general character of our assumption (1.2), we can apply Theorem 1.1 (more precisely, some of the results behind its proof) to study the HeleShaw flow. This is a well-known 2D model which describes a flow between two parallel flat plates following Darcy law [HS1898, CJK07]. After a transformation of the type $u(x, t)=\int_{0}^{t} p(x, \tau) d \tau$-where $p(x, t)$ is the pressure-the problem becomes
(1.3) $\quad u=t \quad$ in $\mathrm{K} \times(0, \mathrm{~T}) \subset \mathbf{R}^{2} \times \mathbf{R}$

$$
u \geq 0,
$$

where $\mathrm{K} \subset \mathbf{R}^{2}$ is a given compact set, and $\mathrm{K}^{c}:=\mathbf{R}^{2} \backslash \mathrm{~K}$. Since the singular set is closed inside the free boundary (see for instance Lemma 6.2(a)), as a consequence of our fine analysis of singular points, we can also show the following:

Theorem 1.2. - Let $\mathrm{K} \subset \mathbf{R}^{2}$ be any compact set, and $u(x, t)$ be any solution to the Hele-Shaw flow (1.3). Let $\Sigma^{t} \subset \mathrm{~K}^{c}$ be the set of singular points of $\partial\{u(\cdot, t)>0\}$, and let $\mathcal{S}:=\{t \in(0, \mathrm{~T})$ : $\left.\Sigma^{t} \neq \emptyset\right\}$ be the set of singular times. Then $\mathcal{S}$ is relatively closed inside $(0, \mathrm{~T})$ and

$$
\operatorname{dim}_{\mathcal{H}}(\mathcal{S}) \leq \frac{1}{4}
$$

In particular, the free boundary is $\mathrm{C}^{\infty}$ for a.e. time $t \in(0, \mathrm{~T})$.
Prior to our result, it was an open question to decide whether singularities in such model could persist in time or not. Theorem 1.2 answers this question, and provides for the first time an estimate on the set of singular times.
1.2. Higher-order expansions at most singular points. - A key tool in the proof of Theorem 1.1 is a very fine understanding of singular points, as explained next.

For the obstacle problem (1.1), a classical result of Caffarelli [Caf98] states that at every singular point $x_{0}$ we have an expansion of the form

$$
\begin{equation*}
u(x)=p_{2, x_{0}}\left(x-x_{0}\right)+o\left(\left|x-x_{0}\right|^{2}\right), \tag{1.4}
\end{equation*}
$$

where $p_{2, x_{0}}$ is a nonnegative, homogeneous, quadratic polynomial satisfying $\Delta p_{2, x_{0}} \equiv 1$.
In dimension $n=2$ this estimate was improved in [W99] by replacing $o\left(\left|x-x_{0}\right|^{2}\right)$ with $\mathrm{O}\left(\left|x-x_{0}\right|^{2+\alpha}\right)$ for some $\alpha>0$, and in arbitrary dimensions it was shown in [CSV18] that $o\left(\left|x-x_{\circ}\right|^{2}\right)$ can be replaced by $\mathrm{O}\left(\left|x-x_{\circ}\right|^{2}|\log | x-\left.x_{0}\right|^{-\epsilon}\right)$, for some $\epsilon>0$. More recently, it was proved by the first and third authors [FS19] that, in every dimension $n$, one actually has

$$
u(x)=p_{2, x_{0}}\left(x-x_{0}\right)+\mathrm{O}\left(\left|x-x_{0}\right|^{3}\right)
$$

possibly outside a set of "anomalous" singular points whose Hausdorff dimension is at most $n-3$.

Here, in order to prove our main result, we need to improve substantially the understanding of singular points, establishing a new higher order expansion at most singular points for monotone families of solutions to the obstacle problem. Here and in the sequel, $\operatorname{dim}_{\mathcal{H}}$ will denote the Hausdorff dimension (see Section 7 for a definition).

Theorem 1.3. - Let $\left\{u^{t}\right\}_{t \in(-1,1)}$ be a family of solutions to (1.1) in $\mathbf{B}_{1} \subset \mathbf{R}^{n}$ which is continuous and nondecreasing in $t$ (in particular, they could be independent of $t$ ). Let $\Sigma^{t} \subset \partial\left\{u^{t}>0\right\} \cap \mathbf{B}_{1}$ be the set of singular points of $u^{t}$, and $\hat{\Sigma}:=\bigcup_{t \in(-1,1)} \Sigma^{t} \subset \mathrm{~B}_{1}$.

Then there exists a set $\mathrm{E} \subset \hat{\Sigma}$, with $\operatorname{dim}_{\mathcal{H}}(\mathrm{E}) \leq n-2$, such that for every $t_{0} \in(-1,1)$ and every $x_{0} \in \Sigma^{t_{0}} \backslash \mathrm{E}$ we have

$$
\begin{equation*}
u^{t_{0}}(x)=\mathrm{P}_{4, x_{0}, t_{0}}\left(x-x_{0}\right)+\mathrm{O}\left(\left|x-x_{0}\right|^{5-\zeta}\right) \tag{1.5}
\end{equation*}
$$

for all $\zeta>0$, where $\mathrm{P}_{4, x_{0}, t_{0}}$ is a fourth order polynomial satisfying $\Delta \mathrm{P}_{4, x_{0}, t_{0}} \equiv 1$.
An important point here is that the dimension $n-2$ of the "bad" set E is sharp. Indeed, by well known examples in $\mathbf{R}^{2}$ (see e.g. [Sak93]), one can construct solutions $u$ whose singular set contains a $(n-2)$-dimensional subset E for which (1.5) does not hold at any point in E .

As the reader will see from the proof, when $p_{2, x_{0}, t_{0}}(x)=\frac{1}{2}(x \cdot e)^{2}$ for some unit vector $e \in \mathbf{R}^{n}$ then the expansion (1.5) can be written alternatively as

$$
u\left(x_{\circ}+x\right)=\frac{1}{2}(e \cdot x+p(x))^{2}+\mathrm{O}\left(|x|^{5-\zeta}\right)
$$

for a certain polynomial $p$ of degree 3 with no linear or constant terms. Geometrically, this expansion-together with a Lipschitz estimate that we will establish later-yields
that, around most singular points, the contact set is contained inside a set of the form $\left\{|e \cdot x+p(x)| \leq \mathrm{C}|x|^{4-\zeta}\right\}$. Thus, if the free boundary has a cusp, then at most points this cusp must be very thin. It is worth noticing that the expression of $p$ (or equivalently, of $\mathrm{P}_{4, x_{0}, t_{0}}$ ) is related to the curvature of the free boundary near a singular point. In particular, whenever the solution is even with respect to the hyperplane $\{e \cdot x=0\}$, then $p \equiv 0$ (and thus $\left.\mathrm{P}_{4, x_{0}, t_{0}}(x) \equiv \frac{1}{2}(e \cdot x)^{2}\right)$, since there are no curvature terms.

To establish Theorem 1.1 we need to introduce a variety of new ideas, combining Geometric Measure Theory tools, PDE estimates, several dimension reduction arguments, and new monotonicity formulas. This is explained in more detail next.
1.3. On the proofs of the main results. - Let us give an overview of the main ideas introduced in this paper.
1.3.1. From expansion to cleaning: a Sard-type approach. - The starting idea to prove Theorem 1.1 is the following: denote by $\Sigma^{t} \subset \partial\left\{u^{t}>0\right\} \cap \mathrm{B}_{1}$ the set of singular points of $u^{t}$. Assume that, for some fixed $t_{0}$, we have $x_{0} \in \Sigma^{t_{0}}$ and $u^{t_{0}}$ has an expansion of the form

$$
\begin{equation*}
u^{t_{0}}\left(x_{\circ}+x\right)=\mathrm{P}(x)+\mathrm{O}\left(|x|^{\lambda}\right) \tag{1.6}
\end{equation*}
$$

for some $\lambda \geq 2$ and some polynomial P such that $\Delta \mathrm{P} \equiv 1$. Note that, since $u^{t_{0}} \geq 0$, the expansion above implies that $\mathrm{P}(x) \geq-\mathrm{O}\left(|x|^{\lambda}\right)$. Hence, for any $r>0, \mathrm{P}+\mathrm{C} r^{\lambda}$ is a solution to the obstacle problem in $\mathrm{B}_{r}$ with an empty contact set, and $u^{t_{0}}\left(x_{\circ}+\cdot\right)$ is $\mathrm{O}\left(r^{\lambda}\right)$-close to it. This suggests that, thanks to the monotonicity assumption (1.2), by slightly increasing the value of the parameter $t$ the contact set of $u^{t}$ inside $\mathrm{B}_{r}\left(x_{0}\right)$ will disappear.

To make this argument quantitative we need to introduce a series of delicate barrier constructions which actually depend on the fine structure of the singular point $x_{0}$, see Section 9. In this way we are able to prove that, for a increment of $t$ of size $\delta t \sim r^{\lambda-1}$, the contact set of $u^{t+\delta t}$ is $\mathbf{B}_{r}$ disappears: more precisely we can show that, for some $\mathbf{C}>0$,

$$
\begin{equation*}
\Sigma^{t_{0}+\mathrm{C}_{\gamma^{\lambda-1}}} \cap \mathrm{~B}_{r}\left(x_{\mathrm{o}}\right)=\emptyset \quad \text { for } r>0 \text { sufficiently small. } \tag{1.7}
\end{equation*}
$$

As we will explain better below, we can prove an expansion as in (1.6) for $\lambda$ belonging a discrete set $\Lambda$ (see (1.9) for a definition of this set).

Hence, for $t_{\circ} \in(-1,1)$ and $x_{\circ} \in \Sigma^{t_{0}} \subset \partial\left\{u^{t_{0}}>0\right\}$, we define $\lambda_{x_{0}, t_{0}}$ to be the maximal $\lambda \in \Lambda$ for which we can prove an expansion as in (1.6). Then, for each $\lambda \in \Lambda$ and $t \in(-1,1)$, we define $\Sigma^{t, \lambda}$ as the set of $x_{0} \in \Sigma_{t}$ for which $\lambda_{x_{0}, t}=\lambda$. In other words, $\Sigma^{t, \lambda}$ is defined as the set of points at which we have a polynomial expansion up to order $\lambda \in \Lambda$ but not better. Then, a covering argument "à la Sard" yields
(1.8)

$$
\operatorname{dim}_{\mathcal{H}}\left(\left\{t \in(-1,1): \Sigma^{t, \lambda} \neq \emptyset\right\}\right) \leq \frac{\operatorname{dim}_{\mathcal{H}}\left(\bigcup_{t \in(-1,1)} \Sigma^{t, \lambda}\right)}{\lambda-1}
$$

(see Proposition 7.7(a) for a more refined statement). In particular, if the right hand side is strictly less than 1 , then $\Sigma^{t, \lambda}=\emptyset$ for a.e. $t$. On the other hand, if the right hand side is larger or equal to 1 , then a coarea-type argument allows us to show that $\Sigma^{t, \lambda}$ is very small for a.e. $t$ (see Proposition 7.7(b)).

In view of the given description of our approach, our goals are the following:
(1) given a singular point, prove an expansion up to order $\mathrm{O}\left(|x|^{\lambda}\right)$ with $\lambda \in \Lambda$ as large as possible;
(2) given $\lambda \in \Lambda$, estimate the dimension of $\bigcup_{t \in(-1,1)} \Sigma^{t, \lambda}$, i.e., the set of points where the expansion stops at $\lambda$.
1.3.2. A higher-order expansion at singular points: the case of a single solution. - To understand these questions in a simplified situation, one can first look at the problem without the parameter $t$. So, given a solution $u$ to the obstacle problem and a singular point $x_{0}$, we want to obtain a Taylor expansion around $x_{\circ}$ at the highest possible order. This will require several steps, described below.
(a) Second blow-up: a cubic expansion at most points. - This first part is essentially contained in [FS19]. Recall first that, as proven in [Caf98], for any singular point $x_{\circ}$ we have (1.4), that is, $p_{2, x_{o}}$ is tangent up to second order to $u\left(x_{\circ}+\cdot\right)$ at 0 . Equivalently

$$
\frac{u\left(x_{\circ}+r \cdot\right)}{r^{2}} \rightarrow p_{2, x_{0}} \quad \text { as } r \rightarrow 0
$$

and $p_{2, x_{0}}$ is called the "first blow-up" of $u$ at $x_{0}$.
One can then catalog singular points according to the dimension of the kernel of $p_{2, x_{0}}$ : given $m \in\{0, \ldots, n-1\}$, we say

$$
x_{0} \in \Sigma_{m} \Longleftrightarrow \operatorname{dim}\left(\left\{p_{2, x_{0}}=0\right\}\right)=m .
$$

We then consider the "second blow-ups", namely, the possible limits of the functions

$$
\tilde{w}_{2, r}(x):=\frac{w_{2}(r x)}{\left\|w_{2}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}}, \quad \text { where } w_{2}:=u\left(x_{\circ}+\cdot\right)-p_{2, x_{0}}
$$

as $r \rightarrow 0$. As shown in [FS19], $r \mapsto\left\|\nabla \tilde{w}_{2, r}\right\|_{L^{2}\left(\mathrm{~B}_{1}\right)}$ is monotonically increasing (equivalently, the so-called Almgren frequency formula is monotone on $w_{2}$ ). Thanks to this fact, setting $\lambda_{2}:=\lim _{r \rightarrow 0}\left\|\nabla \tilde{w}_{2, r}\right\|_{L^{2}\left(\mathrm{~B}_{1}\right)}$, one can characterize all the blow-ups (namely, the accumulation points of $\left\{\tilde{w}_{2, r}\right\}$ as $r \rightarrow 0$ ): they are $\lambda_{2}$-homogeneous functions $q$

- either satisfying

$$
\Delta q=0 \quad \text { in } \mathbf{R}^{n}, \quad \text { if } x_{\circ} \in \Sigma_{m} \text { with } m \leq n-2,
$$

- or solving the Signorini problem

$$
\begin{aligned}
& \Delta q \leq 0, \quad q \Delta q \equiv 0,\left.\quad \Delta q\right|_{\mathbf{R}^{n} \backslash \mathrm{~L}}=0,\left.\quad q\right|_{\mathrm{L}} \geq 0 \\
& \left.\quad \text { if } \mathrm{L}:=\left\{p_{2, x_{0}}=0\right\} \text { is a hyperplane (i.e., } x_{0} \in \Sigma_{n-1}\right) .
\end{aligned}
$$

In the first case (i.e., $m \leq n-2$ ), since $q$ is harmonic in $\mathbf{R}^{n}$ it must be $\lambda_{2} \in$ $\{2,3,4, \ldots\}$. Also, following [FS19], one can show that $\lambda_{2} \geq 3$ up to an "anomalous set" of dimension $m-1$ inside $\Sigma_{m}$. This implies that, for $x_{\circ}$ outside this anomalous set, we have $u=p_{2, x_{0}}\left(x-x_{0}\right)+\mathrm{O}\left(\left|x-x_{0}\right|^{3}\right)$. Note that applying this result to $u=u^{t_{0}}$ gives an expansion as in (1.6) for $\lambda=3$. As we will explain in Section 1.3.4(a) below, when $m \leq n-2$ we are able to improve (1.7) so that this expansion suffices for proving our main theorem.

The real challenge is to understand the set $\Sigma_{n-1}$. In this case, since $q$ solves the Signorini problem, thanks to a classification result for 2D solutions and a dimension reduction argument, we can show that

$$
\begin{equation*}
\lambda_{2} \in \Lambda:=\{2,3,4, \ldots\} \cup\left\{\frac{7}{2}, \frac{11}{2}, \frac{15}{2}, \ldots\right\} \tag{1.9}
\end{equation*}
$$

outside a set of singular points of dimension $n-3$. Also it is easy to prove that, in this case, $\lambda_{2} \neq 2$; thus, except for a small set, the lowest possible value for $\lambda_{2}$ is 3 . The main challenge is now to improve this cubic expansion to higher order.
(b) Third blow-up: a delicate dichotomy. - From now on, we focus on points of $\Sigma_{n-1}$ where $\lambda_{2}=3$, i.e., $q$ is a 3rd-order homogeneous solution of Signorini (as one can see from the coming argument, the other cases can be considered as a particular case of this taking $\lambda_{2}=3$ and $p_{3, x_{0}} \equiv 0$ in the definition of $w_{3}$ below). Two possibilities arise, depending whether some accumulation point $q$ of $\left\{\tilde{w}_{2, r}\right\}$ as $r \rightarrow 0$ is harmonic or not. These two cases have to be analyzed separately.

- The third blow-up is not harmonic: a new uniqueness result. By another dimension reduction argument, we can prove that the set where $q$ is not harmonic has dimension $n-2$. However this is not enough, and here comes one of the key arguments introduced in this paper: as explained in Section 1.3.4(b), in order to obtain Schaeffer's conjecture in $\mathbf{R}^{4}$ we need to prove that the limit $q$ of $\tilde{w}_{2, r}$ is unique, and that this set is $(n-2)$-rectifiable. To accomplish, in Section 5 we introduce new differential formulae, compactness and barrier arguments, and a delicate ODE-type lemma, that allow us to obtain the uniqueness of blow-ups (taking quotients of suitable qualities) even if we lack a monotonicity formula.
- The third blow-up is harmonic: a monotonicity argument at nondegenerate points. Assume that there exists a harmonic accumulation point $q$. Then (thanks to a Monneau-type monotonicity formula) we can show that the limit $\lim _{r \downarrow 0} \tilde{w}_{2, r}$ exists (i.e., all accumulation points coincide) and that $u\left(x_{0}+\cdot\right)=p_{2, x_{0}}+p_{3, x_{0}}+o\left(|x|^{3}\right)$ for some 3-homogeneous harmonic polynomial $p_{3, x_{0}}\left(p_{3, x_{0}}\right.$ being a multiple of $\left.q\right)$. This suggests to iterate the previous
blow-up procedure by defining

$$
\tilde{w}_{3, r}(x):=\frac{w_{3}(r x)}{\left\|w_{3}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}}, \quad \text { where } w_{3}:=u-p_{2, x_{0}}-p_{3, x_{0}},
$$

and try to mimicking the argument described above. Unfortunately, in this case it is not true anymore that $r \mapsto\left\|\nabla \tilde{w}_{3, r}\right\|_{L^{2}\left(\mathbf{B}_{1}\right)}$ is monotonically increasing. Still, by a delicate bootstrapping argument (cf. Lemma 4.3) we can prove that ${ }^{1}$

$$
\begin{gathered}
r \mapsto\left\|\nabla \tilde{w}_{3, r}\right\|_{L^{2}\left(\mathrm{~B}_{1}\right)} \text { is almost increasing, provided } \\
\left\|w_{3}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \gtrsim r^{4-\varepsilon} \text { for some } \varepsilon>0 .
\end{gathered}
$$

Therefore, under this nondegeneracy assumption, we can consider accumulation points of $\tilde{w}_{3, r}$ and prove that they are $\lambda_{3}$-homogeneous solution of Signorini for some $\lambda_{3} \in$ $[3,4)$. Then, by a dimension reduction argument (based again on the classification of 2D solutions), we can prove that $\lambda_{3} \in\{3,7 / 2\}=\Lambda \cap[3,4)$ in a set of dimension $n-2$, and the remaining points are of codimension 3 . So, to summarize:
(i) for most points in $\Sigma_{n-1}$ where the limit of $\tilde{w}_{2, r}$ is harmonic, the assumption $\left\|w_{3}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \gtrsim r^{4-\varepsilon}$ fails for every $\varepsilon>0$, except perhaps in a set of dimension $n-2$;
(ii) if $\left\|w_{3}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \gtrsim r^{4-\varepsilon}$ holds, then the blow-up is $\lambda_{3}$-homogeneous for $\lambda_{3} \in$ $\{3,7 / 2\}$, except for a set of dimension at most $n-3$.

In order to prove Schaeffer's conjecture in $\mathbf{R}^{4}$, it is important for us to rule out the possibility that $\lambda_{3}=3$ in case (ii). This is highly nontrivial, and follows from the analysis performed in Section 5 to understand points where $\tilde{w}_{2, r}$ converges to a non-harmonic function.
(c) Fourth blow-up: monotonicity via a new ansatz and proof of (1.5). - To go further in our analysis and prove our main theorem, we now need to investigate the set of points where case (i) happens, namely $\left\|w_{3}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \gtrsim r^{4-\varepsilon}$ fails for every $\varepsilon>0$. In this case Almgren's monotonicity formula fails on $w_{3, r}$, and therefore a new approach needs to be found.

The key idea here is to replace $w_{3}=u\left(x_{\circ}+\cdot\right)-p_{2, x_{0}}-p_{3, x_{0}}$ with a much more refined Ansatz $\mathrm{W}_{3}:=u\left(x_{\circ}+\cdot\right)-\mathscr{P}_{x_{0}}$ which takes into account both the curvature of the free boundary and the non-negativity of the solution - this is done in Definition 4.5. Doing so, and defining $\tilde{W}_{3, r}$ in analogy to what done before, we can show (again after a bootstrap argument) that

$$
\begin{aligned}
& r \mapsto\left\|\nabla \tilde{\mathrm{~W}}_{3, r}\right\|_{\mathrm{L}^{2}\left(\mathbf{B}_{1}\right)} \text { is almost increasing, provided } \\
& \quad\left\|\mathrm{W}_{3}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathbf{B}_{1}\right)} \gtrsim r^{5-\varepsilon} \text { for some } \varepsilon>0
\end{aligned}
$$

[^1](see Lemma 4.9). Let us note that obtaining this almost monotonicity is much more involved than in the previous case (when we had $4-\varepsilon$ instead of $5-\varepsilon$ ). The reason is technical and rather delicate: we need to show that the size of $\mathrm{W}_{3, r}$ controls the one of its gradient (see Lemma 4.7), and this follows from a semiconvexity estimate along some rotational derivatives.

Once this almost monotonicity is proved, we can consider accumulation points of $\tilde{\mathrm{W}}_{3, r}$ at all points where $\left\|\mathrm{W}_{3}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \gtrsim r^{5-\varepsilon}$, and we prove that (up to a codimension 2 set) the only possible limit is a harmonic polynomial $p_{4, x_{0}}$ of degree 4 . This leads to an expansion of the form

$$
w_{4}:=u\left(x_{0}+\cdot\right)-\mathscr{P}_{x_{0}}-p_{4, x_{0}}, \quad w_{4}(x)=o\left(x^{4}\right) .
$$

Hence, to finally obtain (1.5) with $\mathrm{P}_{4, x_{0}}=\mathscr{P}_{x_{0}}+p_{4, x_{0}}$, we only need to prove that $w_{4}(x)=$ $\mathrm{O}\left(|x|^{5-\zeta}\right)$ for all $\zeta>0$, up to a set of dimension $n-2$. This is again nontrivial: indeed, we can show that $r \mapsto\left\|\nabla \tilde{w}_{4, r}\right\|_{L^{2}\left(B_{1}\right)}$ is almost increasing provided that $p_{4, x_{0}}$ vanishes on $\left\{p_{2, x_{0}}=0\right\}$ and $\left\|w_{4}(r \cdot)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \gtrsim r^{5-\varepsilon}$ for some $\varepsilon>0$. Hence we need to ensure that these assumptions are satisfied in a large enough set, and for this we exploit some recent results on the size of the zero set of harmonic functions (see [NV17]) and another dimension reduction argument. In this way, we conclude the validity of (1.5).

It is important to remark here that the expansion (1.5) up to order $5-\zeta$ is exactly what we need in order to prove Theorem 1.1. Even if one could improve such an expansion to a higher order, the estimate on the singular set in Theorem 1.1 would not change. Indeed, the bounds on the sizes of the sets where the expansions stop before $5-\zeta$ would not improve, and the conclusion in Theorem 1.1 would remain exactly the same (see also Remark 9.6).
1.3.3. From one solution to a monotone 1-parameter family of solutions. - Note that the analysis performed above holds only for one solution. If now we have an increasing family of solutions $\left\{u^{t}\right\}_{t \in(-1,1)}$, we need to understand for each $t$ the size of points where the expansion stops at some fixed order $\lambda \in \Lambda$.

If one simply applies the previous analysis to each solution $u^{t}$ one would not be able to conclude. Indeed, if for each solution the set $\Sigma^{t, \lambda}$ has dimension bounded by some $s \geq 0$, then a simple argument (using the structure of our problem) would show that their union over $t \in(-1,1)$ has dimension bounded by $s+1$. Unfortunately, this estimate would be absolutely too weak for our scope. Indeed, in order to prove our result, we need to show the following: if the analysis performed on a single solution implies that a set $\Sigma^{t, \lambda}$ has dimension bounded by $s$, then also $\bigcup_{t \in(-1,1)} \Sigma^{t, \lambda}$ has dimension bounded by $s$. In other words, the bound on the union should be exactly the same as the one obtained for each single set!

To achieve this, we have to exploit the fact that we have an increasing family $u^{t}$ of solutions to obtain very refined estimates on the possible blow-ups of a fixed solution at a free boundary point. More precisely, the idea is the following: if a sequence of singular
points $x_{k} \in \Sigma^{t_{k}}$ converges to $0 \in \Sigma^{0}$, and if both solutions $u^{t_{k}}$ and $u^{0}$ have a Taylor expansion up to the same order $\lambda$ at these points, then this implies some extra information on the possible Taylor expansion of $u^{0}$ at 0 (more precisely, this implies some symmetry properties on its higher order term). This analysis is performed in Section 6 and it introduces a complete series of new ideas and techniques with respect to [FS19], where only one fixed solution was considered. We want to emphasize that, with respect to the case of only one solution (where one can still deduce symmetry properties of blow-ups as a consequence of being an accumulation point of other singular points), in this case this analysis is made particularly delicate by the fact that we do not have any equation in $t$ relating the solutions: we only know that they are ordered and strictly increasing with respect to $t$. Still, we are able to deduce some strong symmetry properties of blow-up at all points where the Almgren's frequency is continuous (see the results in Section 6), and from these properties we obtain a very precise description of the structure of singular points. This is then combined with a series of covering and dimension reduction arguments (see Sections 7 and 8) to estimate the size of the singular points where blow-ups have few symmetries, which allow us to show the desired dimensional bounds on $\bigcup_{t \in(-1,1)} \Sigma^{t, \lambda}$. To our knowledge, this is the first dimension reduction argument for a 1-parameter family of solutions to elliptic PDEs, and we expect these ideas and techniques to be useful in many other problems.
1.3.4. Extra comments. - The previous description syntheses well the main ideas behind our strategy, and what explained until now suffices for proving the Schaeffer conjecture in $\mathbf{R}^{3}$. However, for the $\mathbf{R}^{4}$ case, other extra ideas (that we only briefly mentioned before) are required. In particular, we need a more refined analysis that depends on the type of singular point, having to distinguish two cases:
(a) The case of the "lower dimensional strata". - For singular points in $\bigcup_{t \in(-1,1)} \Sigma_{m}^{t}$ with $m \leq n-2$, we know that the expansion (1.6) stops at $\lambda=2$ at "anomalous points", and these points have dimension bounded by $m-1$. Hence, denoting this set of anomalous points by $\bigcup_{t \in(-1,1)} \Sigma_{m}^{t, 2}$, and the remaining $m$-dimensional set (where the expansion stops at $\lambda=3$ ) by $\bigcup_{t \in(-1,1)} \Sigma_{m}^{t, 3}$, applying (1.8) for $n=4$ and $m=2$ we get the trivial bound

$$
\operatorname{dim}_{\mathcal{H}}\left(\left\{t \in(-1,1): \Sigma_{m}^{t, \lambda} \neq \emptyset\right\}\right) \leq \frac{\operatorname{dim}_{\mathcal{H}}\left(\bigcup_{t \in(-1,1)} \Sigma_{m}^{t, \lambda}\right)}{\lambda-1}=1,
$$

$$
\text { for } \lambda=2,3 .
$$

To improve this estimate, we refine our barrier arguments and show that (1.7) can be improved to $\Sigma_{m}^{t_{t}+r^{\lambda-\varepsilon}, i} \cap \mathrm{~B}_{r}\left(x_{0}\right)=\emptyset$ for any $\varepsilon>0$, for $\lambda=2,3$. This increased speed at which the contact set clears near singular points for "lower strata" is not difficult to prove, but is fundamental to establish Schaeffer's conjecture in $\mathbf{R}^{4}$.
(b) The case of the "top stratum". - A big difficulty arises from points in $\bigcup_{t \in(-1,1)} \Sigma_{n-1}^{t}$ where the expansion stops at $\lambda=3$. More precisely, as mentioned in Section 1.3.2(b) when discussing the case where the third blow-up is not harmonic, there may exist a set $\bigcup_{t \in(-1,1)} \Sigma_{n-1}^{t, 3}$ of dimension $n-2$ at which the expansion $u^{t_{\circ}}(x)=p_{2, x_{0}, t_{0}}\left(x-x_{0}\right)+\mathrm{O}(\mid x-$ $\left.x_{0}| |^{3}\right)$ is sharp-i.e., the third derivatives of $u^{t_{0}}$ do not exist at those points. Then, if we use (1.8) for $\lambda=3$ and $n=4$, we obtain the (trivial) estimate

$$
\operatorname{dim}_{\mathcal{H}}\left(\left\{t \in(-1,1): \Sigma_{n-1}^{t, 3} \neq \emptyset\right\}\right) \leq \frac{n-2}{\lambda-1}=1
$$

while to establish Schaeffer's conjecture in $\mathbf{R}^{4}$ we need to prove $\mathcal{H}^{1}\left(\left\{t \in(-1,1): \Sigma^{t, 3} \neq\right.\right.$ $\emptyset\})=0$. Since it is possible to construct examples where the set $\bigcup_{t \in(-1,1)} \Sigma_{n-1}^{t, 3}$ is $(n-2)$ dimensional, there is no hope to improve its dimensional bound. In addition, unlike in the "lower strata", one can see that (1.7) for $\lambda=3$ is sharp at these points. Consequently, we need a completely different argument to conclude.

Here the idea consists in taking negative increments of $t$-which make the contact set become thicker instead of thinner-and use barrier arguments to show that all free boundary points in a neighborhood become regular at a slightly enhanced speed (see Lemma 9.4). However, we can only take advantage of this improvement if we can prove that the set $\bigcup_{t \in(-1,1)} \Sigma_{n-1}^{t, 3}$ is ( $n-2$ )-rectifiable (the information that its Hausdorff dimension is bounded by $n-2$ is not sufficient). This rectifiability result is crucial, and its proof relies on the existence of $\lim _{r \downarrow 0} \tilde{w}_{2, r}$ in the non-harmonic case. As mentioned before, this fact requires completely new ideas with respect to the classical tools known in this kind of problems, and it is the focus of Section 5. It is worth observing that such arguments lead to new interesting results even when applied to the Signorini problem, see Appendix B.
1.4. Organization of the paper. - The paper is organized as follows. In Section 2 we introduce a series useful functionals that will be used in the proof of some of our monotonicity formulae. In Section 3 we present some preliminary results that will be needed throughout the paper. Then, in Section 4 we develop our higher order analysis of singular points. In Section 5 we study in detail singular points at which the blow up of $u\left(x_{\circ}+\cdot\right)-p_{2, x_{0}}$ is 3 -homogeneous (cf. Section 1.3.2). In Section 6 we consider a 1parameter family of solutions to the obstacle problem and study symmetry properties of their blow-ups. In Section 7 we prove a series of lemmas of geometric measure theorytype, and in Section 8 we establish several dimension reduction results. Finally, in Section 9 we prove our main result, Theorem 1.1, as well as Theorem 1.2. In addition, at the end of the paper we provide two appendices on the Signorini problem: one with some basic results that are needed throughout the paper, and a second one with new results (uniqueness and nondegeneracy of blow-ups at odd frequencies) that are a consequence of our analysis in Section 5 and that we believe to be of independent interest.

## 2. Useful functionals and formulae

In all this section $r \in(0,1)$, and $w: \mathrm{B} \rightarrow \mathbf{R}$ denotes an arbitrary $\mathrm{C}^{1,1}$ function defined in a ball $\mathrm{B} \subset \mathbf{R}^{n}$ (specified in each statement). Throughout the paper we shall use the following dimensionless quantities:

$$
\begin{aligned}
\mathrm{D}(r, w) & :=r^{2-n} \int_{\mathbf{B}_{r}}|\nabla w|^{2}=r^{2} \int_{\mathbf{B}_{1}}|\nabla w|^{2}(r \cdot), \\
\mathrm{H}(r, w) & :=r^{1-n} \int_{\partial \mathbf{B}_{r}} w^{2}=\int_{\partial \mathbf{B}_{1}} w^{2}(r \cdot) .
\end{aligned}
$$

We also introduce here a useful notation for rescaling and normalization. Given $w: \mathrm{B} \rightarrow \mathbf{R}$ and $r>0$ we define $w_{r}$ and $\tilde{w}_{r}$ as

$$
\begin{equation*}
w_{r}(x):=w(r x) \quad \text { and } \quad \tilde{w}_{r}(x):=\frac{w_{r}}{\mathrm{H}\left(1, w_{r}\right)^{\frac{1}{2}}}=\frac{w(r \cdot)}{\mathrm{H}(r, w)^{\frac{1}{2}}} . \tag{2.1}
\end{equation*}
$$

We start by computing the derivatives of H and D .
Lemma 2.1. - Let $w \in \mathrm{C}^{1,1}\left(\mathrm{~B}_{2}\right)$. Then

$$
\begin{equation*}
\left.\frac{d}{d r}\right|_{r=1} \mathrm{H}(r, w)=2 \int_{\partial \mathrm{B}_{1}} w w_{v}=2 \int_{\mathbf{B}_{1}} w \Delta w+2 \int_{\mathbf{B}_{1}}|\nabla w|^{2} . \tag{2.2}
\end{equation*}
$$

Proof. - This is a standard computation, that can be found for instance in [FS19].

Lemma 2.2. - Let $w \in \mathrm{C}^{1,1}\left(\mathrm{~B}_{2}\right)$. Then

$$
\begin{equation*}
\left.\frac{d}{d r}\right|_{r=1} \mathrm{D}(r, w)=2 \int_{\partial \mathrm{B}_{1}} w_{v}^{2}-2 \int_{\mathrm{B}_{1}} \Delta w(x \cdot \nabla w) . \tag{2.3}
\end{equation*}
$$

Proof. - For convenience, we set $\mathrm{D}(r):=\mathrm{D}(r, w)$. It holds

$$
\begin{aligned}
\mathrm{D}^{\prime}(1) & =\sum_{i, j} \int_{\mathrm{B}_{1}} 2 w_{i} x_{j} w_{i j}+2 \mathrm{D}(1) \\
& =\sum_{i, j} \int_{\partial \mathrm{B}_{1}} 2 w_{i} x_{j} w_{j} v_{i}-\sum_{i, j} \int_{\mathrm{B}_{1}} 2\left(w_{i} x_{j}\right)_{i} w_{j}+2 \mathrm{D}(1)
\end{aligned}
$$

$$
\begin{aligned}
& =2 \int_{\partial \mathrm{B}_{1}} w_{v}^{2}-2 \int_{\mathrm{B}_{1}} \Delta w(x \cdot \nabla w)-2 \int_{\mathrm{B}_{1}}|\nabla w|^{2}+2 \mathrm{D}(1) \\
& =2 \int_{\partial \mathrm{B}_{1}} w_{v}^{2}-2 \int_{\mathrm{B}_{1}} \Delta w(x \cdot \nabla w)
\end{aligned}
$$

Let us introduce the functions

$$
\begin{equation*}
\phi(r, w):=\frac{\mathrm{D}(r, w)}{\mathrm{H}(r, w)}, \quad \phi^{\gamma}(r, w):=\frac{\mathrm{D}(r, w)+\gamma r^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} . \tag{2.4}
\end{equation*}
$$

The quantity $\phi$ is often known as the Almgren frequency function. Instead, $\phi^{\gamma}$ is a new truncated frequency function, that to our knowledge has never been introduced before. It will be used throughout the paper and will be extremely useful in our arguments, as it will allow us to deal with the cases when H may be too small. ${ }^{2}$ The choice of the constant $\gamma$ in front of $r^{2 \gamma}$ in the numerator is important to make the following lemma work.

Lemma 2.3. - Let $w \in \mathrm{C}^{1,1}\left(\mathrm{~B}_{1}\right)$. Then for $r \in(0,1)$ we have

$$
\frac{d}{d r} \phi(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}+\mathrm{E}(r, w)}{(\mathrm{H}(r, w))^{2}}
$$

and

$$
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}+\mathrm{E}^{\gamma}(r, w)}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}},
$$

where

$$
\begin{equation*}
\mathrm{E}(r, w):=\left(r^{2-n} \int_{\mathbf{B}_{r}} w \Delta w\right) \mathrm{D}(r, w)-\left(r^{2-n} \int_{\mathbf{B}_{r}}(x \cdot \nabla w) \Delta w\right) \mathrm{H}(r, w) \tag{2.5}
\end{equation*}
$$

and
(2.6)

$$
\begin{aligned}
\mathrm{E}^{\gamma}(r, w): & \left(r^{2-n} \int_{\mathbf{B}_{r}} w \Delta w\right)\left(\mathrm{D}(r, w)+\gamma r^{2 \gamma}\right) \\
& -\left(r^{2-n} \int_{\mathbf{B}_{r}}(x \cdot \nabla w) \Delta w\right)\left(\mathrm{H}(r, w)+r^{2 \gamma}\right) .
\end{aligned}
$$

[^2]Proof. - We first observe that, for $r \in(0,1)$, the formula for $\phi$ can be deduced from the one for $\phi^{\gamma}$ letting $\gamma \uparrow+\infty$.

By scaling it is enough to compute, for $a>0$,

$$
\left.\frac{d}{d r}\right|_{r=1} \log \phi^{\gamma, a}(r, w), \quad \text { for } \phi^{\gamma, a}(r, w):=\frac{\mathrm{D}(r, w)+\gamma(a r)^{2 \gamma}}{\mathrm{H}(r, w)+(a r)^{2 \gamma}} .
$$

Using Lemmas 2.1 and 2.2 we obtain

$$
\frac{\left.\frac{d}{d r}\right|_{r=1}\left(\mathrm{D}(r, w)+\gamma(a r)^{2 \gamma}\right)}{\mathrm{D}(1, w)+\gamma a^{2 \gamma}}=2 \frac{\int_{\partial \mathrm{B}_{1}} w_{v}^{2}-\int_{\mathrm{B}_{1}}(x \cdot \nabla w) \Delta w+\gamma^{2} a^{2 \gamma}}{\int_{\partial \mathrm{B}_{1}} w w_{v}-\int_{\mathrm{B}_{1}} w \Delta w+\gamma a^{2 \gamma}}
$$

and

$$
\frac{\left.\frac{d}{d r}\right|_{r=1}\left(\mathrm{H}(r, w)+a r^{2 \gamma}\right)}{\mathrm{H}(1, w)+a^{2 \gamma}}=2 \frac{\int_{\partial \mathrm{B}_{1}} w w_{v}+\gamma a^{2 \gamma}}{\int_{\partial \mathrm{B}_{1}} w^{2}+a^{2 \gamma}} .
$$

Therefore,

$$
\begin{aligned}
\frac{d}{d r} \log \phi^{\gamma, a}(1, w) & =\frac{\left.\frac{d}{d r}\right|_{r=1}\left(\mathrm{D}(r, w)+\gamma(a r)^{2 \gamma}\right)}{\mathrm{D}(1, w)+\gamma a^{2 \gamma}}-\frac{\left.\frac{d}{d r}\right|_{r=1}\left(\mathrm{H}(r, w)+a r^{2 \gamma}\right)}{\mathrm{H}(1, w)+a^{2 \gamma}} \\
& =2 \frac{\mathrm{X}^{2}+\text { rest }}{\left(\mathrm{D}(1, w)+\gamma a^{2 \gamma}\right)\left(\mathrm{H}(1, w)+a^{2 \gamma}\right)},
\end{aligned}
$$

where

$$
\mathrm{X}^{2}:=\left(\int_{\partial \mathrm{B}_{1}} w_{v}^{2}+\gamma^{2} a^{2 \gamma}\right)\left(\int_{\partial \mathrm{B}_{1}} w^{2}+a^{2 \gamma}\right)-\left(\int_{\partial \mathrm{B}_{1}} w w_{v}+\gamma a^{2 \gamma}\right)^{2} \geq 0
$$

(the non-negativity of $\mathrm{X}^{2}$ follows by Hölder inequality), and

$$
\text { rest }:=\left(\int_{\mathrm{B}_{1}} w \Delta w\right)\left(\int_{\partial \mathrm{B}_{1}} w w_{\nu}+\gamma a^{2 \gamma}\right)-\left(\int_{\mathrm{B}_{1}}(x \cdot \nabla w) \Delta w\right)\left(\int_{\partial \mathrm{B}_{1}} w^{2}+a^{2 \gamma}\right) .
$$

Using again Lemma 2.1 we have

$$
\begin{aligned}
\text { rest }= & \left(\int_{\mathbf{B}_{1}} w \Delta w\right)^{2}+\left(\int_{\mathbf{B}_{1}} w \Delta w\right)\left(\mathrm{D}(1, w)+\gamma a^{2 \gamma}\right) \\
& -\left(\int_{\mathbf{B}_{1}}(x \cdot \nabla w) \Delta w\right)\left(\mathrm{H}(1, w)+a^{2 \gamma}\right) .
\end{aligned}
$$

By scaling (applying the previous formulas with $w$ replaced by $w_{r}=w(r \cdot)$ and $a$ replaced by $r$ ) we obtain

$$
\frac{d}{d r} \log \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}+\mathrm{E}^{\gamma}(r, w)}{\left(\mathrm{D}(r, w)+\gamma r^{2 \gamma}\right)\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)},
$$

where $\mathrm{E}^{\gamma}$ is defined in (2.6). Since $\frac{d}{d r} \log \phi^{\gamma}(r, w)=\phi^{\gamma}(r, w)^{-1} \frac{d}{d r} \phi^{\gamma}(r, w)$, the lemma follows immediately by recalling the definition of $\phi^{\gamma}$ in (2.4).

## 3. Preliminaries: first and second blow-up analysis

In this section we collect some known results and basic tools that will be used throughout the paper. Let $u: \mathrm{B}_{1} \rightarrow \mathbf{R}$ be a solution to the obstacle problem

$$
\begin{equation*}
\Delta u=\chi_{\{u>0\}} \quad \text { and } \quad u \geq 0 \quad \text { in } \mathrm{B}_{1} . \tag{3.1}
\end{equation*}
$$

By the classical theory of Caffarelli on the obstacle problem [Caf77, Caf98], any solution $u$ of (3.1) with $0 \in \partial\{u>0\}$ satisfies
(3.2) $\quad\|u\|_{\mathrm{C}^{1,1}\left(\mathbf{B}_{1 / 2}\right)} \leq \mathrm{C} \quad$ and $\quad \sup _{\mathbf{B}_{r}(0)} u \geq c r^{2} \quad \forall r \in\left(0, \frac{1}{2}\right)$,
where $\mathrm{C}, c>0$ are positive dimensional constants. Moreover, points of the free boundary $\partial\{u>0\}$ can be split into two classes:

- Regular points: $x_{\circ} \in \partial\{u>0\}$ is a regular point if

$$
\lim _{r \rightarrow 0} r^{-2} u\left(x_{\circ}+r x\right)=\frac{1}{2}(\max \{0, \boldsymbol{e} \cdot x\})^{2}
$$

for some $\boldsymbol{e} \in \mathbf{S}^{n-1}$.

- Singular points: $x_{\circ} \in \partial\{u>0\}$ is a singular point if

$$
p_{2, x_{0}}(x):=\lim _{r \rightarrow 0} r^{-2} u\left(x_{\circ}+r x\right)
$$

exists and $p_{2, x_{0}}$ is a quadratic polynomial belonging to the set

$$
\mathcal{P}:=\{\text { convex 2-homogeneous polynomials } p \text { with } \Delta p \equiv 1\}
$$

It is well known that the free boundary is analytic in a neighborhood of regular points. So, the main goal is to understand the structure of singular points.

When $x_{\circ}=0$ is a singular point, we will simplify the notation $p_{2, x_{0}}$ to $p_{2}$. A well known result due to Caffarelli is the following estimate at singular points.

Lemma 3.1. - There exists a modulus of continuity $\omega: \mathbf{R}^{+} \rightarrow \mathbf{R}^{+}$, depending only on the dimension $n$, such that if $u$ is a solution of the obstacle problem (3.1) and 0 is a singular free boundary point, then

$$
\left\|u-p_{2}\right\|_{\mathrm{L}^{\infty}\left(\mathbf{B}_{r}\right)} \leq r^{2} \omega(r), \quad \forall r \in(0,1)
$$

Proof. - This result, with an abstract (dimensional) modulus of continuity $\omega$, is contained in [Caf98, Theorem 8]. A stronger quantitative version of the estimate (with independent proofs) giving an explicit $\mathrm{C}|\log r|^{-\varepsilon}$ modulus of continuity is given in [CSV18, FS19].

Remark 3.2. - Let $p \in \mathcal{P}$. Since $\Delta u=\Delta p=1$ in $\{u>0\}$, we have

$$
\begin{equation*}
(u-p) \Delta(u-p)=p \chi_{\{u=0\}} \geq 0 \tag{3.3}
\end{equation*}
$$

Similarly,

$$
\begin{equation*}
x \cdot \nabla(u-p) \Delta(u-p)=x \cdot \nabla p \chi_{\{u=0\}}=2 p \chi_{\{u=0\}} \geq 0 \tag{3.4}
\end{equation*}
$$

We recall Weiss' monotonicity formula (proved in [W99] for $\lambda=2$, and in [FS19] in the general case) and a useful consequence of it.

Lemma 3.3 (Weiss' formula). - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1), and let 0 be a singular point. Given $p \in \mathcal{P}$, set $w:=u-p$. Also, for $\lambda>0$ set

$$
\mathrm{W}_{\lambda}(r, w):=r^{-2 \lambda}(\mathrm{D}(r, w)-\lambda \mathrm{H}(r, w)) .
$$

Then:
(a) For all $\lambda \geq 2$

$$
\frac{d}{d r} \mathrm{~W}_{\lambda}(r, w) \geq 2 r^{-2 \lambda-1} \int_{\partial \mathrm{B}_{r}}(x \cdot \nabla w-\lambda w)^{2} \geq 0
$$

(b) $\mathrm{W}_{2}\left(0^{+}, w\right)=0$ and

$$
\begin{equation*}
\mathrm{D}(r, w)-2 \mathrm{H}(r, w) \geq 0, \quad \forall r \in(0,1) . \tag{3.5}
\end{equation*}
$$

Proof. - (a) By scaling it is enough to compute $\frac{d}{d r} \mathrm{~W}_{\lambda}(r, w)$ at $r=1$. Using Lemmas 2.1-2.2, we obtain

$$
\begin{aligned}
\mathrm{W}_{\lambda}^{\prime}(1, w) & =\left(\mathrm{D}^{\prime}(1, w)-\lambda \mathrm{H}^{\prime}(1, w)\right)-2 \lambda \mathrm{D}(1, w)+2 \lambda^{2} \mathrm{H}(1, w) \\
& =2 \int_{\partial \mathrm{B}_{1}} w_{v}^{2}-2 \int_{\mathrm{B}_{1}} \Delta w(x \cdot \nabla w)
\end{aligned}
$$

$$
\begin{aligned}
& -2 \lambda \int_{\partial \mathrm{B}_{1}} w w_{v}-2 \lambda \mathrm{D}(1, w)+2 \lambda^{2} \mathrm{H}(1, w) \\
= & 2 \int_{\partial \mathrm{B}_{1}} w_{v}^{2}+2 \int_{\mathrm{B}_{1}}(\lambda w-x \cdot \nabla w) \Delta w-4 \lambda \int_{\partial \mathrm{B}_{1}} w w_{v}+2 \lambda^{2} \int_{\partial \mathrm{B}_{1}} w^{2} \\
= & 2 \int_{\partial \mathrm{B}_{1}}\left(w_{v}-\lambda w\right)^{2}+2 \int_{\mathrm{B}_{1}}(\lambda w-x \cdot \nabla w) \Delta w \\
= & 2 \int_{\partial \mathrm{B}_{1}}\left(w_{v}-\lambda w\right)^{2}+2 \int_{\{u(r \cdot)=0\} \cap \mathrm{B}_{1}}(\lambda p-x \cdot \nabla p) .
\end{aligned}
$$

One concludes noticing that, for $\lambda \geq 2$, it holds $(\lambda p-x \cdot \nabla p)=(\lambda-2) p \geq 0$.
(b) Since 0 is a singular point then $w_{r}(x)=(u-p)(r x)=\left(p_{2}-p\right)(r x)+o\left(r^{2}\right)$, thus

$$
\mathrm{W}_{2}\left(0^{+}, w\right)=\lim _{r \downarrow 0} \mathrm{~W}\left(1, r^{-2} w_{r}\right)=\mathrm{W}_{2}\left(1, p_{2}-p\right)=0
$$

As a consequence, (3.5) follows integrating (a) for $\lambda=2$ between 0 and $r$.
We recall from [FS19] that the frequency function $\phi$ applied to the function $w=$ $u-p$, with $p \in \mathcal{P}$, is monotone increasing in $r$. More precisely we have the following:

Proposition 3.4 (Frequency formula). - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1), and let 0 be a singular point. Given $p \in \mathcal{P}$, set $w:=u-p$. Then

$$
\frac{d}{d r} \phi(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\mathrm{H}(r, w)^{2}} \geq 0, \quad \forall r \in(0,1) .
$$

Proof. - By Lemma 2.3 we just need to show that

$$
\mathrm{E}(r, w):=\left(r^{2-n} \int_{\mathbf{B}_{r}} w \Delta w\right) \mathrm{D}(r, w)-\left(r^{2-n} \int_{\mathbf{B}_{r}}(x \cdot \nabla w) \Delta w\right) \mathrm{H}(r, w) \geq 0 .
$$

Using Remark 3.2 for $w=u-p$ we have

$$
\left(r^{2-n} \int_{\mathbf{B}_{r}}(x \cdot \nabla w) \Delta w\right)=2\left(r^{2-n} \int_{\mathbf{B}_{r}} w \Delta w\right),
$$

thus

$$
\mathrm{E}(r, w)=\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)(\mathrm{D}(r, w)-2 \mathrm{H}(r, w))
$$

which by (3.5) and Remark 3.2 is nonnegative.

The following observation, also contained in [FS19], follows immediately from (3.5).

Lemma 3.5. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1), and let 0 be a singular point. Given $p \in \mathcal{P}$, set $w:=u-p$. Then $\phi\left(0^{+}, w\right) \geq 2$.

A new important estimate that we will use throughout the paper is the following:
Lemma 3.6. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1), and let 0 be a singular point. Given $p \in \mathcal{P}$, set $w:=u-p$. Suppose that for $0<r<\mathrm{R}<1$ we have $\underline{\lambda} \leq \phi(r, w) \leq$ $\phi(\mathrm{R}, w) \leq \bar{\lambda}$. Then, for any given $\delta>0$ we have

$$
\left(\frac{\mathrm{R}}{r}\right)^{2 \underline{\lambda}} \leq \frac{\mathrm{H}(\mathrm{R}, w)}{\mathrm{H}(r, w)} \leq \mathrm{C}_{\delta}\left(\frac{\mathrm{R}}{r}\right)^{2 \bar{\lambda}+\delta}
$$

where $\mathrm{C}_{\delta}$ depends only on $n, \bar{\lambda}, \delta$.
Proof. - Define

$$
\mathrm{F}(r):=\frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)}
$$

By Proposition 3.4 we have

$$
\begin{equation*}
\frac{d}{d r} \phi(r, w) \geq \frac{2}{r}(\mathrm{~F}(r))^{2} \tag{3.6}
\end{equation*}
$$

On the other hand, thanks to Lemma 2.1,

$$
\begin{equation*}
\frac{\frac{d}{d r} \mathrm{H}(r, w)}{\mathrm{H}(r, w)}=\frac{2}{r} \frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w+r^{2-n} \int_{\mathrm{B}_{r}}|\nabla w|^{2}}{\mathrm{H}(r, w)}=\frac{2}{r} \phi(r, w)+\frac{2}{r} \mathrm{~F}(r) . \tag{3.7}
\end{equation*}
$$

Integrating (3.6) and using Cauchy-Schwartz inequality, since $\underline{\lambda} \leq \phi(\rho, w) \leq \bar{\lambda}$ for all $\rho \in(r, \mathbf{R})$ we get

$$
\begin{aligned}
(\bar{\lambda} & -\underline{\lambda})^{1 / 2}(\log (\mathrm{R} / r))^{1 / 2} \\
& \geq\left(\int_{r}^{\mathrm{R}} \frac{d}{d \rho} \phi(\rho, w) d \rho\right)^{1 / 2}\left(\int_{r}^{\mathrm{R}} \frac{d \rho}{\rho}\right)^{1 / 2} \\
& =\left(\int_{r}^{\mathrm{R}} \frac{1}{\rho}(\mathrm{~F}(\rho))^{2} d \rho\right)^{1 / 2}\left(\int_{r}^{\mathrm{R}} \frac{d \rho}{\rho}\right)^{1 / 2} \geq \int_{r}^{\mathrm{R}} \frac{1}{\rho} \mathrm{~F}(\rho) d \rho \geq 0 .
\end{aligned}
$$

Hence, integrating (3.7), we obtain

$$
\begin{aligned}
\log \frac{\mathrm{H}(\mathrm{R}, w)}{\mathrm{H}(r, w)} & \leq \int_{r}^{\mathrm{R}} \frac{2}{\rho}(\bar{\lambda}+\mathrm{F}(\rho)) d \rho \leq \log \left((\mathrm{R} / r)^{2 \bar{\lambda}}\right)+\mathrm{C}(\log (\mathrm{R} / r))^{1 / 2} \\
& \leq \log \left((\mathrm{R} / r)^{2(\bar{\lambda}+\delta)}\right)+\mathrm{C}
\end{aligned}
$$

where C depends only on $n, \bar{\lambda}$, and $\delta$.
For the lower bound we recall that, since $w \Delta w \geq 0$, we have $\mathrm{F}(\rho) \geq 0$. Therefore, integrating (3.7) over [ $r, \mathrm{R}]$,

$$
\log \frac{\mathrm{H}(\mathrm{R}, w)}{\mathrm{H}(r, w)} \geq 2 \underline{\lambda} \int_{r}^{\mathrm{R}} \frac{d \rho}{\rho}=\log (\mathrm{R} / r)^{2 \underline{\lambda}} .
$$

We will also need the following result, which allows us to control the $L^{\infty}$ norm of the difference of two solutions with the $L^{2}$ norm.

Lemma 3.7. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ and $v: \mathrm{B}_{1} \rightarrow[0, \infty)$ be solutions of the obstacle problem (3.1). Then

$$
\|u-v\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{1 / 2}\right)} \leq \mathrm{C}(n)\|u-v\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)} .
$$

Proof. - On the one hand, from

$$
\Delta(u-v)=1-\Delta v \geq 0 \quad \text { in }\{u>0\} \quad \text { and } \quad u-v=-v \leq 0 \quad \text { in }\{u=0\}
$$

we obtain that $(u-v)_{+}=\max (u-v, 0)$ is subharmonic in $B_{1}$. Exchanging the role of $u$ and $v$, the same argument shows that $(v-u)_{+}=(u-v)_{-}$is subharmonic. Thus also $|u-v|=(u-v)_{+}+(u-v)_{-}$is subhamonic in $\mathrm{B}_{1}$, and using the mean value property we obtain, for $x \in \mathrm{~B}_{1 / 2}$,

$$
|u-v|(x) \leq \int_{\mathrm{B}_{1 / 2}(x)}|u-v| \leq \mathrm{C}(n)\|u-v\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{1}\right)} \leq \mathrm{C}(n)\|u-v\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)} .
$$

We now start investigating the structure of possible second blow-ups. The next few results are a small modification of those in [FS19].

The following Lipschitz estimate for the rescaled difference $u-p$, with $p \in \mathcal{P}$, will be useful in the sequel. We recall that $w_{r}$ and $\tilde{w}_{r}$ have been defined in (2.1).

Lemma 3.8. -Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1) with $u \not \equiv p_{2}$, and let 0 be a singular point. Given $p \in \mathcal{P}$, set $w:=u-p$. Let $\mathrm{R} \geq 1$, and $r \in\left(0, \frac{1}{10 \mathrm{R}}\right)$. Then

$$
\begin{equation*}
\partial_{e \boldsymbol{e}} \tilde{w}_{r} \geq-\mathrm{C} \quad \text { in } \mathrm{B}_{\mathrm{R}}, \quad \forall \boldsymbol{e} \in \mathbf{S}^{n-1} \cap\{p=0\} \tag{3.8}
\end{equation*}
$$

where C depends only on $n, \mathrm{R}$, and $\phi(1 / 2, u-p)$. In addition, if $\operatorname{dim}(\{p=0\})=n-1$, then

$$
\begin{equation*}
\left|\nabla \tilde{w}_{r}\right| \leq \mathrm{C} \quad \text { in } \mathrm{B}_{\mathrm{R} / 2} \tag{3.9}
\end{equation*}
$$

where C depends only on $n, \mathrm{R}$, and $\phi\left(\frac{1}{2}, u-p\right)$.
Proof. - This proof is essentially contained in Step 3 from the Proof of Proposition 2.10 in [FS19]. However, since some small changes are needed in our setting, we reproduce the main steps for the convenience of the reader.

Given a function $f: \mathbf{R}^{n} \rightarrow \mathbf{R}$, a vector $\boldsymbol{e} \in \mathbf{S}^{n-1}$, and $h>0$, let

$$
\delta_{\boldsymbol{e}, h}^{2} f:=\frac{f(\cdot+h \boldsymbol{e})+f(\cdot-h \boldsymbol{e})-2 f}{h^{2}}
$$

denote a second order incremental quotient. For $\boldsymbol{e} \in\{p=0\} \cap \mathbf{S}^{n-1}$ we have $\delta_{\boldsymbol{e}, h}^{2} p \equiv 0$. Thus, since $\Delta u=1$ outside of $\{u=0\}$ and $\Delta u \leq 1$ in $\mathrm{B}_{1}$, we have

$$
\Delta\left(\delta_{\boldsymbol{e}, h}^{2} w\right)=\frac{\Delta u(\cdot+h \boldsymbol{e})+\Delta u(\cdot-h \boldsymbol{e})-2 \Delta u}{h^{2}} \leq 0 \quad \text { in } \mathrm{B}_{\mathrm{R}} \backslash\{u=0\}
$$

On the other hand, since $u \geq 0$ we have

$$
\delta_{e, h}^{2} w=\delta_{e, h}^{2} u(\cdot) \geq 0 \quad \text { in }\{u=0\} .
$$

As a consequence, the negative part of the second order incremental quotient $\left(\delta_{e, h}^{2} \tilde{w}_{r}\right)_{-}$ is subharmonic, and so is its limit $\left(\partial_{e e}^{2} \tilde{w}_{r}\right)_{-}$(recall that $u$ is semiconvex, and thus $\left(\delta_{e, h}^{2} \tilde{w}_{r}\right)_{-} \rightarrow\left(\partial_{e e}^{2} \tilde{w}_{r}\right)_{-}$a.e. as $h \rightarrow 0$ ). Therefore, by weak Harnack inequality (see for instance [CC95, Theorem 4.8(2)]) there exists $\varepsilon=\varepsilon(n) \in(0,1)$ such that

$$
\left\|\left(\partial_{e e}^{2} \tilde{w}_{r}\right)_{-}\right\|_{L^{\infty}\left(\mathrm{B}_{\mathrm{R}}\right)} \leq \mathrm{C}(n)\left(f_{\mathrm{B}_{\mathrm{R}}}\left(\partial_{e e} \tilde{w}_{r}\right)_{-}^{\varepsilon}\right)^{1 / \varepsilon} \leq \mathrm{C}(n)\left(f_{\mathrm{B}_{2 \mathrm{R}}}\left|\partial_{e e} \tilde{w}_{r}\right|^{\varepsilon}\right)^{1 / \varepsilon}
$$

Also, by standard interpolation inequalities, the $L^{\varepsilon}$ norm (here we use $\varepsilon<1$ ) can be controlled by the weak $\mathrm{L}^{1}$ norm, namely

$$
\left(\int_{\mathrm{B}_{2 \mathrm{R}}}\left|\partial_{e e} \tilde{w}_{r}\right|^{\varepsilon}\right)^{1 / \varepsilon} \leq \mathrm{C}(n, \mathrm{R}) \sup _{\theta>0} \theta\left|\left\{\left|\partial_{e e} \tilde{w}_{r}\right|>\theta\right\} \cap \mathrm{B}_{2 \mathrm{R}}\right| .
$$

Furthermore, by Calderón-Zygmund theory, the right hand side above is controlled by

$$
\left\|\Delta \tilde{w}_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{3 \mathrm{R}}\right)}+\left\|\tilde{w}_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{3 \mathrm{R}}\right)}
$$

Thus, since $\Delta w_{r} \leq 0$ in $\mathrm{B}_{3 \mathrm{R}},\left\|\Delta \tilde{w}_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{\mathrm{R}}\right)}$ is controlled by the $\mathrm{L}^{1}$ norm of $\tilde{w}_{r}$ inside $\mathrm{B}_{4 \mathrm{R}}$ : indeed, if $\chi$ is a smooth nonnegative cut-off function that is equal to 1 in $B_{3 \mathrm{R}}$ and vanishes outside $\mathrm{B}_{4 \mathrm{R}}$, then

$$
\begin{equation*}
\left\|\Delta \tilde{w}_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{3 \mathrm{R}}\right)} \leq-\int_{\mathrm{B}_{4 \mathrm{R}}} \chi \Delta \tilde{w}_{r}=-\int_{\mathrm{B}_{4 \mathrm{R}}} \Delta \chi \tilde{w}_{r} \leq \mathrm{C}(n, \mathrm{R}) \int_{\mathrm{B}_{4 \mathrm{R}} \backslash \mathrm{~B}_{3 \mathrm{R}}}\left|\tilde{w}_{r}\right| . \tag{3.10}
\end{equation*}
$$

Also, for $8 r \mathrm{R}<1$, as a consequence of Lemma 3.6 we have

$$
\mathrm{H}\left(4 \mathrm{R}, \tilde{w}_{r}\right) \leq \mathrm{CR}^{2 \phi(1, u-p)+1} \mathrm{H}\left(1, \tilde{w}_{r}\right)=\mathrm{CR}^{2 \phi\left(\frac{1}{2}, u-p\right)+1}
$$

and thus

$$
\left\|\tilde{w}_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4 \mathrm{R}}\right)} \leq \mathrm{C}\left(n, \mathrm{R}, \phi\left(\frac{1}{2}, u-p\right)\right) .
$$

In conclusion, we obtain

$$
\left\|\left(\partial_{e e} \tilde{w}_{r}\right)_{-}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{\mathrm{R}}\right)} \leq \mathrm{C}(n, \mathrm{R})\left\|\tilde{w}_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4 \mathrm{R}}\right)} \leq \mathrm{C}\left(n, \mathrm{R}, \phi\left(\frac{1}{2}, u-p\right)\right) .
$$

Finally note that, when $\{p=0\}$ is $(n-1)$-dimensional, as a consequence of (3.8) the tangential Laplacian of $\tilde{w}_{r}$ (in the directions of $\{p=0\}$ ) is uniformly bounded from below. Thus, since $\Delta \tilde{w}_{r} \leq 0$, we have

$$
\partial_{e^{\prime} e^{\prime}} \tilde{w}_{r} \leq \mathrm{C} \quad \text { in } \mathrm{B}_{\mathrm{R}}, \quad \text { for } \boldsymbol{e}^{\prime} \in\{p=0\}^{\perp} \text { with }\left|\boldsymbol{e}^{\prime}\right|=1
$$

where, as before, $\mathrm{C}=\mathrm{C}\left(n, \mathrm{R}, \phi\left(\frac{1}{2}, u-p\right)\right)$. Thanks to these semiconvexity and semiconcavity estimates, we deduce the Lipschitz bound (3.9).

The next result corresponds to [FS19, Proposition 2.10]. However the statement there has a small mistake (that anyhow does not affect any of the arguments in [FS19]) and for convenience we correct it here.

Proposition 3.9. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1) with $u \not \equiv p_{2}$, let 0 be a singular point, and set $w:=u-p_{2}$. Denote $m:=\operatorname{dim}\left(\left\{p_{2}=0\right\}\right) \in\{0,1,2, \ldots n-1\}$, and $\lambda^{2 n d}:=$ $\phi\left(0^{+}, w\right)$.

Then, for every sequence $r_{k} \downarrow 0$ there is a subsequence $r_{k_{\ell}}$ such that $\tilde{w}_{r_{k \ell}} \rightharpoonup q$ as $\ell \rightarrow \infty$ in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$, where $q \not \equiv 0$ is a $\lambda^{2 n d}$-homogeneous function satisfying the following:
(a) If $0 \leq m \leq n-2$ then $q$ is a harmonic polynomial, and in particular $\lambda^{2 n d} \in\{2,3,4, \ldots\}$. In addition, if $\lambda^{2 n d}=2$, then in some appropriate coordinates we have

$$
\begin{equation*}
p_{2}(x)=\frac{1}{2} \sum_{i=m+1}^{n} \mu_{i} x_{i}^{2} \quad \text { and } \quad q(x)=v \sum_{i=m+1}^{n} x_{i}^{2}-\sum_{j=1}^{m} v_{j} x_{j}^{2}, \tag{3.11}
\end{equation*}
$$

where $\mu_{i}, v>0$, and they satisfy $\sum_{i=m+1}^{n} \mu_{i}=1,(n-m) v=\sum_{j=1}^{m} v_{j}$, and $\left|v_{j}\right| \leq v$ for $a n y j=1, \ldots, m$.
(b) If $m=n-1$ then we have $\tilde{w}_{r_{k e}} \rightarrow q$ in $\mathrm{C}_{\mathrm{loc}}^{0}\left(\mathbf{R}^{n}\right)$ and we have $\lambda^{2 n d} \geq 2+\alpha_{o}$, where $\alpha_{\circ}>0$ is a dimensional constant. In addition, $q$ is a global solution of the Signorini problem:
(3.12)

$$
\left\{\begin{array}{lll}
\Delta q \leq 0 & \text { and } & q \Delta q=0 \\
\Delta q=0 & \text { in } \mathbf{R}^{n} \\
q \geq 0 & & \text { in } \mathbf{R}^{n} \backslash\left\{p_{2}=0\right\} \\
q \geq & \text { on }\left\{p_{2}=0\right\}
\end{array}\right.
$$

Proof. - The statement here is almost identical to that of [FS19, Proposition 2.10]. The only differences are the following:
(1) In [FS19] it is incorrectly stated that $\boldsymbol{v}_{j}>0$. Instead, [FS19, Lemma 2.12] proves that $v$ is the largest eigenvalue of $\mathrm{D}^{2} q$, hence the correct conclusion is that $\left|v_{j}\right| \leq v$ for $\operatorname{each} j=1, \ldots, m$.
(2) In the above statement we said that $\tilde{w}_{k} \rightharpoonup q$ weakly in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$, while [FS19, Proposition 2.10] states the convergence only in $W^{1,2}\left(B_{1}\right)$. The reason why we may replace $B_{1}$ by any larger ball is Lemma 3.6, as it allows us to control $H\left(R, \tilde{w}_{r}\right)$ by $\mathrm{C}(n, \phi(1, w), \mathrm{R}) \mathrm{H}\left(1, \tilde{w}_{r}\right)$ for any $r<1 / \mathrm{R}$. Hence, using a diagonal argument, the proof of [FS19, Proposition 2.10] yields the desired result.

We now recall another important estimate from [FS19]:
Proposition 3.10. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1) with $u \not \equiv p_{2}$, let 0 be a singular point, and set $w:=u-p_{2}, \lambda^{2 n d}:=\phi\left(0^{+}, w\right)$. Let $\lambda \in\left(0, \lambda^{2 n d}\right]$. Then

$$
\frac{\left|\{u=0\} \cap \mathrm{B}_{r}\right|}{\left|\mathrm{B}_{r}\right|} \leq \mathrm{C} r^{\lambda-2} \quad \forall r \in(0,1 / 2) .
$$

Moreover, if $\operatorname{dim}\left(\left\{p_{2}=0\right\}\right)=n-1$ then

$$
\{u=0\} \cap \mathrm{B}_{r} \subset\left\{x: \operatorname{dist}\left(x,\left\{p_{2}=0\right\}\right) \leq \mathrm{C} r^{\lambda-1}\right\}
$$

The constant C depends only on $n$ and $\lambda$.
Proof. - The first part is exactly [FS19, Proposition 2.13]. The second part on $\Sigma_{n-1}$ follows by the argument in [FS19, Remark 2.14], as a consequence of the Lipschitz estimate in Lemma 3.8.

Following the notation introduced in [FS19], we denote

$$
\begin{equation*}
\Sigma_{m}:=\left\{x_{\circ} \text { singular points with } \operatorname{dim}\left(\left\{p_{2, x_{0}}=0\right\}\right)=m\right\}, \quad 0 \leq m \leq n-1 \tag{3.13}
\end{equation*}
$$

and, for $m \leq n-2$,
(3.14)

$$
\Sigma_{m}^{a}:=\left\{x_{\circ} \in \Sigma_{m} \text { such that } \phi\left(0^{+}, u\left(x_{\circ}+\cdot\right)-p_{2, x_{0}}\right)=2\right\}, \quad 0 \leq m \leq n-2 .
$$

Moreover, for $m=n-1$ we introduce some further notation: we define

$$
\begin{equation*}
\Sigma_{n-1}^{<3}:=\left\{x_{\circ} \in \Sigma_{n-1} \text { such that } \phi\left(0^{+}, u\left(x_{\circ}+\cdot\right)-p_{2, x_{0}}\right)<3\right\}, \tag{3.15}
\end{equation*}
$$

and

$$
\begin{equation*}
\Sigma_{n-1}^{\geq 3}:=\Sigma_{n-1} \backslash \Sigma_{n-1}^{<3} . \tag{3.16}
\end{equation*}
$$

Finally, we will need the following:
Definition 3.11.-Let $u: \mathbf{B}_{1} \rightarrow[0, \infty)$ solve (3.1). For $1 \leq m \leq n-1$ we denote by $\Sigma_{m}^{3 r d}$ the set of points $x_{0} \in \Sigma_{m}$ such that, for $w:=u\left(x_{\circ}+\cdot\right)-p_{2, x_{0}}$, the following two conditions hold:
(i) $\phi\left(0^{+}, w\right) \geq 3$;
(ii) there exists some sequence $r_{k} \downarrow 0$ along which $r_{k}^{-3} w\left(r_{k} \cdot\right)$ converges, weakly in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$, to some 3-homogeneous harmonic polynomial vanishing on $\left\{p_{2, x_{0}}=0\right\}$-possibly the polynomial zero.

Notice that, by Proposition 3.9(a), for $m \leq n-2$ we have $\Sigma_{m} \backslash \Sigma_{m}^{a}=\Sigma_{m}^{3 r d}=\Sigma_{m}^{\geq 3}$. On the other hand, this is not true for $m=n-1$, and later on in the paper we will need to understand the set $\Sigma_{n-1}^{\geq 3} \backslash \Sigma_{n-1}^{3 r d}$.

We conclude this section by recalling that if $0 \in \Sigma_{m}^{3 r d}$ then the limit

$$
\begin{equation*}
\lim _{r \downarrow 0} r^{-3}\left(u-p_{2}\right)(r \cdot) \tag{3.17}
\end{equation*}
$$

exists. Indeed, as shown in [FS19], this is a consequence of the following Monneau-type monotonicity formula.

Lemma 3.12. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ satisfy (3.1), and let $0 \in \Sigma_{m}^{\geq 3}$ for some $0 \leq m \leq$ $n-1$. Let $w:=u-p_{2}-\mathrm{P}$, where P is any 3-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$. Then

$$
\begin{equation*}
\mathrm{D}(r, w) \geq 3 \mathrm{H}(r, w) \quad \forall r \in(0,1) \tag{3.18}
\end{equation*}
$$

and
(3.19)

$$
\frac{d}{d r}\left(r^{-6} \mathrm{H}(r, w)\right) \geq-\mathrm{C} \sup _{\partial \mathrm{B}_{1}} \frac{\mathrm{P}^{2}}{p_{2}}
$$

where C is some dimensional constant.
Proof. - The proof is contained in [FS19, Lemma 4.1].
The next result provides the existence of a unique limit in (3.17) for all points in $\Sigma_{m}^{3 r d}$, which follows immediately from Lemma 3.12 (see [FS19, Proposition 4.5]):

Lemma 3.13. - Let $u: \mathbf{B}_{1} \rightarrow[0, \infty)$ solve (3.1). Then, for all $x_{\circ}$ in $\Sigma_{m}^{3 r d}$ with $0 \leq m \leq$ $n-1$, the limit

$$
p_{3, x_{0}}:=\lim _{r \downarrow 0} \frac{1}{r^{3}}\left(u\left(x_{\circ}+r \cdot\right)-p_{2, x_{o}}(r \cdot)\right)
$$

exists, and $p_{3, x_{0}}$ is a 3-homogeneous harmonic polynomial vanishing on $\left\{p_{2, x_{0}}=0\right\}$.
When $x_{\circ}=0$ we simplify the notation $p_{3,0}$ to $p_{3}$.

## 4. Higher order blow-ups on the maximal stratum

As explained in Section 1.2, in order to prove the main result of this paper (Theorem 1.1) we need to obtain - among other things - an expansion up to order $\mathrm{O}\left(|x|^{5-\zeta}\right)$ at "most points" of $\Sigma_{n-1}^{3 r d}$. This requires a very detailed analysis of such set, which is the goal of this section. From now on, we will only study the points of $\Sigma_{n-1}$ (hence, $m=n-1$ ).

In order to study the higher regularity properties of the set $\Sigma_{n-1}^{3 r d}$, we need a new frequency function for $u-p_{2}-p_{3}$. The following lemma is a more flexible version of Lemma 3.6. It will be useful later in order to prove the almost monotonicity of $\phi^{\gamma}(r, w)$ for a suitable $\gamma$, where $w$ will be the difference between $u$ and its polynomial expansions at singular points.

Lemma 4.1. - Let $\mathrm{R} \in(0,1)$, and let $w: \mathrm{B}_{\mathrm{R}} \rightarrow[0, \infty)$ be a $\mathrm{C}^{1,1}$ function. Assume that for some $\kappa \in(0,1)$ we have

$$
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-r^{\kappa-1} \quad \forall r \in(0, \mathrm{R}) .
$$

Then the following holds:
(a) Suppose that $0<\underline{\lambda} \leq \phi^{\gamma}(r, w) \leq \bar{\lambda}$ for all $r \in(0, \mathrm{R})$. Then, for any given $\delta>0$ we have

$$
\frac{1}{\mathrm{C}_{\delta}}\left(\frac{\mathrm{R}}{r}\right)^{2 \underline{\lambda}-\delta} \leq \frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} \leq \mathrm{C}_{\delta}\left(\frac{\mathrm{R}}{r}\right)^{2 \bar{\lambda}+\delta} \quad \text { for all } r \in(0, \mathrm{R}) \text {, }
$$

where $\mathrm{C}_{\delta}$ depends only on $n, \gamma, \kappa, \bar{\lambda}, \delta$.
(b) Assume in addition that

$$
\frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq-r^{\kappa} \quad \forall r \in(0, \mathrm{R})
$$

Then, for $\lambda_{*}:=\phi^{\gamma}\left(0^{+}, w\right)$, we have

$$
e^{-\frac{4}{k^{2}}}\left(\frac{\mathrm{R}}{r}\right)^{2 \lambda_{*}} \leq \frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}}
$$

Proof. - (a) Define

$$
\mathrm{F}(r):=\frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}}
$$

so that, by assumption, we have
(4.1)

$$
\frac{d}{d r} \phi^{\gamma}(r, w)+r^{k-1} \geq \frac{2}{r}(\mathrm{~F}(r))^{2}
$$

It follows by Lemma 2.1 that

$$
\begin{align*}
\frac{\frac{d}{d r}\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)} & =\frac{2}{r} \frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w+r^{2-n} \int_{\mathrm{B}_{r}}|\nabla w|^{2}+\gamma r^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}}  \tag{4.2}\\
& =\frac{2}{r} \phi^{\gamma}(r, w)+\frac{2}{r} \mathrm{~F}(r) .
\end{align*}
$$

Integrating (4.1) and using Cauchy-Schwarz inequality, since $0 \leq \phi^{\gamma}(r, w) \leq \bar{\lambda}$ for all $r \in(0, \mathrm{R})$, we get
(4.3)

$$
\begin{aligned}
\left|\int_{r}^{\mathrm{R}} \frac{1}{\rho} \mathrm{~F}(\rho) d \rho\right| & \leq\left(\int_{r}^{\mathrm{R}} \frac{1}{\rho}(\mathrm{~F}(\rho))^{2} d \rho\right)^{1 / 2}\left(\int_{r}^{\mathrm{R}} \frac{d \rho}{\rho}\right)^{1 / 2} \\
& \leq\left(\int_{r}^{\mathrm{R}}\left(\frac{d}{d r} \phi^{\gamma}(\rho, w)+\rho^{\kappa-1}\right) d \rho\right)^{1 / 2}\left(\int_{r}^{\mathrm{R}} \frac{d \rho}{\rho}\right)^{1 / 2} \\
& \leq\left(\bar{\lambda}+\frac{1}{\kappa}\left(\mathrm{R}^{\kappa}-r^{\kappa}\right)\right)^{1 / 2}(\log (\mathrm{R} / r))^{1 / 2} \\
& \leq \mathrm{C}(\log (\mathrm{R} / r))^{1 / 2} .
\end{aligned}
$$

Hence, integrating (4.2) between $r$ and R (recall $0<r<\mathrm{R}<1$ ) and using $\phi^{\gamma}(\rho, w) \leq \bar{\lambda}$ and (4.3) we obtain

$$
\begin{aligned}
& \log \frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} \\
& \quad \leq \int_{r}^{\mathrm{R}} \frac{2}{\rho}(\bar{\lambda}+\mathrm{F}(\rho)) d \rho \leq 2 \bar{\lambda} \log (\mathrm{R} / r)+\mathrm{C}(\log (\mathrm{R} / r))^{1 / 2} \\
& \quad \leq(2 \bar{\lambda}+\delta) \log (\mathrm{R} / r)+\mathrm{C}_{\delta} .
\end{aligned}
$$

Similarly (now using $\phi^{\gamma}(\rho, w) \geq \underline{\lambda}$ ) we obtain

$$
\log \frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq(2 \underline{\lambda}-\delta) \log (\mathrm{R} / r)-\mathrm{C}_{\delta} .
$$

(b) In this case we have $\mathrm{F}(r) \geq-r^{\kappa}$ for all $r \in(0, \mathrm{R})$. Hence, integrating (4.2) between 0 and $\rho \in(0, R)$ we obtain

$$
\phi^{\gamma}(\rho, w)-\lambda_{*} \geq-\frac{1}{\kappa} \rho^{\kappa} .
$$

Thus,

$$
\begin{aligned}
\log \frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} & \geq \int_{r}^{\mathrm{R}} \frac{2}{\rho}\left(\phi^{\gamma}(\rho, w)+\mathrm{F}(\rho)\right) d \rho \\
& \geq 2 \lambda_{*} \int_{r}^{\mathrm{R}} \frac{d \rho}{\rho}-\left(\frac{2}{\kappa}+1\right) \int_{r}^{\mathrm{R}} \rho^{\kappa-1} d \rho \\
& \geq 2 \lambda_{*} \log (\mathrm{R} / r)-\frac{4}{\kappa^{2}}
\end{aligned}
$$

Remark 4.2. - An interesting consequence of Lemma 4.1(a) is the following. If $w$ is as in of Lemma 4.1(a) then $\phi^{\gamma}\left(0^{+}, w\right) \leq \gamma$. Indeed, otherwise we would have $\phi^{\gamma}(r, w) \geq \underline{\lambda}:=\gamma+\delta>\gamma$ for all $r>0$ small, and Lemma 4.1(a) would imply that $\mathrm{H}(r, w)+r^{2 \gamma} \leq \mathrm{C} r^{2 \lambda}-\delta=\mathrm{C}^{2 \gamma+\delta}$ for $r \ll 1$, impossible.

The following lemma gives the (approximate) monotonicity of $\phi^{\gamma}$ when applied to $w:=u-p_{2}-\mathrm{P}$, where P is any 3-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$.

Lemma 4.3. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1), with $0 \in \Sigma_{n-1}^{\geq 3}$. Let $w:=$ $u-p_{2}-\mathrm{P}$, where P is a 3-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$. Then, given $\gamma \in(3,4)$, for all $r \in(0,1 / 2)$ we have

$$
\begin{aligned}
& \frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-\mathrm{C} r^{3-\gamma} \quad \text { and } \\
& \frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq-\mathrm{C} r^{4-\gamma}
\end{aligned}
$$

where C depends only on $n, \gamma,\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$.
In particular, assuming $0 \in \Sigma_{n-1}^{3 r d}$, the previous inequalities hold for $w:=u-p_{2}-p_{3}$.

Proof. - The proof will rely on a iteration argument where one enlarge the value of $\gamma$, starting from 3 and increasing it up to the desired $\gamma \in(3,4)$. We split the proof in two steps.

- Step 1. We first show that

$$
\begin{equation*}
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-\mathrm{C} r^{3-\gamma} \phi^{\gamma}(r, w) g^{\gamma}(r) \tag{4.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq-\mathrm{C}^{4-\gamma} g^{\gamma}(r) \tag{4.5}
\end{equation*}
$$

where

$$
g^{\gamma}(r):=\frac{\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{2} \backslash \mathrm{~B}_{1}\right)}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{1 / 2}}
$$

and C depends only on $n$ and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\boldsymbol{B}_{1}\right)}$.
Indeed, by Lemma 2.3 we have

$$
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}+\overline{\mathrm{E}}^{\gamma}(r, w),
$$

where

$$
\begin{aligned}
& \overline{\mathrm{E}}^{\gamma}(r, w):=\frac{2}{r} \frac{r^{2-n} \int_{\mathrm{B}_{r}}\left(\lambda_{r} w-x \cdot \nabla w\right) \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}} \text { and } \\
& \lambda_{r}=\phi^{\gamma}(r, w)=\frac{\mathrm{D}(r, w)+\gamma r^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} .
\end{aligned}
$$

Note that $\Delta w=\Delta\left(u-p_{2}-p_{3}\right)=\chi_{\{u>0\}}-1-0=-\chi_{\{u=0\}}$. Also, since $\lambda_{r} \geq 3$ by (3.18), using the inequality $p_{2}+\mathrm{P} \geq-\frac{\mathrm{P}^{2}}{2 p_{2}}$, since $\frac{\mathrm{p}^{2}}{p_{2}}$ is homogeneous of degree 4 we have

$$
\begin{aligned}
\left(\lambda_{r}-x \cdot \nabla\right)\left(p_{2}+\mathrm{P}\right) & \geq\left(\lambda_{r}-2\right) p_{2}+\left(\lambda_{r}-3\right) \mathrm{P} \geq\left(\lambda_{r}-3\right)\left(p_{2}+\mathrm{P}\right) \\
& \geq-\frac{\left(\lambda_{r}-3\right)}{2}\left(\sup _{\partial \mathrm{B}_{1}} \frac{\mathrm{P}^{2}}{p_{2}}\right)|x|^{4} .
\end{aligned}
$$

Therefore we obtain
(4.6)

$$
\begin{aligned}
\overline{\mathrm{E}}^{\gamma}(r, w) & =\frac{2}{r} \frac{r^{2-n} \int_{\mathrm{B}_{r} \cap\{u=0\}}\left(\lambda_{r}-x \cdot \nabla\right)\left(p_{2}+\mathrm{P}\right)}{\mathrm{H}(r, w)+r^{2 \gamma}} \\
& \geq-\left(\lambda_{r}-3\right) r^{3-\gamma}\left(\sup _{\partial \mathrm{B}_{1}} \frac{\mathrm{P}^{2}}{p_{2}}\right) \frac{r^{2-n}\left|\mathrm{~B}_{r} \cap\{u=0\}\right|}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{1 / 2}} .
\end{aligned}
$$

Also, since $\Delta w=-\chi_{\{u=0\}} \leq 0$, choosing $\chi \in \mathrm{C}_{c}^{\infty}\left(\mathrm{B}_{2}\right)$ a nonnegative cut-off satisfying $\chi=1$ in $\mathrm{B}_{1}$, integrating by parts we obtain

$$
\begin{aligned}
r^{2-n}\left|\{u=0\} \cap \mathrm{B}_{r}\right| & =\int_{\mathrm{B}_{1}}-\Delta w_{r} \leq \int_{\mathrm{B}_{2}}-\Delta w_{r} \chi \\
& =-\int_{\mathrm{B}_{2}} w_{r} \Delta \chi \leq \mathrm{C} \int_{\mathrm{B}_{2} \backslash \mathrm{~B}_{1}}\left|w_{r}\right| \leq \mathrm{C}\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{2} \backslash \mathrm{~B}_{1}\right)} .
\end{aligned}
$$

Thus, since $\lambda_{r}=\phi^{\gamma}(r, w)$ and recalling (4.6), we have shown that

$$
\overline{\mathrm{E}}^{\gamma}(r, w) \geq-\mathrm{C} r^{3-\gamma} \phi^{\gamma}(r, w) g^{\gamma}(r)
$$

and (4.4) follows. Note that, since by assumption $\mathrm{P}^{2}$ is divisible by $p_{2}$, we have that $\sup _{\partial \mathrm{B}_{1}} \frac{\mathrm{P}^{2}}{p_{2}}$ is bounded by a constant depending only on $n$ and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathbf{B}_{1}\right)}$, and thus the constant C above depends only on $n$ and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$.

Similarly, using again $p_{2}+\mathrm{P} \geq-\frac{\mathrm{P}^{2}}{2 p_{2}}$, we obtain

$$
\frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}}=\frac{r^{2-n} \int_{\{u=0\} \cap \mathrm{B}_{r}}\left(p_{2}+\mathrm{P}\right)}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq-\mathrm{C}^{4-\gamma} \frac{r^{2-n}\left|\{u=0\} \cap \mathrm{B}_{r}\right|}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{1 / 2}},
$$

which gives (4.5).

- Step 2. Next we show that (4.4) implies that, for all $\gamma<4$,

$$
\begin{equation*}
\phi^{\gamma}(r, w) \leq \mathrm{C}_{\gamma} \quad \text { and } \quad g^{\gamma}(r) \leq \mathrm{C}_{\gamma}, \tag{4.7}
\end{equation*}
$$

with $\mathrm{C}_{\gamma}$ depending only on $n, \gamma$, and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$.
We prove (4.7) for all $\gamma \in[3,4)$ by iteratively increasing the value of $\gamma$ at each iteration, starting from $\gamma=3$, in order to always have (along the iteration) a uniform bound on $\phi^{\nu}(r, w)$ and $g^{\gamma}(r)$.

First, we observe that since $0 \in \Sigma_{n-1}^{\geq 3}$ we have $\phi\left(0^{+}, u-p_{2}\right) \geq 3$, hence $\left|u-p_{2}\right| \leq$ $\mathrm{C}|x|^{3}$. Therefore $\left|u-p_{2}-\mathrm{P}\right| \leq \mathrm{C}|x|^{3}$, which immediately implies that $g^{3}(r) \leq \mathrm{C}_{3}$, and then it follows by $(4.4)$ that $\log \left(\phi^{3}(\cdot, w)\right)$ is almost monotonically increasing, so in particular it is uniformly bounded.

We show next that if (4.7) holds for some $\gamma \geq 3$ then (4.7) holds also with $\gamma$ replaced by $\gamma+\beta$ for any $\beta>0$ such that $3 \beta<4-\gamma$.

Indeed, we can bound

$$
\begin{aligned}
\phi^{\gamma+\beta}(r, w) & =\frac{\mathrm{D}(r, w)+(\gamma+\beta) r^{2 \gamma+2 \beta}}{\mathrm{H}(r, w)+r^{2 \gamma+2 \beta}} \leq \frac{2}{r^{2 \beta}} \frac{\mathrm{D}(r, w)+\gamma r^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} \leq \frac{\phi^{\gamma}(r, w)}{r^{2 \beta}} \\
& \leq 2 \frac{\mathrm{C}_{\gamma}}{r^{2 \beta}},
\end{aligned}
$$

and similarly

$$
g^{\gamma+\beta}(r)=\frac{\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{2} \backslash \mathrm{~B}_{1}\right)}}{\left(\mathrm{H}(r, w)+r^{2 \gamma+2 \beta}\right)^{1 / 2}} \leq \frac{1}{r^{\beta}} \frac{\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{2} \backslash \mathrm{~B}_{1}\right)}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{1 / 2}}=\frac{g^{\gamma}(r)}{r^{\beta}} \leq \frac{\mathrm{C}_{\gamma}}{r^{\beta}} .
$$

Then (4.4)-with $\gamma$ replaced by $\gamma+\beta$-yields

$$
\begin{equation*}
\frac{d}{d r} \phi^{\gamma+\beta}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-\mathrm{CC}_{\gamma}^{2} r^{4-\gamma-1-3 \beta} \tag{4.8}
\end{equation*}
$$

Noticing that $r^{4-\gamma-1-3 \beta}$ is integrable over $r \in(0,1)$ provided $3 \beta<4-\gamma$, (4.8) implies that $\phi^{\gamma+\beta}(r, w)$ is almost monotonically increasing. In particular $\phi^{\gamma+\beta}(r, w) \leq \mathrm{C}^{\prime}$ for $r \in(0,1 / 2)$, where $\mathrm{C}^{\prime}$ is a constant depending only on $n$ and $\gamma+\beta$.

In addition, (4.8) and Lemma 4.1(a) imply that

$$
\frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma+2 \beta}}{\mathrm{H}(r, w)+r^{2 \gamma+2 \beta}} \leq \mathrm{C}^{\prime} \quad \forall \mathrm{R} \in(r, 2 r),
$$

thus

$$
\left(g^{\gamma+\beta}(r)\right)^{2}=\frac{\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{2} \backslash \mathrm{~B}_{1}\right)}^{2}}{\mathrm{H}(r, w)+r^{2 \gamma+2 \beta}} \leq \mathrm{C} \int_{r}^{2 r} \frac{\mathrm{H}(\mathrm{R}, w) d \mathrm{R}}{\mathrm{H}(r, w)+r^{2 \gamma+2 \beta}} \leq \mathrm{CC}^{\prime},
$$

and therefore (4.7) holds for $\gamma$ replaced by $\gamma+\beta$.
Having proven that if (4.7) holds for some $\gamma \geq 3$ then (4.7) holds also with $\gamma$ replaced by $\gamma+\beta$ for any $\beta>0$ such that $3 \beta<4-\gamma$, iterating finitely many times we conclude that (4.7) holds for any $\gamma<4$, as claimed.

Combining this information with (4.4) we obtain

$$
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-\mathrm{CC}_{\gamma}^{2} r^{3-\gamma} \quad \forall r \in(0,1 / 2)
$$

Also, it follows by (4.5) and (4.7) that $\frac{r^{2-n} \int_{\mathrm{B}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq-\mathrm{CC}_{\gamma} r^{4-\gamma}$. Finally, when $0 \in \Sigma_{n-1}^{3 r d}$ we can take $\mathrm{P}=p_{3}$ (since, by definition of $\Sigma_{n-1}^{3 r d}, p_{3}$ vanishes on $\left\{p_{2}=0\right\}$ ).

We have proved that in $\Sigma_{n-1}^{3 r d}$ we have almost monotonicity of the new truncated frequency function $\phi^{\gamma}\left(r, u-p_{2}-p_{3}\right)$ for any $\gamma \in(3,4)$. This means that $\phi^{\gamma}\left(0^{+}, u-p_{2}-\right.$ $p_{3}$ ) exists, and satisfies $3 \leq \phi^{\gamma}\left(0^{+}, u-p_{2}-p_{3}\right) \leq \gamma$ (see Remark 4.2).

Hence, we can now introduce the following:
Definition 4.4. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1). We denote by $\Sigma_{n-1}^{>3}$ the set of points $x_{\circ} \in \Sigma_{n-1}^{3 r d}$ such that, for $w:=u\left(x_{0}+\cdot\right)-p_{2, x_{0}}-p_{3, x_{0}}$, we have $\phi^{\gamma}\left(0^{+}, w\right)>3$ for any $\gamma \in(3,4)$.

Moreover, we will need the following:

Definition 4.5. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and assume that $0 \in \Sigma_{n-1}^{3 r d}$. Choose a coordinate system such that

$$
\begin{equation*}
p_{2}(x)=\frac{1}{2} x_{n}^{2} \quad \text { and } \quad p_{3}(x)=\sum_{\alpha=1}^{n-1} \frac{a_{\alpha}}{2} x_{\alpha}^{2} x_{n}+\frac{a_{n}}{6} x_{n}^{3} . \tag{4.9}
\end{equation*}
$$

(Here we used that $p_{3}$ is harmonic and vanishes on $\left\{x_{n}=0\right\}$.) We define the fourth order polynomial Ansatz at 0 , and we denote it by $\mathscr{P}$, as ${ }^{3}$
(4.10)

$$
\mathscr{P}(x):=\frac{1}{2} x_{n}^{2}+p_{3}+\frac{1}{2}\left(\frac{p_{3}}{x_{n}}\right)^{2}+x_{n} \mathrm{Q}=\frac{1}{2}\left(x_{n}+\frac{p_{3}}{x_{n}}+\mathrm{Q}\right)^{2}+\mathrm{O}\left(|x|^{5}\right) .
$$

Here Q is a 3-homogeneous polynomial which depends only on $p_{2}$ and $p_{3}$, and is defined as follows

$$
\begin{equation*}
\mathrm{Q}(x):=\sum_{\alpha=1}^{n-1}\left(a_{\alpha}^{2}-\frac{a_{\alpha} a_{n}}{3}\right)\left(\frac{x_{n}^{3}}{12}-\frac{x_{\alpha}^{2} x_{n}}{2}\right) . \tag{4.11}
\end{equation*}
$$

When $u$ is a solution of (3.1) and $x_{0} \in \Sigma_{n-1}^{3 r d}$, we define $\mathscr{P}_{x_{0}}$ to be the fourth oder polynomial Ansatz at 0 of $u\left(x_{0}+\cdot\right)$ (note that $\mathscr{P}_{x_{0}}$ depends only on $p_{2, x_{0}}$ and $p_{3, x_{0}}$ ). In addition, for $\alpha \in\{1,2, \ldots, n-1\}$ we define the osculating rotation vector fields at 0 as

$$
\begin{equation*}
\mathbf{X}_{\alpha}:=\left(1+a_{\alpha} x_{n}\right) \boldsymbol{e}_{\alpha}-a_{\alpha} x_{\alpha} \boldsymbol{e}_{n}, \quad \text { where } \boldsymbol{e}_{i}=(0, \ldots, 0, \stackrel{i}{1}, 0, \ldots, 0), \tag{4.12}
\end{equation*}
$$

where $a_{\alpha}$ are as in (4.9).
We will use the following notation throughout the section. Given $f \in \mathrm{C}^{1}\left(\mathbf{R}^{n}\right)$, we denote by $\mathbf{X}_{\alpha} f$ the derivative of $f$ in the direction of $\mathbf{X}_{\alpha}$, namely $\mathbf{X}_{\alpha} f=\left(1+a_{\alpha} x_{n}\right) \partial_{\alpha} f-$ $a_{\alpha} x_{\alpha} \partial_{n} f$.

Lemma 4.6. - Given $p_{2}$ and $p_{3}$ as in (4.9), define $\mathbf{Q}$ as (4.11). Then, $\mathscr{P}$ given by (4.10) satisfies

$$
\Delta \mathscr{P}=1 \quad \text { and } \quad \mathbf{X}_{\alpha} \mathbf{X}_{\alpha} \mathscr{P}=\mathrm{O}\left(|x|^{3}\right) \quad \forall \alpha \in\{1,2, \ldots, n-1\} .
$$

Proof. - Let $p_{2}, p_{3}, \mathscr{P}$, and $\mathbf{X}_{\alpha}$, be as in (4.9), (4.10), (4.11), (4.12). We compute

$$
\begin{aligned}
& \mathbf{X}_{\alpha} \mathbf{X}_{\alpha}\left(\frac{x_{n}^{2}}{2}\right)=-a_{\alpha} x_{n}+a_{\alpha}^{2}\left(x_{\alpha}^{2}-x_{n}^{2}\right), \\
& \mathbf{X}_{\alpha} \mathbf{X}_{\alpha} p_{3}=a_{\alpha} x_{n}+2 a_{\alpha}^{2}\left(x_{n}^{2}-x_{\alpha}^{2}\right)-\sum_{\beta=1}^{n-1} \frac{a_{\alpha} a_{\beta}}{2} x_{\beta}^{2}-\frac{a_{\alpha} a_{n}}{2} x_{n}^{2}+\mathrm{O}\left(|x|^{3}\right),
\end{aligned}
$$

[^3]\[

$$
\begin{aligned}
& \mathbf{X}_{\alpha} \mathbf{X}_{\alpha}\left(\frac{1}{2}\left(\frac{p_{3}}{x_{n}}\right)^{2}\right)=\sum_{\alpha=1}^{n-1} \frac{a_{\alpha} a_{\beta}}{2} x_{\beta}^{2}+\frac{a_{\alpha} a_{n}}{6} x_{n}^{2}+a_{\alpha}^{2} x_{\alpha}^{2}+\mathrm{O}\left(|x|^{3}\right) \\
& \mathbf{X}_{\alpha} \mathbf{X}_{\alpha}\left(x_{n} \mathrm{O}\right)=-\left(a_{\alpha}^{2}-\frac{a_{\alpha} a_{n}}{3}\right) x_{n}^{2}+\mathrm{O}\left(|x|^{3}\right)
\end{aligned}
$$
\]

and thus adding them we obtain

$$
\mathbf{X}_{\alpha} \mathbf{X}_{\alpha} \mathscr{P}=\mathrm{O}\left(|x|^{3}\right) .
$$

Similarly, using that $\sum_{\alpha=1}^{n-1} a_{\alpha}=-a_{n}$ (as a consequence of the fact that $p_{3}$ is harmonic), a direct (but tedious) computation shows that

$$
\Delta\left(\frac{1}{2}\left(\frac{p_{3}}{x_{n}}\right)^{2}+x_{n} \mathrm{Q}\right)=0
$$

therefore $\Delta \mathscr{P} \equiv 1$.
We will need the following semiconvexity estimate in the spirit of Lemma 3.8.
Lemma 4.7. - Assume that $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ is a solution of (3.1) and that $0 \in \Sigma_{n-1}^{3 r d}$. Let $w:=u-\mathscr{P}-\mathrm{P}$, where P is a 4-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$. Then

$$
\inf _{\mathrm{B}_{r}} r^{2} \mathbf{X}_{\alpha} \mathbf{X}_{\alpha} w \geq-\mathrm{C}(\mathrm{P})\left(\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)}+r^{5}\right)
$$

and

$$
\sup _{\mathbf{B}_{r / 2}} r|\nabla w| \leq \mathrm{C}(\mathrm{P})\left(\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)}+r^{5}\right)
$$

for all $r \in(0,1 / 5)$.
Proof. - Let $p_{2}, p_{3}, \mathscr{P}$, and $\mathbf{X}_{\alpha}$, be as in (4.9), (4.10), (4.11), (4.12), and fix $\alpha \in$ $\{1,2, \ldots, n-1\}$.

- Step 1. For $r>0$ small, we consider the rescaled vector field $\mathbf{X}_{\alpha}^{r}=\mathbf{X}_{\alpha}(r \cdot)$, and denote $w_{r}=w(r \cdot)$ and $v:=\mathbf{X}_{\alpha}^{r} \mathbf{X}_{\alpha}^{r} w_{r}$. We consider

$$
\bar{v}(x):=\min \left\{v(x),-\mathrm{C}(\mathrm{P}) r^{5}\right\}
$$

for some constant $\mathrm{C}(\mathrm{P})>0$ depending on P , to be chosen. We claim that

$$
\Delta \bar{v} \leq 0 \quad \text { in } \mathrm{B}_{5} .
$$

Since $w_{r}$ is $\mathrm{C}^{1,1}$ (for fixed $r>0$ ) but not $\mathrm{C}^{2}$, to prove that $\bar{v}$ is subharmonic we need to proceed similarly to the proof of Lemma 3.8, now taking second order "rotational"
incremental quotients. More precisely, let $\phi_{\mathbf{X}_{\alpha}^{r}}^{h}$ denote the integral flow of the vector field $\mathbf{X}_{\alpha}^{r}$ at time $h$, and define

$$
v^{h}:=\frac{w_{r} \circ \phi_{\mathbf{X}_{\alpha}^{r}}^{h}+w_{r} \circ \phi_{\mathbf{X}_{\alpha}}^{-h}-2 w_{r}}{h^{2}}
$$

On the one hand, since $\phi_{\mathbf{X}_{\alpha}^{\prime}}^{h}$ is a rotation (and thus it commutes with $\Delta$ ), noticing that $\Delta w=\chi_{\{u>0\}}-1 \leq 0$ in $\mathrm{B}_{1}$ and $\Delta w=0$ in $\{u>0\}$ we obtain

$$
\begin{equation*}
\Delta v^{h} \leq 0 \quad \text { in }\{u(r \cdot)>0\} \cap \mathrm{B}_{1 / r} \tag{4.13}
\end{equation*}
$$

On the other hand, we claim that
(4.14)

$$
v^{h} \geq-\mathrm{C}(\mathrm{P}) r^{5} \quad \text { in }\{u(r \cdot)=0\} \cap \mathrm{B}_{5} .
$$

Indeed, recalling that $\mathbf{X}_{\alpha} \mathbf{X}_{\alpha} \mathscr{P}=\mathrm{O}\left(|x|^{3}\right)$ (by Lemma 4.6) and since $\mathbf{X}_{\alpha} \mathbf{X}_{\alpha} \mathrm{P}=\partial_{\alpha \alpha} \mathrm{P}+$ $\mathrm{O}\left(|x|^{3}\right)$, we obtain

$$
\mathbf{X}_{\alpha} \mathbf{X}_{\alpha}(-\mathscr{P}-\mathrm{P}) \geq-\partial_{\alpha \alpha} \mathrm{P}+\mathrm{O}\left(|x|^{3}\right)
$$

In addition, since P is 4-homogeneous and vanishes on $\left\{x_{n}=0\right\}$, we wave $\partial_{\alpha \alpha} \mathrm{P}=x_{n} \ell\left(x^{\prime}\right)$ where $\ell$ is some linear function, thus

$$
\left|\partial_{\alpha \alpha} \mathrm{P}\right| \leq \mathrm{C}(\mathrm{P})|x|\left|x_{n}\right| .
$$

Therefore, combining all these estimates, we get

$$
\begin{equation*}
\left|\mathbf{X}_{\alpha} \mathbf{X}_{\alpha}(\mathscr{P}+\mathrm{P})\right| \geq \mathrm{C}(\mathrm{P})\left(r^{3}+r\left|x_{n}\right|\right) \quad \text { in } \mathrm{B}_{5 r} . \tag{4.15}
\end{equation*}
$$

In addition, thanks to Proposition 3.10 (recall that $\lambda^{2 n d} \geq 3$, since by assumption $0 \in \Sigma_{n-1}^{3 r d}$ ) we have

$$
\begin{equation*}
\{u=0\} \cap \mathrm{B}_{1 / 2} \subset\left\{\left|x_{n}\right| \leq \mathrm{C}\left|x^{\prime}\right|^{2}\right\} \tag{4.16}
\end{equation*}
$$

thus

$$
\begin{aligned}
\left(x^{\prime}, x_{n}\right) \in\{u=0\} \cap \mathrm{B}_{5 r} & \Rightarrow\left|x^{\prime}\right| \leq 5 r,\left|x_{n}\right| \leq \mathrm{C} r^{2} \\
& \Rightarrow\left|\phi_{\mathbf{x}_{\alpha}}^{\text {}}(x) \cdot \boldsymbol{e}_{n}\right| \leq \mathrm{C} r^{2} \quad \forall h \in(0,2) .
\end{aligned}
$$

As a consequence, combining this bound with (4.15) we obtain

$$
\begin{aligned}
& \frac{\left|(\mathscr{P}+\mathrm{P}) \circ \phi_{\mathbf{x}_{\alpha}}^{r h}+(\mathscr{P}+\mathrm{P}) \circ \phi_{\mathbf{x}_{\alpha}}^{-r h}-2(\mathscr{P}+\mathrm{P})\right|}{(r h)^{2}} \leq \mathrm{C}(\mathrm{P}) r^{3} \\
& \quad \text { in }\{u=0\} \cap \mathrm{B}_{5 r} \quad \forall h \in(0,1) .
\end{aligned}
$$

Rescaling this estimate one gets (4.14), that combined with (4.13) implies that

$$
\bar{v}^{h}:=\min \left\{v^{h}(x),-\mathrm{C}(\mathrm{P}) r^{5}\right\} \quad \text { is superharmonic. }
$$

Therefore, since $\bar{v}=\lim _{h \downarrow 0} \bar{v}^{h}$ a.e. then the function $\bar{v}$ is superharmonic too, as claimed.

- Step 2. Note that if we consider $\mathrm{V}:=\partial_{\alpha \alpha} w_{r}$ instead of rotational derivatives (as we did in the proof of Lemma 3.8), then the same argument as the one above gives

$$
\left|\partial_{\alpha \alpha}(\mathscr{P}+\mathrm{P})\right| \leq \mathrm{C}(\mathrm{P}) r^{2} \quad \text { in } \mathrm{B}_{5 r}
$$

(cf. (4.15)), from which one deduces that the function $\overline{\mathrm{V}}:=\min \left\{\mathrm{V},-\mathrm{C}(\mathrm{P}) r^{4}\right\}$ is superharmonic (notice the difference in the power of $r$ in the definitions of $\bar{v}$ and $\overline{\mathrm{V}}$ ).

- Step 3. Now, as in the proof of Lemma 3.8, by weak Harnack inequality, interpolation, and the Calderón-Zygmund theory, we have

$$
\begin{aligned}
\|\bar{v}\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{1}\right)} & \leq \mathrm{C}(n)\left(f_{\mathrm{B}_{3 / 2}}|\bar{v}|^{\varepsilon}\right)^{1 / \varepsilon} \leq \mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{W}_{\text {watk }}^{2,1}\left(\mathrm{~B}_{2}\right)}+r^{5}\right) \\
& \leq \mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{3} \backslash \mathrm{~B}_{2}\right)}+\left\|\Delta w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{3}\right)}+r^{5}\right) .
\end{aligned}
$$

On the other hand, since $\Delta w_{r} \leq 0$ (because $\left.\Delta w=\Delta u-\Delta(\mathscr{P}-\mathrm{P})=\Delta u-1 \leq 0\right)$, reasoning as in the proof of Lemma 3.8 (cf. (3.10)) we obtain

$$
\left\|\Delta w_{r}\right\|_{\mathrm{L}^{1}\left(\mathbb{B}_{3}\right)} \leq \mathrm{C}(n)\left\|w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4} \backslash \mathrm{~B}_{3}\right)} .
$$

Hence $\|\bar{v}\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{1}\right)} \leq \mathrm{C}\left\|w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4} \backslash \mathrm{~B}_{3}\right)}$, which yields

$$
\begin{equation*}
\inf _{\mathrm{B}_{1}} \mathbf{X}_{\alpha} \mathbf{X}_{\alpha} w_{r} \geq-\mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4} \backslash \mathrm{~B}_{3}\right)}+r^{5}\right) \tag{4.17}
\end{equation*}
$$

and the first part of the lemma (semiconvexity) follows easily by scaling.
In order to prove the Lipschitz bound, we note that if we repeat the same reasoning with $\overline{\mathrm{V}}$ instead of $\bar{v}$, we find instead the semiconvexity estimates

$$
\inf _{\mathrm{B}_{1}} \partial_{\alpha \alpha} w_{r} \geq-\mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4} \backslash \mathrm{~B}_{3}\right)}+r^{4}\right), \quad \text { for } 1 \leq \alpha \leq n-1
$$

Although this estimate is less precise (it has an error of size $r^{4}$ instead of $r^{5}$ ) it is still useful. Indeed, using that $\Delta w_{r} \leq 0$, it implies the semiconcavity estimate

$$
\sup _{\mathrm{B}_{1}} \partial_{n n} w_{r} \leq \mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4} \backslash \mathrm{~B}_{3}\right)}+r^{4}\right) .
$$

Combined together, these semiconvexity and semiconcavity estimates imply a bound on the Lipschitz constant of $w_{r}$ in terms of its $\mathrm{L}^{\infty}$ norm in $\mathrm{B}_{4} \backslash \mathrm{~B}_{3}$ and its semiconvexity/semiconcavity constants, that is

$$
\begin{equation*}
\left\|w_{r}\right\|_{\operatorname{Lip}\left(\mathrm{B}_{3 / 4}\right)} \leq \mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{4} \backslash \mathrm{~B}_{3}\right)}+r^{4}\right) . \tag{4.18}
\end{equation*}
$$

Although this is not the desired bound, this will be useful in the next step to obtain the sharp bound.

- Step 4. To conclude the proof, we need to improve (4.18) and get Lipschitz bound for $w_{r}$ in $\mathrm{B}_{1 / 2}$ with an error $\mathrm{O}\left(r^{5}\right)$. For this, we note that for each $\alpha=1, \ldots, n-1$ the unit vector field $\mathrm{E}_{\alpha}=\mathbf{X}_{\alpha}^{r} /\left|\mathbf{X}_{\alpha}^{r}\right|$ satisfies $\left|\mathrm{E}_{\alpha}-\boldsymbol{e}_{\alpha}\right| \leq \mathrm{C} r$ in $\mathrm{B}_{1}$, and thus we can complete $\left\{\mathrm{E}_{\alpha}\right\}_{\alpha=1}^{n-1}$ to obtain a orthonormal moving frame by adding a vectorfield $\mathrm{E}_{n}$ satisfying $\left|\mathrm{E}_{n}-\boldsymbol{e}_{n}\right| \leq \mathrm{C} r$ in $\mathrm{B}_{1}$.

Note that, since $\nabla_{\mathbf{X}_{\alpha}^{r}} \mathbf{X}_{\alpha}^{r}=\mathrm{O}(r)$ and $\nabla_{\mathrm{E}_{\alpha}} \mathrm{E}_{\alpha}=\mathrm{O}(r)$, we can choose $\mathrm{E}_{n}$ satisfying $\nabla_{\mathrm{E}_{n}} \mathrm{E}_{n}=\mathrm{O}(r)$. Hence, using these bounds and $\Delta w_{r} \leq 0$, we obtain

$$
\begin{aligned}
\mathrm{E}_{n} \mathrm{E}_{n} w_{r}-\mathrm{C} r\left\|w_{r}\right\|_{\mathrm{Lip}\left(\mathrm{~B}_{3 / 4}\right)} & \leq \mathrm{D}^{2} w_{r}\left(\mathrm{E}_{n}, \mathrm{E}_{n}\right) \leq-\sum_{\alpha=1}^{n-1} \mathrm{D}^{2} w_{r}\left(\mathrm{E}_{\alpha}, \mathrm{E}_{\alpha}\right) \\
& \leq-\sum_{\alpha=1}^{n-1} \mathbf{X}_{\alpha}^{r} \mathbf{X}_{\alpha}^{r} w_{r}+\mathrm{C} r\left\|w_{r}\right\|_{\mathrm{Lip}\left(\mathrm{~B}_{3 / 4}\right)} \quad \text { in } \mathrm{B}_{3 / 4} .
\end{aligned}
$$

Thus, recalling (4.17) and (4.18), we get

$$
\begin{equation*}
\sup _{\mathrm{B}_{3 / 4}} \mathrm{E}_{n} \mathrm{E}_{n} w_{r} \leq \mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{4} \backslash \mathrm{~B}_{3}\right)}+r\left\|w_{r}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{4} \backslash \mathrm{~B}_{3}\right)}+r^{5}\right) \tag{4.19}
\end{equation*}
$$

It follows by (4.17) (resp. (4.19)) that the restriction of $w_{r}$ to the integral curves of $\mathbf{X}_{\alpha}$ (resp. $\mathrm{E}_{n}$ ) is semiconvex (resp. semiconcave), and hence Lipschitz along these curves. Since the directions of these curves span $\mathbf{R}^{n}$, this yields

$$
\sup _{\mathrm{B}_{1 / 2}}\left|\nabla w_{r}\right| \leq \operatorname{Cisup}_{\mathrm{B}_{1 / 2}}\left(\sum_{\alpha=1}^{n-1}\left|\nabla w_{r} \cdot \mathbf{X}_{\alpha}\right|+\left|\nabla w_{r} \cdot \mathrm{E}_{n}\right|\right) \leq \mathrm{C}\left(\left\|w_{r}\right\|_{L^{\infty}\left(\mathrm{B}_{4} \backslash \mathrm{~B}_{3}\right)}+r^{5}\right) .
$$

Finally, to conclude the proof, it suffices to show that the $L^{\infty}\left(B_{4} \backslash B_{3}\right)$-norm above can be replaced by $\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{2}\right)}+\mathrm{C}(\mathrm{P}) r^{5}$. Indeed, recalling that $\Delta w=-\chi_{\{u=0\}}$, the function $w$ is superharmonic everywhere, and harmonic outside $\{u=0\}$. In particular, this gives the desired control on the $\mathrm{L}^{\infty}$ norm of $w_{-}$. In addition, since $\mathscr{P} \geq \mathrm{O}\left(|x|^{5}\right)$ and P is a 4-th order polynomial vanishing on $\left\{x_{n}=0\right\}$, it follows from (4.16) that

$$
w=u-\mathscr{P}-\mathrm{P} \leq \mathrm{P} \leq \mathrm{C}(\mathrm{P}) r^{5} \quad \text { inside }\{u=0\} \cap \mathrm{B}_{r},
$$

hence the function $\max \left\{w_{+}, \mathrm{C}(\mathrm{P}) r^{5}\right\}$ is subharmonic. Thus, the mean value inequalities allow us to control the $\mathrm{L}^{\infty}$ norm of $\left(w_{r}\right)_{ \pm}$with the $\mathrm{L}^{1}$ (or $\mathrm{L}^{2}$ ) norm of $\left|w_{r}\right|+\mathrm{Cr} r^{5}$ in a slightly larger domain, concluding the proof.

Remark 4.8. - We note that Lemma 2.3 can be rewritten as

$$
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}+\frac{2}{r} \int_{\mathrm{B}_{1}}\left(\lambda_{r} \hat{w}_{r}^{(\gamma)}-x \cdot \nabla \hat{w}_{r}^{(\gamma)}\right) \Delta \hat{w}_{r}^{(\gamma)},
$$

where

$$
\lambda_{r}:=\phi^{\gamma}(r, w) \quad \text { and } \quad \hat{w}_{r}^{(\gamma)}:=\frac{w(r \cdot)}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{1 / 2}} .
$$

Also, we observe that Lemma 4.7 yields $\left\|\nabla \hat{w}_{r}^{(\gamma)}\right\|_{L^{\infty}} \leq \mathrm{C}$ for all $\gamma \leq 5$ when $w:=u-$ $\mathscr{P}-\mathrm{P}$. This will be crucial in the proof of Lemma 4.9 below.

Lemma 4.9. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ be a solution of (3.1), and assume that $0 \in \Sigma_{n-1}^{3 r d}$. Set $w:=u-\mathscr{P}-\mathrm{P}$, where $\mathscr{P}$ is defined in (4.10), and P is a 4-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$. Then, given $\gamma \in(4,5)$, for all $r \in(0,1 / 2)$ we have

$$
\begin{aligned}
& \frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-\mathrm{C} r^{4-\gamma} \quad \text { and } \\
& \frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq-\mathrm{Cr} r^{5-\gamma}
\end{aligned}
$$

where C is a constant depending only on $n, \gamma$, and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$.
Proof. - With no loss of generality, we assume that $\left\{p_{2}=0\right\}=\left\{x_{n}=0\right\}$. By Remark 4.8 we have

$$
\begin{equation*}
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}+\frac{2}{r} \int_{\mathrm{B}_{1}}\left(\lambda_{r} \hat{w}_{r}^{(\gamma)}-x \cdot \nabla \hat{w}_{r}^{(\gamma)}\right) \Delta \hat{w}_{r}^{(\gamma)} . \tag{4.20}
\end{equation*}
$$

- Step 1. We show that for some $\mathrm{C}=\mathrm{C}(n, \mathrm{P})$ we have

$$
\begin{equation*}
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-\mathrm{C} r^{4-\gamma} \phi^{\gamma}(r, w)\left(g^{\gamma}(r)+1\right) \tag{4.21}
\end{equation*}
$$

where

$$
g^{\gamma}(r):=\frac{\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)}^{2}}{\mathrm{H}(r, w)+r^{2 \gamma}} .
$$

Indeed, recall that by Lemma 4.7 we have

$$
r|\nabla w| \leq \mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathbf{B}_{5} \backslash \mathrm{~B}_{1}\right)}+r^{5}\right) \quad \text { in } \mathbf{B}_{r} .
$$

Now, notice that (recall that $p_{3}$ and P are divisible by $x_{n}$ )

$$
\frac{1}{2}\left(x_{n}+p_{3} / x_{n}+\mathrm{Q}+\mathrm{P} / x_{n}\right)^{2}=\mathscr{P}+\mathrm{P}+\mathrm{O}\left(|x|^{5}\right)
$$

we obtain

$$
\begin{aligned}
\left(x_{n}\right. & \left.+p_{3} / x_{n}+\mathrm{Q}+\mathrm{P} / x_{n}\right)(1+\mathrm{O}(|x|)) \\
& =\left(x_{n}+p_{3} / x_{n}+\mathrm{Q}+\mathrm{P} / x_{n}\right) \partial_{n}\left(x_{n}+p_{3} / x_{n}+\mathrm{Q}+\mathrm{P} / x_{n}\right) \\
& =\left|\partial_{n}(\mathscr{P}+\mathrm{P})\right|+\mathrm{O}\left(|x|^{4}\right) .
\end{aligned}
$$

Therefore, since Q is 3-homogenous and also divisible by $x_{n}$ and $\left|x_{n}\right| \leq \mathrm{C}|x|^{2}$ in $\{u=0\}$ (cf. (4.16)) we obtain

$$
\frac{1}{2}\left|x_{n}+p_{3} / x_{n}+\mathrm{P} / x_{n}\right| \leq \partial_{n}(\mathscr{P}+\mathrm{P})+\mathrm{O}\left(|x|^{4}\right) \quad \text { in }\{u=0\} .
$$

Thus, since $w:=u-\mathscr{P}-\mathrm{P}$, inside $\mathrm{B}_{r} \cap\{u=0\}$ (for $r$ small) we have

$$
\begin{align*}
r \frac{1}{2}\left|x_{n}+p_{3} / x_{n}+\mathrm{P} / x_{n}\right| & \leq r\left|\partial_{n}(\mathscr{P}+\mathrm{P})\right|+\mathrm{O}\left(r^{5}\right)=r\left|\partial_{n} w\right|+\mathrm{O}\left(r^{5}\right)  \tag{4.22}\\
& \leq \mathrm{C}\left(\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)}+r^{5}\right)
\end{align*}
$$

Therefore

$$
\begin{align*}
|w| & =|-(\mathscr{P}+\mathrm{P})|=\left|x_{n}+p_{3} / x_{n}+\mathrm{P} / x_{n}\right|^{2}+\mathrm{O}\left(r^{5}\right)  \tag{4.23}\\
& \leq 2 r^{2}\left|x_{n}+p_{3} / x_{n}+\mathrm{P} / x_{n}\right|+\mathrm{O}\left(r^{5}\right) \\
& \leq \mathrm{C}\left(r\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)}+r^{5}\right) \quad \text { in } \mathrm{B}_{r} \cap\{u=0\}
\end{align*}
$$

and, for $\alpha=1,2, \ldots, n-1$,

$$
\begin{align*}
r\left|\partial_{\alpha} w\right| & =r\left|-\partial_{\alpha}(\mathscr{P}+\mathrm{P})\right|=r\left|\partial_{\alpha}\left(p_{3} / x_{n}\right)\right|\left|x_{n}+p_{3} / x_{n}+\mathrm{P} / x_{n}\right|+\mathrm{O}\left(r^{5}\right)  \tag{4.24}\\
& \leq \mathrm{C} r^{2}\left|x_{n}+p_{3} / x_{n}+\mathrm{P} / x_{n}\right|+\mathrm{O}\left(r^{5}\right) \\
& \leq \mathrm{C}\left(r\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)}+r^{5}\right) \quad \text { in } \mathrm{B}_{r} \cap\{u=0\},
\end{align*}
$$

from which it follows that

$$
\begin{align*}
\left|\lambda_{r} \hat{w}_{r}^{(\gamma)}-x \cdot \nabla \hat{w}_{r}^{(\gamma)}\right| & \leq \frac{\mathrm{C}\left(1+\lambda_{r}\right)\left(r\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathbf{B}_{5} \backslash \mathbf{B}_{1}\right)}+r^{5}\right)}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{1 / 2}}  \tag{4.25}\\
& \leq \mathrm{C}\left(1+\lambda_{r}\right)\left(r g^{\gamma}(r)^{1 / 2}+r^{5-\gamma}\right) \\
& \leq \mathrm{C}\left(1+\lambda_{r}\right) r^{5-\gamma}\left(g^{\gamma}(r)^{1 / 2}+1\right) \quad \text { in } \mathrm{B}_{1} \cap\left\{u_{r}=0\right\} .
\end{align*}
$$

Also, since $\Delta \mathscr{P} \equiv 1$ and $\Delta \mathrm{P} \equiv 0$ we have $\Delta w=-\chi_{\{u=0\}} \leq 0$. Hence, arguing as in (3.10),

$$
\begin{equation*}
0 \geq \int_{\mathrm{B}_{1}} \Delta \hat{w}_{r}^{(\gamma)} \geq-\mathrm{C}\left\|\hat{w}_{r}^{(\gamma)}\right\|_{\mathrm{L}^{1}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)} \geq-\mathrm{C} \frac{\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)}}{\left(\mathrm{H}(r, w)+r^{2} \gamma\right)^{1 / 2}} \geq-\mathrm{C}^{\gamma}(r)^{1 / 2} . \tag{4.26}
\end{equation*}
$$

Combining this information with (4.25) we obtain

$$
\left|\int_{\mathrm{B}_{1}}\left(\lambda_{r} \hat{w}_{r}^{(\gamma)}-x \cdot \nabla \hat{w}_{r}^{(\gamma)}\right) \Delta \hat{w}_{r}^{(\gamma)}\right| \leq \mathrm{C}\left(1+\lambda_{r}\right) r^{5-\gamma}\left(g^{\gamma}(r)+1\right),
$$

that together with (4.20) yields (4.21).

- Step 2. Next we show that (4.21) implies that, for all $\gamma<5$, we have

$$
\begin{equation*}
\phi^{\gamma}(r, w) \leq \mathrm{C}_{\gamma} \quad \text { and } \quad g^{\gamma}(r) \leq \mathrm{C}_{\gamma}, \tag{4.27}
\end{equation*}
$$

where $\mathrm{C}_{\gamma}$ depends only on $n, \gamma$, and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathbf{B}_{1}\right)}$.
We prove $(4.27)$ for all $\gamma \in[3,5)$ by iteratively increasing the value of $\gamma$ at each iteration, starting from $\gamma=3$, in order to always have (along the iteration) a uniform bound on $\phi^{\gamma}(r, w)$ and $g^{\gamma}(r)$.

First, we observe that since $0 \in \Sigma_{n-1}^{3 r d}$ we have $\phi\left(0^{+}, u-p_{2}\right) \geq 3$, hence $\mid u-\mathscr{P}-$ $\mathrm{P} \mid \leq \mathrm{Cr}^{3}$ in $\mathrm{B}_{r}$, with a bound depending only on $n$ and P . This immediately implies that $g^{3}(r) \leq \mathrm{C}_{3}$, and then it follows by $(4.21)$ that $\phi^{3}(\cdot, w)$ is almost monotonically increasing, so in particular it is uniformly bounded.

Then, by the very same argument as the one used in Step 2 in the proof of Lemma 4.3 (using (4.21) in place of (4.4)) we deduce that if (4.27) holds for some $\gamma \geq 3$, then (4.27) holds also with $\gamma$ replaced by $\gamma+\beta$ for any $\beta>0$ such that $4 \beta<5-\gamma$. Thanks to this fact, with finitely many iterations we conclude that (4.27) holds for any $\gamma<5$, as claimed.

Combining (4.27) with (4.21), we obtain

$$
\frac{d}{d r} \phi^{\gamma}(r, w) \geq \frac{2}{r} \frac{\left(r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w\right)^{2}}{\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)^{2}}-\mathrm{CC}_{\gamma}^{2} r^{4-\gamma} \quad \forall r \in(0,1)
$$

Moreover, recalling (4.23) and (4.26), we conclude that

$$
\frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}}=\int_{\mathrm{B}_{1} \cap\left\{u_{r}=0\right\}} \hat{w}_{r} \Delta \hat{w}_{r} \geq-\mathrm{Cr}^{5-\gamma} g^{\gamma}(r) \geq-\mathrm{CC}^{\gamma} r^{5-\gamma} .
$$

Thanks to Lemma 4.9 we know that the truncated frequency function $\phi^{\gamma}$ is almost monotone for $\gamma<5$, and we can use this to study finer properties for points in $\Sigma_{n-1}^{3 r d}$. In particular, we introduce the following:

Definition 4.10. - Let $u: B_{1} \rightarrow[0, \infty)$ solve (3.1). We denote by $\Sigma_{n-1}^{4 t h}$ the set of points $x_{0} \in \Sigma_{n-1}^{3 r d}$ such that the following holds:

Set $w=u\left(x_{0}+\cdot\right)-\mathscr{P}_{x_{0}}$, where $\mathscr{P}_{x_{0}}$ is defined as in (4.10) starting from $p_{2, x_{0}}$ and $p_{3, x_{0}}$. Then there exists some sequence $r_{k} \downarrow 0$ along which $r_{k}^{-4} w\left(r_{k} \cdot\right)$ converges, weakly in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$, to some 4-homogeneous harmonic polynomial vanishing on $\left\{p_{2, x_{0}}=0\right\}$-possibly the polynomial zero.

We can now prove the existence of a unique 4-th order limit at points of $\Sigma_{n-1}^{4 \text { th }}$.
Lemma 4.11. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{\geq 4}$ (see Definition 4.13 below). Set $w:=u-\mathscr{P}-\mathrm{P}$, where $\mathscr{P}$ is defined in (4.10), and P is a 4-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$. Then, for any $\gamma \in(4,5)$ we have

$$
\begin{equation*}
\frac{d}{d r} \log \left(r^{-8}\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)\right) \geq-\mathrm{C}^{4-\gamma} \tag{4.28}
\end{equation*}
$$

where C is a constant depending only on $n, \gamma$, and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$.
As a consequence, for all $x_{\circ} \in \Sigma_{n-1}^{4 \text { th }}$ the limit

$$
p_{4, x_{0}}:=\lim _{r \downarrow 0} \frac{1}{r^{4}}\left(u\left(x_{\circ}+r \cdot\right)-\mathscr{P}_{x_{0}}(r \cdot)\right)
$$

exists, and it is a 4-homogeneous harmonic polynomial vanishing on $\left\{p_{2, x_{0}}=0\right\}$.
Proof. - For every 4-homogeneous harmonic polynomial P vanishing on $\left\{p_{2}=0\right\}$, we have

$$
\frac{d}{d r} \log \left(r^{-8}\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)\right)=\frac{2}{r}\left(\phi^{\gamma}(r, w)-4\right)+\frac{r^{2-n} \int_{\mathrm{B}_{r}} w \Delta w}{\mathrm{H}(r, w)+r^{2 \gamma}} \geq-\mathrm{C} r^{4-\gamma}
$$

where we used Lemma 4.9. This proves (4.28).
Now, if $0 \in \Sigma_{n-1}^{4 \text { th }}$ then we have that, for some $r_{k} \downarrow 0$ and some P which is 4homogeneous harmonic and vanishes on $\left\{p_{2}=0\right\}$,

$$
\log \left(r_{k}^{-8}\left(\mathrm{H}\left(r_{k}, w\right)+r_{k}^{2 \gamma}\right)\right) \rightarrow-\infty
$$

Thus, thanks to (4.28) we have

$$
\lim _{r \downarrow 0} \log \left(r^{-8}\left(\mathrm{H}(r, w)+r^{2 \gamma}\right)\right)=-\infty
$$

which implies that $r^{-4}(u-\mathscr{P})(r x) \rightarrow \mathrm{P}=: p_{4,0}$.
When $x_{\circ}=0$ we will simplify the notation $p_{4,0}$ to $p_{4}$.
We can now prove an enhanced version of Proposition 3.9 for higher-order blowups in $\Sigma_{n-1}^{3 r d}$ and $\Sigma_{n-1}^{4 t h}$.

Proposition 4.12. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and let $\mathscr{P}$ be as in Definition 4.5.
(a) Let $0 \in \Sigma_{n-1}^{3 r d} \backslash \Sigma_{n-1}^{4 t h}$ and set $w:=u-\mathscr{P}$. Then the limit $\lambda^{3 r d}:=\lim _{r \downarrow 0} \phi(r, w)$ exists and satisfies $\lambda^{3 r d} \in[3,4]$. Moreover, for every sequence $r_{k} \downarrow 0$ there is a subsequence $r_{k e}$ such that $\tilde{w}_{r_{k}} \rightarrow q$ as $\ell \rightarrow \infty$ in $\mathrm{C}_{\mathrm{loc}}^{0}\left(\mathbf{R}^{n}\right)$ and weakly in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$, where $q \not \equiv 0$ is a global $\lambda^{3 r d}$-homogeneous solution of the Signorini problem (3.12). In addition, if $\lambda^{3 \text { th }}<4$ then $q$ is even with respect to $\left\{p_{2}=0\right\}$.
(b) Let $0 \in \Sigma_{n-1}^{4 t h}$ and set $w:=u-\mathscr{P}-p_{4}$. Then:
(b1) either $\mathrm{H}(r, w)^{1 / 2} \leq \mathrm{C}_{\zeta} r^{5-\zeta}$ for all $\zeta \in(0,1)$, for some $\mathrm{C}_{\zeta}$ depending on $\zeta$;
(b2) or the limit $\lambda^{4 t h}:=\lim _{r \downarrow 0} \phi(r, w)$ exists and satisfies $\lambda^{4 \text { th }} \in[4,5)$. Moreover, for every sequence $r_{k} \downarrow 0$ there is a subsequence $r_{k \ell}$ such that $\tilde{w}_{r_{k \ell}} \rightharpoonup q$ as $\ell \rightarrow \infty$ in $\mathbf{C}_{\mathrm{loc}}^{0}\left(\mathbf{R}^{n}\right)$ and weakly in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$, where $q \not \equiv 0$ is a $\lambda^{4 \text { th }}-$ homogeneous solution of (3.12), even with respect to $\left\{p_{2}=0\right\}$.

Proof. - (a) Let $0 \in \Sigma_{n-1}^{3 r d} \backslash \Sigma_{n-1}^{4 t h}$, and $w:=u-\mathscr{P}$. Let $\gamma:=5-\varepsilon \in(4,5)$.
We note that $\phi^{\gamma}\left(0^{+}, w\right) \leq 4$. Indeed, if by contradiction $\phi^{\gamma}\left(0^{+}, w\right)>4$ then by Lemma 4.1 (which can be applied thanks to Lemma 4.9) we would have

$$
\mathrm{H}(r, w)+r^{2 \gamma} \leq \mathrm{C} r^{\phi^{\gamma}\left(0^{+}, w\right)}(\mathrm{H}(1, w)+1) \ll r^{8} \quad \text { as } r \downarrow 0,
$$

and hence we would have $r^{-4} w_{r} \rightarrow 0$ and in particular $0 \in \Sigma_{n-1}^{4 \text { th }}$ (with $p_{4} \equiv 0$ ), contradicting our assumption.

Note now that since $\gamma>4$ we have

$$
\begin{equation*}
\phi^{\gamma}\left(0^{+}, w\right)=\lim _{r \downarrow 0} \frac{\mathrm{D}(r, w)+\gamma r^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} \leq 4<\gamma . \tag{4.29}
\end{equation*}
$$

Also, using Lemma 4.1,

$$
(4.29) \quad \Rightarrow \quad \frac{r^{2 \gamma}}{\mathrm{H}(r, w)} \downarrow 0 \quad \Rightarrow \quad \lambda^{3 r d}:=\phi^{\gamma}\left(0^{+}, w\right)=\phi\left(0^{+}, w\right)
$$

Thus, the limit $\phi\left(0^{+}, w\right)$ exists and equals again $\lambda^{3 r d} \leq 4$. In addition, by Lemmas 4.1 and 4.9 we have

$$
\left\|\tilde{w}_{r}\right\|_{\mathrm{w}^{1,2}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R}) \quad \forall r>0
$$

for each $\mathrm{R} \geq 1$, which gives weak compactness in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$ of $\tilde{w}_{r}$ as $r \downarrow 0$.
We now show that any "accumulation point"

$$
q:=\lim _{k} \tilde{w}_{r_{k}}
$$

satisfies

$$
\begin{equation*}
c r^{-\lambda^{3 r d}} \mathrm{H}(r, q)^{1 / 2} \leq \mathrm{H}(1, q)^{1 / 2}=1 \quad \forall r \in(0,1) \tag{4.30}
\end{equation*}
$$

and, for all $\delta>0$,

$$
\begin{equation*}
c_{\delta} \mathrm{R}^{-\lambda^{3 / d}-\delta} \mathrm{H}(\mathrm{R}, q)^{1 / 2} \leq \mathrm{H}(1, q)^{1 / 2}=1 \quad \forall \mathrm{R} \in(1, \infty), \tag{4.31}
\end{equation*}
$$

where $c, c_{\delta}$ are positive constants.

Indeed, since $\phi^{\gamma}\left(0^{+}, w\right)=\lambda^{3 r d}$, Lemma 4.1 gives (for $0<r<\mathrm{R} \ll 1$ )

$$
c(\mathrm{R} / r)^{2 \lambda^{3 r d}} \leq \frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}}=\frac{\mathrm{H}\left(\mathrm{R} / r_{k}, \tilde{w}_{r_{k}}\right)+\frac{\mathrm{R}^{2 \gamma}}{\mathrm{H}\left(r_{r} w\right)}}{\mathrm{H}\left(r / r_{k}, \tilde{w}_{r_{k}}\right)+\frac{r^{2 \gamma}}{\mathrm{H}\left(r_{k}, w\right)}} .
$$

In particular, replacing R by $r_{k}$ and $r$ by $r_{k} r$ (with $r<1$ ) we obtain

$$
c r^{-2 \lambda^{3 r d}} \leq \frac{\mathrm{H}\left(1, \tilde{w}_{r_{k}}\right)+\frac{r_{k}^{2 \gamma}}{\mathrm{H}\left(r_{r} w\right)}}{\mathrm{H}\left(r, \tilde{w}_{r_{k}}\right)+\frac{\left(\frac{\left(r_{k}\right)^{2 \gamma}}{\mathrm{H}\left(r_{k}, w\right)}\right.}{\left.\mathrm{l}, \tilde{w}^{2}\right)}} \leq \frac{\mathrm{H}\left(1, \tilde{w}_{r_{k}}\right)+\frac{r_{k}^{2 \gamma}}{\mathrm{H}\left(r_{k}, w\right)}}{\mathrm{H}\left(r, \tilde{w}_{r_{k}}\right)} \rightarrow \frac{\mathrm{H}(1, q)}{\mathrm{H}(r, q)},
$$

proving (4.30).
Similarly, by the other inequality in Lemma 4.1 we have (for $0<r<\mathrm{R} \ll 1$ )

$$
c_{\delta}(\mathrm{R} / r)^{2 \lambda^{3 r d}}+\delta \geq \frac{\mathrm{H}(\mathrm{R}, w)+\mathrm{R}^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}}=\frac{\mathrm{H}\left(\mathrm{R} / r_{k}, \tilde{w}_{r_{k}}\right)+\frac{\mathrm{R}^{2 \gamma}}{\mathrm{H}\left(r_{k} w\right)}}{\mathrm{H}\left(r / r_{k}, \tilde{w}_{r_{k}}\right)+\frac{r^{2 \gamma}}{\mathrm{H}\left(r_{k}, w\right)}} .
$$

In particular, replacing $\mathbf{R}$ by $r_{k} \mathbf{R}$ and $r$ by $r_{k}($ with $\mathbf{R}>1)$ we obtain

$$
c_{\delta} \mathrm{R}^{2 \gamma} \geq \frac{\mathrm{H}\left(\mathrm{R}, \tilde{w}_{r_{k}}\right)+\mathrm{R}^{2 \gamma} \frac{r_{k_{k}}^{2 \gamma}}{\mathrm{H}\left(\mathrm{R}_{r_{k}} w\right)}}{\mathrm{H}\left(1, \tilde{w}_{r_{k}}\right)+\frac{\left(r_{k}\right)}{\mathrm{H}\left(r_{k}, w\right)}} \geq \frac{\mathrm{H}\left(\mathrm{R}, \tilde{w}_{r_{k}}\right)}{\mathrm{H}\left(1, \tilde{w}_{r_{k}}\right)+\frac{\left(r_{k}\right)^{2 \gamma}}{\mathrm{H}\left(r_{k}, w\right)}} \rightarrow \frac{\mathrm{H}(\mathrm{R}, q)}{\mathrm{H}(1, q)},
$$

which proves (4.31).
We now note that $\Delta w=-\chi_{\{u=0\}}$ implies that $\Delta \tilde{w}_{k}$ and $\Delta q$ are nonpositive measures. Also, the Lipschitz estimate in Lemma 4.7 (with $\mathrm{P} \equiv 0)$ implies that

$$
\left\|\tilde{w}_{r_{k}}\right\|_{\mathrm{Lip}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R}), \quad \text { and thus } \quad\|q\|_{\mathrm{Lip}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R})
$$

for all $\mathrm{R} \geq 1$. As a consequence, the convergence of $\tilde{w}_{r_{k}}$ to $q$ is uniform on compact sets. Furthermore, by (4.10) we have $w=u-\mathscr{P} \geq-\mathscr{P} \geq-\mathrm{O}\left(|x|^{5}\right)$ on $\left\{x_{n}=p_{3} / x_{n}\right\}$, so by uniform convergence we obtain

$$
q \geq 0 \quad \text { on }\left\{x_{n}=0\right\}
$$

In addition, by Lemma 4.9 we have

$$
\int_{\mathrm{B}_{\mathrm{R}}} \tilde{w}_{r_{k}} \Delta \tilde{w}_{r_{k}} \downarrow 0
$$

and since $\tilde{w}_{r_{k}} \rightarrow q$ in $\mathrm{C}^{0}$ and $0 \geq \Delta w_{k} \rightharpoonup^{*} \Delta q$ (up to extracting a further subsequence) we obtain

$$
\int_{\mathrm{B}_{\mathrm{R}}} q \Delta q=0 \quad \forall r \geq 1
$$

But since $\Delta \tilde{w}_{r}$ is supported in $\left\{u_{r}=0\right\} \subset\left\{\left|x_{n}\right| \leq o(1)\right\}$ as $r \downarrow 0$, the support the nonpositive measure $\Delta q$ is contained on $\left\{x_{n}=0\right\}$ where $q \geq 0$. We have thus shown that $q: \mathbf{R}^{n} \rightarrow \mathbf{R}$ is a solution of the Signorini problem (3.12).

Finally, recalling (4.30) and (4.31) we have that $\phi\left(0^{+}, q\right) \geq \lambda^{3 r d}$ while $\phi(+\infty, q) \leq$ $\lambda^{3 r d}+\delta$ for all $\delta>0$. This implies that $\phi(r, q)=\lambda^{3 r d}$ for all $r>0$, and thus $q$ must be $\lambda^{3 r d}$-homogeneous (see Lemma A.3).

Note also that, by Lemma 4.1, we have $\mathrm{H}(w, r) \gg r^{\lambda^{3 n d}}+\delta$ for every $\delta>0$. Thus, since by definition of $\Sigma_{n-1}^{3 r d}$ we have $\phi\left(0^{+}, u-p_{2}\right) \geq 3$, this implies $\left|u-p_{2}\right| \leq \mathrm{C}|x|^{3}$ and thus $|u-\mathscr{P}| \leq \mathrm{C}|x|^{3}$. Therefore it must be $\lambda^{3 r d} \geq 3$.

To conclude part (a), we need to show that if $\lambda^{3 r d}<4$ then $q$ is even. For this, notice that if one writes $q$ as the sum of its even and odd part, then the odd part is harmonic. Thus, if $\lambda^{3 r d} \in(3,4)$ then any $\lambda^{3 r d}$-homogeneous solution of the Signorini problem is even (since the homogeneity of a harmonic function is always an integer), so we only need to understand the case $\lambda^{3 r d}=3$.

Assume $\lambda^{3 r d}=3$, and let us show that $q$ is even. We have (see Lemma 4.1) that $\mathrm{H}(w, r) \gg r^{3+\delta}$ for all $\delta>0$ as $r \downarrow 0$ therefore

$$
q:=\lim _{k} \tilde{w}_{r_{k}}=\lim _{k} \frac{\left(u-p_{2}-p_{3}\right)\left(r_{k} \cdot\right)}{\mathrm{H}\left(r_{k}, u-p_{2}-p_{3}\right)^{1 / 2}} .
$$

Moreover, using (3.19) from Lemma 3.12 we obtain (note that $w$ in this proof and in Lemma 3.12 are different)

$$
\begin{equation*}
\int_{\partial \mathrm{B}_{1}}\left(\frac{\left(u-p_{2}-p_{3}\right)_{r}}{r^{3}}+\mathrm{P}\right)^{2}+\mathrm{C}(\mathrm{P}) r \geq \lim _{r \downarrow 0} \int_{\partial \mathrm{B}_{1}}\left(\frac{\left(u-p_{2}-p_{3}\right)_{r}}{r^{3}}+\mathrm{P}\right)^{2} \geq \int_{\partial \mathrm{B}_{1}} \mathrm{P}^{2} \tag{4.32}
\end{equation*}
$$

for all P harmonic 3-homogeneous vanishing on $\left\{p_{2}=0\right\}$, therefore

$$
-\mathrm{C}(\mathrm{P}) r \leq \int_{\partial \mathrm{B}_{1}}\left[\left(\frac{\left(u-p_{2}-p_{3}\right)_{r}}{r^{3}}+\mathrm{P}\right)^{2}-\mathrm{P}^{2}\right] .
$$

As a consequence, expanding the square and dividing by

$$
\varepsilon_{r}:=\left(\int_{\partial \mathrm{B}_{1}}\left(\frac{\left(u-p_{2}-p_{3}\right)_{r}}{r^{3}}\right)^{2}\right)^{1 / 2}=o(1),
$$

since $r^{\delta} \ll \varepsilon_{r}$ as $r \downarrow 0$ we obtain

$$
-\mathrm{C}(\mathrm{P}) \frac{r_{k}}{\varepsilon_{r_{k}}} \leq \int_{\partial \mathbf{B}_{1}}\left[\varepsilon_{r}\left(\frac{\left(u-p_{2}-p_{3}\right)_{r}}{\mathrm{H}\left(r_{k}, u-p_{2}-p_{3}\right)^{1 / 2}}\right)^{2}+2 \frac{\left(u-p_{2}-p_{3}\right)_{r}}{\mathrm{H}\left(r_{k}, u-p_{2}-p_{3}\right)^{1 / 2}} \mathrm{P}\right],
$$

and in the limit as $r \downarrow 0$ we get

$$
0 \leq 2 \int_{\partial \mathrm{B}_{1}} q \mathrm{P}
$$

for every odd harmonic 3-homogeneous polynomial P . Since if P is an odd harmonic 3 -homogeneous polynomial then so is -P , we deduce that $q$ must be orthogonal to all odd harmonic polynomials, hence $q$ is even.
(b) We assume that (bl) fails and we prove (b2). If (bl) fails then there exist $\zeta \in(0,1)$ and a sequence $r_{k} \rightarrow 0$ such that $\mathrm{H}\left(r_{k}, w\right)^{1 / 2} \geq \mathrm{C}_{\zeta} r_{k}^{5-\zeta}$. In particular, there exists some $\gamma \in(4,5)$ such that

$$
\begin{equation*}
\phi^{\gamma}\left(0^{+}, w\right)=\lim _{r \downarrow 0} \frac{\mathrm{D}(r, w)+\gamma r^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}}<\gamma \tag{4.33}
\end{equation*}
$$

Thus, as in the proof of (a),

$$
\begin{aligned}
(4.33) & \Rightarrow \quad \frac{r^{2 \gamma}}{\mathrm{H}(r, w)} \downarrow 0 \Rightarrow \lambda^{4 t h}:=\phi\left(0^{+}, w\right)=\phi^{\gamma}\left(0^{+}, w\right) \\
\forall \gamma & \in\left(\lambda^{4 t h}, 5\right) .
\end{aligned}
$$

In addition, combining Lemmas 4.1 and 4.9 we obtain that

$$
\left\|\tilde{w}_{r}\right\|_{\mathrm{w}^{1,2}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R}) \quad \forall r>0
$$

for each $\mathrm{R} \geq 1$, which gives compactness of sequences $\tilde{w}_{r_{k}}$ as $r_{k} \downarrow 0$ - they converge weakly in $\mathrm{W}^{1,2}\left(\mathrm{~B}_{\mathrm{R}}\right)$ for every R up to extracting a subsequence. Also, as in (a), it follows by Lemma 4.1 that any "accumulation point"

$$
q:=\lim _{k} \tilde{w}_{r_{k}}
$$

satisfies (4.30) and (4.31) with $\lambda^{3 r d}$ replaced by $\lambda^{4 t h}$. Also, exactly as in (a) we have that $\Delta \tilde{w}_{k}$ and $\Delta q$ are nonpositive measures, and

$$
\left\|\tilde{w}_{r_{k}}\right\|_{\mathrm{Lip}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R}) \quad \text { and thus } \quad\|q\|_{\mathrm{Lip}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R})
$$

for all $\mathrm{R} \geq 1$. As a consequence, the convergence is uniform on compact sets. Furthermore (4.10) yields $w=u-\mathscr{P}-p_{4} \geq-\mathscr{P}-p_{4} \geq-\mathrm{O}\left(|x|^{5}\right)$ on $\left\{x_{n}=p_{3} / x_{n}\right\}$, so by uniform convergence of $\tilde{w}_{r_{k}}$ to $q$ we obtain

$$
q \geq 0 \quad \text { on }\left\{x_{n}=0\right\} .
$$

Also, by Lemma 4.9 we obtain $\int_{\mathrm{BR}_{\mathrm{R}}} \tilde{w}_{r_{k}} \Delta \tilde{w}_{r_{k}} \downarrow 0$ from which we deduce that $\int_{\mathrm{B}_{\mathrm{R}}} q \Delta q=0$ for all $\mathrm{R}>1$. As a consequence, $q$ is a solution of the Signorini problem (3.12). Finally,
recalling (4.30) and (4.31) we have that $\phi\left(0^{+}, q\right) \geq \lambda^{4 \text { th }}$ while $\phi(+\infty, q) \leq \lambda^{4 t h}+\delta$ for all $\delta>0$. This implies that $\phi(r, q)=\lambda^{4 t h}$ for all $r>0$, and thus $q$ must be $\lambda^{4 t h}$-homogeneous.

Note also that by Lemma 4.1 we have $\mathrm{H}(w, r) \gg r^{\lambda^{4 h}}+\delta$ for every $\delta>0$. Thus, since by definition of $\Sigma_{n-1}^{4 t h}$ we have $|u-\mathscr{P}| \leq \mathrm{C}|x|^{4}$, it must be $\lambda^{4 \text { th }} \geq 4$.

Finally, we prove that $q$ must be even. As in (a), we only need to understand the case $\lambda^{4 t h}=4$. When $\lambda^{4 t h}=4$ then we have (see Lemma 4.9) $\mathrm{H}(w, r) \gg r^{4+\delta}$ for all $\delta>0$ as $r \downarrow 0$. On the other hand, by definition of $p_{4}$ it follows that $\mathrm{H}\left(r_{k}, u-\mathscr{P}-p_{4}\right)^{1 / 2}=r_{k}^{4} \varepsilon_{k}$, where $r_{k}^{\delta} \ll \varepsilon_{k}=o(1)$. Then, using Lemma 4.11 we obtain

$$
\begin{aligned}
r^{-8} \mathrm{H}(r, u-\mathscr{P}-\mathrm{P}) & \geq-\mathrm{C}(\mathrm{P}) r^{\delta}+\lim _{r \downarrow 0} r^{-8} \mathrm{H}(r, u-\mathscr{P}-\mathrm{P}) \\
& =-\mathrm{C}(\mathrm{P}) r^{\delta}+\mathrm{H}\left(1, p_{4}-\mathrm{P}\right),
\end{aligned}
$$

for all P quartic vanishing on $\left\{x_{n}=0\right\}$. Therefore, similarly to (a), we deduce that $q$ is orthogonal to every odd harmonic 4-homogeneous polynomial. Since $q$ is a solution of Signorini this implies that its odd part (which is harmonic) must vanish, concluding the proof.

We can now introduce the following:
Definition 4.13. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and recall the definition of $\mathscr{P}_{x_{0}}$ in (4.10).

We denote by $\Sigma_{n-1}^{\geq 4}$ the set of points $x_{0} \in \Sigma_{n-1}^{3 r d}$ such that, for $w:=u\left(x_{0}+\cdot\right)-\mathscr{P}_{x_{0}}$, we have $\phi^{\gamma}\left(0^{+}, w\right) \geq 4$ for every $\gamma \in(4,5)$.

We denote by $\Sigma_{n-1}^{>4}$ the set of points $x_{0} \in \Sigma_{n-1}^{4 \text { th }}$ such that, for $w:=u\left(x_{\circ}+\cdot\right)-\mathscr{P}_{x_{0}}-p_{4, x_{0}}$, we have $\phi^{\gamma}\left(0^{+}, w\right)>4$ for every $\gamma \in(4,5)$.

Furthermore, for fixed $\zeta \in(0,1)$ we denote by $\Sigma_{n-1}^{\geq 5-\zeta}$ the set of points $x_{\circ} \in \Sigma_{n-1}^{4 \text { th }}$ such that, for $w:=u\left(x_{\circ}+\cdot\right)-\mathscr{P}_{x_{0}}-p_{4, x_{0}}$, we have $\phi^{\gamma}\left(0^{+}, w\right) \geq 5-\zeta$ for any $\gamma \in(5-\zeta, 5)$.

Our last goal of this section is to show that $\Sigma_{n-1}^{>4}=\Sigma_{n-1}^{4 t h}$. For this, we need a new monotonicity formula.

Lemma 4.14. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{4 \text { th }}$. Let $w:=u-\mathscr{P}-p_{4}$, where $\mathscr{P}$ is defined in (4.10), and let P be any 4-homogeneous harmonic polynomial such that $\mathrm{P} \geq 0$ on $\left\{p_{2}=0\right\}$. Then

$$
\frac{d}{d r}\left(r^{-4} \int_{\partial \mathrm{B}_{1}} w_{r} \mathrm{P}\right) \leq \mathrm{C}
$$

where C is a constant depending only on $n$ and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$.

Proof. - After a rotation, we may assume $p_{2}=\frac{1}{2} x_{n}^{2}$. We have

$$
\begin{aligned}
\frac{d}{d r} \int_{\partial \mathrm{B}_{1}} w_{r} \mathrm{P} & =\int_{\partial \mathrm{B}_{1}} \frac{x}{r} \cdot \nabla w_{r} \mathrm{P}=\frac{1}{r} \int_{\partial \mathrm{B}_{1}} \partial_{v} w_{r} \mathrm{P}=\frac{1}{r} \int_{\mathrm{B}_{1}} \operatorname{div}\left(\nabla w_{r} \mathrm{P}\right) \\
& =\frac{1}{r}\left(\int_{\mathrm{B}_{1}} \nabla w_{r} \nabla \mathrm{P}+\int_{\mathrm{B}_{1}} \Delta w_{r} \mathrm{P}\right) \\
& =\frac{1}{r}\left(\int_{\partial \mathrm{B}_{1}} w_{r} \partial_{v} \mathrm{P}-\int_{\mathrm{B}_{1}} w_{r} \Delta \mathrm{P}+\int_{\mathrm{B}_{1}} \Delta w_{r} \mathrm{P}\right) \\
& =\frac{1}{r}\left(4 \int_{\partial \mathrm{B}_{1}} w_{r} \mathrm{P}+\int_{\mathrm{B}_{1}} \Delta w_{r} \mathrm{P}\right)
\end{aligned}
$$

where we used that $\partial_{\nu} \mathrm{P}=4 \mathrm{P}$ on $\partial \mathrm{B}_{1}$, and that $\Delta \mathrm{P}=0$. Now, since $\Delta w_{r}=-r^{2} \chi_{\left\{u_{r}=0\right\}}$, we deduce that

$$
\frac{d}{d r}\left(r^{-4} \int_{\partial \mathrm{B}_{1}} w_{r} \mathrm{P}\right)=-\frac{1}{r^{3}} \int_{\mathrm{B}_{1} \cap\left\{u_{r}=0\right\}} \mathrm{P} .
$$

Finally notice that (4.22) rescaled implies, using $\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5} \backslash \mathrm{~B}_{1}\right)} \leq \mathrm{C} r^{4}$ since $0 \in \Sigma_{n-1}^{4 \text { th }}$,

$$
\left\{u_{r}=0\right\} \cap \mathrm{B}_{1} \subset\left\{\left|x_{n}+r p_{3} / x_{n}\right| \leq \mathrm{Cr} r^{2}\right\} \quad \text { and thus } \quad\left|\left\{u_{r}=0\right\} \cap \mathrm{B}_{1}\right| \leq \mathrm{Cr} r^{2}
$$

Moreover, since $\mathrm{P} \geq 0$ on $\left\{x_{n}=0\right\}$, we have $\mathrm{P} \geq-\mathrm{C}\left|x_{n}\right|$ in $\mathrm{B}_{1}$. Hence we obtain

$$
-\int_{\mathrm{B}_{1} \cap\left\{u_{r}=0\right\}} \mathrm{P} \leq \mathrm{C} r\left|\left\{u_{r}=0\right\} \cap \mathrm{B}_{1}\right| \leq \mathrm{C} r^{3},
$$

and the lemma follows.
We can now show the following:
Proposition 4.15. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1). Then $\Sigma_{n-1}^{4 \text { th }}=\Sigma_{n-1}^{>4}$.
Proof. - Assume by contradiction that $0 \in \Sigma_{n-1}^{4 t h} \backslash \Sigma_{n-1}^{>4}$. Then, by Proposition $4.12(\mathrm{~b})$, there is a sequence $r_{k} \rightarrow 0$ along which $\tilde{w}_{r_{k}} \rightarrow q$ locally uniformly in $\mathbf{R}^{n}$, where $q$ is a 4-homogeneous even solution of the Signorini problem (3.12). Then, by [GP09, Lemma 1.3.4], $q$ is a harmonic polynomial.

Let $w:=u-\mathscr{P}-p_{4}$. Since $r^{-4} w_{r} \rightarrow 0$ (by definition of $p_{4}$ ), it follows by Lemma 4.14 that

$$
\int_{\partial \mathrm{B}_{1}} r^{-4} w_{r} \mathrm{P} \leq \mathrm{C} r
$$

for any 4-homogeneous harmonic polynomial P vanishing on $\left\{p_{2}=0\right\}$.

Set now $\tilde{w}_{r}=w_{r} / H\left(1, w_{r}\right)^{1 / 2}$ and $\varepsilon_{r}:=r^{-4} H\left(1, w_{r}\right)^{1 / 2}$, and notice that, since $0 \notin$ $\Sigma_{n-1}^{>4}$, for any $\delta>0$ we have $\varepsilon_{r} \gg r^{\delta}$ for $r>0$ small enough. Hence,

$$
\mathrm{C} r \geq \int_{\partial \mathrm{B}_{1}} r^{-4} w_{r} \mathrm{P}=\int_{\partial \mathrm{B}_{1}} \varepsilon_{r} \tilde{w}_{r} \mathrm{P} .
$$

Dividing by $\varepsilon_{r}$, and letting $r=r_{k} \rightarrow 0$, we deduce that

$$
0 \geq \int_{\partial \mathrm{B}_{1}} q \mathrm{P} .
$$

Taking $\mathrm{P}=q$, this provides the desired contradiction.

## 5. Uniqueness and nondegeneracy of non-harmonic cubic blow-ups

The goal of this section is to study the set $\Sigma_{n-1}^{\geq 3} \backslash \Sigma_{n-1}^{>3}$, namely the set of singular points where blow-ups are 3 -homogeneous and non-harmonic. ${ }^{4}$ As explained in the introduction, this study is crucial for our proof of Theorem 1.1.

We will prove that $\Sigma_{n-1}^{\geq 3} \backslash \Sigma_{n-1}^{3 r d}$ is contained in a countable union of $(n-2)$ dimensional Lipschitz manifolds, and that $\Sigma_{n-1}^{3 r d} \backslash \Sigma_{n-1}^{>3}=\emptyset$. For this, we will need to establish the uniqueness and nondegeneracy of blow-ups at these points.

We start by classifying all $\lambda$-homogeneous solutions of the Signorini problem in $\mathbf{R}^{n}$, with $\lambda$ odd.

Lemma 5.1.-Let $q: \mathbf{R}^{n} \rightarrow \mathbf{R}$ be a $\lambda$-homogenous solution of the Signorini problem

with homogeneity $\lambda=2 m+1, m \in \mathbf{N}$. Then $q \equiv 0$ on $\left\{x_{n}=0\right\}$.
Proof. - Using complex variables (so $i$ denotes the imaginary unit), for $\alpha \in$ $\{1,2, \ldots, n-1\}$ define

$$
\psi(x):= \begin{cases}i^{1-\lambda} \operatorname{Re}\left[\left(x_{n}+i x_{\alpha}\right)^{\lambda}\right] & x_{n} \geq 0 \\ -i^{1-\lambda} \operatorname{Re}\left[\left(x_{n}+i x_{\alpha}\right)^{\lambda}\right] & x_{n} \leq 0\end{cases}
$$

[^4]Note that

$$
\psi\left(x^{\prime}, x_{n}\right)=\psi\left(x^{\prime},-x_{n}\right) \quad \text { and } \quad \psi(x)=0 \quad \text { on }\left\{x_{n}=0\right\} .
$$

In addition, on $\left\{x_{n}=0\right\}$ we have $\partial_{n} \psi\left(x^{\prime}, 0^{+}\right)=\lambda\left|x_{\alpha}\right|^{\lambda-1}$ (recall that $\lambda-1$ is even), therefore

$$
\Delta \psi=\left.2 \lambda\left|x_{\alpha}\right|^{\lambda-1} \mathcal{H}^{n-1}\right|_{\left\{x_{n}=0\right\}} .
$$

On the other hand, since both $\psi$ and $q$ are $\lambda$-homogeneous we have $(x \cdot \nabla q) \psi=q(x$. $\nabla \psi)=\lambda q \psi$. Thus $\int_{\partial \mathrm{B}_{1}}\left(q_{\nu} \psi-q \psi_{\nu}\right)=0$, and an integration by parts gives

$$
\int_{\mathrm{B}_{1}} \Delta q \psi=\int_{\mathrm{B}_{1}} q \Delta \psi .
$$

Since $\Delta q$ is concentrated on $\left\{x_{n}=0\right\}$ where $\psi$ vanishes, combining all together we get

$$
0=\int_{\mathrm{B}_{1}} q \Delta \psi=2 \lambda \int_{\mathrm{B}_{1} \cap\left\{x_{n}=0\right\}} q\left|x_{\alpha}\right|^{\lambda-1} .
$$

Since $q \geq 0$ on $\left\{x_{n}=0\right\}$ and the previous equality holds for all $\alpha \in\{1,2, \ldots, n-1\}$, we conclude that $q$ must vanish on $\left\{x_{n}=0\right\}$.

Lemma 5.2. - Assume that $q: \mathbf{R}^{n} \rightarrow \mathbf{R}$ is a 3-homogenous even solution of the Signorini problem (5.1). Then, after a suitable rotation that leaves the hyperplane $\left\{x_{n}=0\right\}$ invariant, we have

$$
q(x)=b\left|x_{n}\right|^{3}-3\left|x_{n}\right|\left(\sum_{\alpha=1}^{n-1} b_{\alpha} x_{\alpha}^{2}\right),
$$

where $b, b_{\alpha} \geq 0$ and $b=\sum_{\alpha=1}^{n-1} b_{\alpha}$.
Proof. - By Lemma $5.1 q$ must vanish everywhere on $\left\{x_{n}=0\right\}$. Thus, $q$ is a 3homogenous harmonic function in $\left\{x_{n}>0\right\}$ vanishing on $\left\{x_{n}=0\right\}$, so its odd extension is a 3-homogeneous harmonic polynomial. This implies, after a rotation, that

$$
q(x)=b x_{n}^{3}-3 x_{n}\left(\sum_{\alpha=1}^{n-1} b_{\alpha} x_{\alpha}^{2}\right) \quad \text { for } x_{n}>0,
$$

where $b, b_{\alpha} \in \mathbf{R}$ and $b=\sum_{\alpha=1}^{n-1} b_{\alpha}$. Finally, since $q$ is an even solution of Signorini, it follows that $\partial_{n} q \leq 0$ on $\left\{x_{n}=0\right\}$. This implies that $b_{\alpha} \geq 0$ (and thus $b \geq 0$ ), concluding the proof.

In order to continue our analysis, we introduce a new monotonicity formula:

Lemma 5.3. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{\geq 3} \backslash \Sigma_{n-1}^{>3}$. Set $w:=u-p_{2}$ and $w_{r}:=w(r \cdot)$. Then, for fixed $\varrho \in(0,1)$ and for any 3-homogeneous solution $q$ of the Signorini problem (3.12), we have

$$
\frac{d}{d r} \int_{\partial \mathrm{B}_{e}} w_{r} q=\frac{3}{r} \int_{\partial \mathrm{B}_{e}} w_{r} q-\frac{\varrho}{r} \int_{\mathrm{B}_{e}} w_{r} \Delta q+\mathrm{O}\left(r^{3}\right) .
$$

In particular

$$
\frac{d}{d r}\left(\frac{1}{r^{3}} \int_{\partial \mathrm{B}_{1}} w_{r} q\right) \geq-\mathrm{C}
$$

Proof. - We have

$$
\begin{aligned}
\frac{d}{d r} \int_{\partial \mathrm{B}_{e}} w_{r} q & =\int_{\partial \mathrm{B}_{e}} \frac{x}{r} \cdot \nabla w_{r} q=\frac{\varrho}{r} \int_{\partial \mathrm{B}_{e}} \partial_{\nu} w_{r} q=\frac{\varrho}{r} \int_{\mathbf{B}_{e}} \operatorname{div}\left(\nabla w_{r} q\right) \\
& =\frac{\varrho}{r}\left(\int_{\mathrm{B}_{e}} \nabla w_{r} \nabla q+\int_{\mathrm{B}_{e}} \Delta w_{r} q\right) \\
& =\frac{\varrho}{r}\left(\int_{\partial \mathrm{B}_{e}} w_{r} \partial_{\nu} q-\int_{\mathrm{B}_{e}} w_{r} \Delta q+\int_{\mathbf{B}_{e}} \Delta w_{r} q\right) .
\end{aligned}
$$

Now, since $q$ is 3-homogeneous, we find that $\varrho \int_{\partial \mathrm{B}_{e}} w_{r} \partial_{\nu} q=3 \int_{\partial \mathrm{B}_{e}} w_{r} q$. To complete the proof of the Lemma we only need to show that $\int_{\mathrm{B}_{e}} \Delta w_{r} q=\mathrm{O}\left(r^{4}\right)$.

With no loss of generality, assume that $p_{2}=\frac{1}{2} x_{n}^{2}$. Then it follows by Proposition 3.10 that $\{u(r \cdot)=0\} \cap \mathrm{B}_{1} \subset\left\{\left|x_{n}\right| \leq \mathrm{Cr}\right\}$, and $|q| \leq \mathrm{C}\left|x_{n}\right|$ in $\mathrm{B}_{1}$ (by Lemma 5.2). Thus, since $\Delta w_{r}=-r^{2} \chi_{\{u(r \cdot)=0\}}$, we get $\int_{\mathrm{B}_{e}} \Delta w_{r} q=\mathrm{O}\left(r^{4}\right)$.

Finally, taking $\varrho=1$ and using that $-w_{r} \Delta q \geq 0$ in $\mathbf{R}^{n}$ (since $w_{r}=u(r \cdot) \geq 0$ on $\left\{x_{n}=0\right\}$ ), we obtain

$$
\frac{d}{d r}\left(\frac{1}{r^{3}} \int_{\partial \mathrm{B}_{1}} w_{r} q\right)=\frac{1}{r^{4}}\left(-\int_{\partial \mathrm{B}_{1}} w_{r} \Delta q+\int_{\mathrm{B}_{1}} \Delta w_{r} q\right) \geq \frac{1}{r^{4}} \int_{\mathrm{B}_{1}} \Delta w_{r} q \geq-\mathrm{C},
$$

as desired.

As a consequence of the previous lemma, we deduce the uniqueness of blow-ups in $\Sigma_{n-1}^{\geq 3} \backslash \Sigma_{n-1}^{3 r d}$. Notice that this is quite surprising, since even in the (simpler) case of the Signorini problem it was not known if cubic blow-ups are unique at every point (see Appendix B).

Proposition 5.4.-Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{\geq 3} \backslash \Sigma_{n-1}^{3 r d}$. Then the limit

$$
\tilde{q}:=\lim _{r \downarrow 0} \frac{\left(u-p_{2}\right)(r \cdot)}{r^{3}}
$$

exists, and it is a 3-homogeneous (non-harmonic) solution of Signorini.
Proof. - Let $w:=u-p_{2}$, and $w_{r}=w(r \cdot)$. Assume that

$$
q^{(i)}=\lim _{r_{k}^{(i)} \downarrow 0} \frac{1}{\left(r_{k}^{(i)}\right)^{3}} w_{r_{k}^{(i)}}, \quad i=1,2,
$$

are two accumulation points along different sequences $r_{k}^{(i)}$. Then, give a 3-homogeneous solution of Signorini $q$, we can apply Lemma 5.3 to deduce that $r \mapsto \frac{1}{r^{3}} \int_{\partial \mathrm{B}_{1}} w_{r} q$ has a unique limit as $r \rightarrow 0$. In particular this implies that

$$
\begin{equation*}
\int_{\partial \mathrm{B}_{1}} q^{(1)} q=\int_{\partial \mathrm{B}_{1}} q^{(2)} q . \tag{5.2}
\end{equation*}
$$

Choosing $q=q^{(1)}-q^{(2)}$ we obtain

$$
\int_{\partial \mathbf{B}_{1}}\left(q^{(1)}-q^{(2)}\right)^{2}=0,
$$

hence $q^{(1)} \equiv q^{(2)}$, as desired.
The next step consists in showing that if $0 \in \Sigma_{n-1}^{3 r d}$ then $\phi\left(0^{+}, u-p_{2}-p_{3}\right)>3$. This is a kind of nondegeneracy property, which implies that $\Sigma_{n-1}^{3 r d} \backslash \Sigma_{n-1}^{>3}$ is empty. This highly non-trivial fact is essential in order to establish Schaeffer conjecture in $\mathbf{R}^{4}$, and it is the core of this section. Its proof require a barrier and ODE-type arguments obtained below.

Lemma 5.5. - Let $u: B_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{3 r d}$. Set $w:=u-p_{2}-p_{3}$, and let $w_{r}$ and $\tilde{w}_{r}$ be defined as in (2.1). Assume that $\left\{p_{2}=0\right\}=\left\{x_{n}=0\right\}$, and given $x=$ $\left(x_{1}, \ldots, x_{n}\right) \in \mathbf{R}^{n}$ let $x^{\prime}:=\left(x_{1}, \ldots, x_{n-1}\right) \in \mathbf{R}^{n-1}$.

For any $\eta>0$ there exists $\delta=\delta(n, \eta)$ such that if

$$
\begin{aligned}
& \left\|\tilde{w}_{r}-q\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{2}\right)} \leq \delta \\
& \qquad \text { for } q=\left|x_{n}\right|\left(\frac{a}{3} x_{n}^{2}-x^{\prime} \cdot \mathrm{A} x^{\prime}\right), \mathrm{A} \in \mathbf{R}^{(n-1) \times(n-1)}, \mathrm{A} \geq 0, a=\operatorname{trace}(\mathrm{A})
\end{aligned}
$$

then

$$
u(r \cdot)=\mathrm{O}\left(r^{4}\right) \quad \text { on }\left\{x_{n}=0\right\} \cap\left(\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}\right) \cap\left\{x^{\prime} \cdot \mathrm{A} x^{\prime} \geq \eta\right\} .
$$

Proof. - Let $z=\left(z^{\prime}, 0\right) \in\left(\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}\right)$ satisfy $z^{\prime} \cdot \mathrm{A} z^{\prime} \geq \eta$, and given $c>0$, denote

$$
\phi_{z, c}(x):=\left(p_{2}+p_{3}\right)(r z+r x)-r^{3}(n-1)\left|x_{n}\right|^{2}+r^{3}\left|x^{\prime}\right|^{2}+c .
$$

Note that, since $q$ is uniformly close to $\tilde{w}_{r}$, the constant $a$ and the matrix A appearing in the definition of $q$ are universally bounded. Hence, there exists $\varrho>0$ small, depending only on $n$ and $\eta$, such that

$$
-n\left|x_{n}\right|^{2} \geq\left|x_{n}\right|\left(\frac{a}{3} x_{n}^{2}-\left(z^{\prime}+x^{\prime}\right) \cdot \mathrm{A}\left(z^{\prime}+x^{\prime}\right)\right) \quad \text { for }|x|<\varrho .
$$

Thus, denoting $h_{r}:=\mathrm{H}(r, w)^{1 / 2}=o\left(r^{3}\right)$, we have

$$
\begin{align*}
\phi_{z, c} & \geq\left(p_{2}+p_{3}\right)(r z+r x)-r^{3}(n-1)\left|x_{n}\right|^{2}+r^{3}\left|x^{\prime}\right|^{2}+c  \tag{5.3}\\
& >\left(p_{2}+p_{3}\right)(r z+r x)+h_{r} q(r z+r x)+r^{3}\left(\left|x_{n}\right|^{2}+\left|x^{\prime}\right|^{2}\right) \\
& \geq\left(u(r z+r x)-\delta h_{r}\right)+r^{3}|x|^{3} \quad \text { for }|x|<\varrho .
\end{align*}
$$

We now compare the two functions $\hat{u}_{z}(x):=u(r z+r x)$ and $\phi_{z, c}$ in $\mathrm{B}_{\varrho}(0)$. Two cases arise:
(1) either $\phi_{z, c} \geq u_{z}$ for each $c>0$, which implies that $0 \leq u(r z)=u_{z}(0) \leq \phi_{z, 0}(0)=$ 0 (since $p_{2}$ and $p_{3}$ vanish on $\left\{x_{n}=0\right\}$ );
(2) or there exists $c_{*}>0$ such that $\phi_{z, c_{*}}$ touches from above $\hat{u}_{z}$ at some point $y=$ $\left(y^{\prime}, y_{n}\right) \in \overline{\mathrm{B}_{\rho}}$. Note that $\Delta \phi_{z, c_{*}}=r^{2}$ in $\mathrm{B}_{\rho}$, and $\Delta \hat{u}_{z}=r^{2} \chi_{\left\{\hat{u}_{z}>0\right\}}$ in $\mathrm{B}_{\rho}$. Also, since $h_{r}:=$ $\mathrm{H}(r, w)^{1 / 2}=o\left(r^{3}\right)$, for $r$ small enough we have $\phi_{z, c} \geq \hat{u}_{z}(x)$ on $\partial \mathrm{B}_{\varrho}$ (by (5.3)). Thus, it follows by the maximum principle that the point $y$ must belong to $\left\{\hat{u}_{z}=0\right\} \cap \mathrm{B}_{\rho} \subset\left\{\left|x_{n}\right| \leq\right.$ $\mathrm{C} r\} \cap \mathrm{B}_{\rho}$, therefore

$$
\begin{aligned}
0 & =\hat{u}_{z}(y)=\phi_{z, c_{*}}(y)=\left(p_{2}+p_{3}\right)\left(r z^{\prime}+r y^{\prime}, r y_{n}\right)-r^{3}(n-1)\left|y_{n}\right|^{2}+r^{3}\left|y^{\prime}\right|^{2}+c_{*} \\
& \geq-\mathrm{C}^{4}+c_{*} .
\end{aligned}
$$

Thus $c_{*} \leq \mathrm{C} r^{4}$, and as a consequence

$$
0 \leq u(r z)=\hat{u}_{z}(0) \leq \phi_{z, c_{*}}(0)=c_{*} \leq \mathrm{C} r^{4}
$$

This proves that in both cases $0 \leq u(r z) \leq \mathrm{C} r^{4}$, and since $z \in\left\{x_{n}=0\right\} \cap\left(\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}\right) \cap\left\{x^{\prime}\right.$. A $\left.x^{\prime} \geq \eta\right\}$ is arbitrary, the result follows.

Another key tool is the following ODE-type formula.
Lemma 5.6. - Let $u: B_{1} \rightarrow[0, \infty)$ satisfy (3.1), and $0 \in \Sigma_{n-1}^{3 r d}$. Set $w:=u-p_{2}-p_{3}$, let $w_{r}$ and $\tilde{w}_{r}$ be defined as in $(2.1)$, and set $h(r):=\mathrm{H}(r, w)^{1 / 2}$. Assume that $\left\{p_{2}=0\right\}=\left\{x_{n}=0\right\}$, and
given a symmetric $(n-1) \times(n-1)$ matrix $\mathrm{A} \geq 0$, we define its "associated solution of the Signorini problem"

$$
\begin{equation*}
q_{\mathrm{A}}(x):=\left|x_{n}\right|\left(\frac{\operatorname{trace}(\mathrm{A})}{3} x_{n}^{2}-x^{\prime} \cdot \mathrm{A} x^{\prime}\right), \quad x=\left(x^{\prime}, x_{n}\right) \in \mathbf{R}^{n-1} \times \mathbf{R} \tag{5.4}
\end{equation*}
$$

and we introduce the quantity

$$
\begin{equation*}
\psi(r ; \mathrm{A}):=\int_{\partial \mathrm{B}_{1}} \tilde{w}_{r} q_{\mathrm{A}}-2 \int_{\partial \mathrm{B}_{1 / 2}} \tilde{w}_{r} q_{\mathrm{A}} . \tag{5.5}
\end{equation*}
$$

Then

$$
\frac{d}{d r} \psi(r ; \mathrm{A})=-\theta(r) \psi(r ; \mathrm{A})-\frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{A}}+\mathrm{O}\left(r^{3} / h(r)\right),
$$

where

$$
\theta(r):=\left(\frac{h^{\prime}(r)}{h(r)}+\frac{3}{r}\right)=\left(\log \left(h(r) / r^{3}\right)\right)^{\prime}
$$

Proof. - As in the proof of Lemma 5.3, we obtain

$$
\frac{d}{d r} \int_{\partial \mathrm{B}_{e}} w_{r} q_{\mathrm{A}}=\frac{3}{r} \int_{\partial \mathrm{B}_{e}} w_{r} q_{\mathrm{A}}-\frac{\varrho}{r} \int_{\mathrm{B}_{e}} w_{r} \Delta q_{\mathrm{A}}+\mathrm{O}\left(r^{3}\right) .
$$

Now, since $\tilde{w}_{r}=w_{r} / h(r)$ we deduce that

$$
\frac{d}{d r} \int_{\partial \mathrm{B}_{e}} \tilde{w}_{r} q_{\mathrm{A}}=\left(-\frac{h^{\prime}(r)}{h(r)}+\frac{3}{r}\right) \int_{\partial \mathrm{B}_{e}} \tilde{w}_{r} q-\frac{\varrho}{r} \int_{\mathrm{B}_{e}} \tilde{w}_{r} \Delta q_{\mathrm{A}}+\mathrm{O}\left(r^{3} / h(r)\right)
$$

and the lemma follows by combining the identities for $\varrho=1$ and $\varrho=1 / 2$.
We shall also need the following formula:
Lemma 5.7. - Given $\mathrm{A}, \overline{\mathrm{A}} \geq 0$ be two symmetric $(n-1) \times(n-1)$ matrices, let $q_{\mathrm{A}}$ and $q_{\overline{\mathrm{A}}}$ as defined in (5.4).

Then

$$
\int_{\partial \mathrm{B}_{e}} q_{\mathrm{A}} q_{\overline{\mathrm{A}}}=\frac{4 \varrho^{n+5}\left|\partial \mathrm{~B}_{1}\right|}{n(n+2)(n+4)}\left\{\operatorname{trace}(\mathrm{A} \cdot \overline{\mathrm{~A}})+\frac{1}{3} \operatorname{trace}(\mathrm{~A}) \operatorname{trace}(\overline{\mathrm{A}})\right\} .
$$

Proof.- Let $\mathrm{A}=\left(a_{\alpha \beta}\right)_{\alpha, \beta=1}^{n-1}, \overline{\mathrm{~A}}=\left(\bar{a}_{\alpha \beta}\right)_{\alpha, \beta=1}^{n-1}, \quad a=\operatorname{trace}(\mathrm{A})=\sum_{\alpha} a_{\alpha \alpha}, \bar{a}=$ $\operatorname{trace}(\overline{\mathrm{A}})=\sum_{\alpha} \bar{a}_{\alpha \alpha}$. Denote for brevity $q=q_{\mathrm{A}}, \bar{q}=q_{\overline{\mathrm{A}}}$. Then

$$
\int_{\partial \mathrm{B}_{e}} q \bar{q}=\sum_{\alpha, \beta, \gamma, \delta=1}^{n-1} \int_{\partial \mathrm{B}_{e}} x_{n}^{2}\left(a \frac{x_{n}^{2}}{3}-a_{\alpha \beta} x_{\alpha} x_{\beta}\right)\left(\bar{a} \frac{x_{n}^{2}}{3}-\bar{a}_{\gamma \delta} x_{\gamma} x_{\delta}\right) .
$$

Up to a rotation in the $\left\{x_{n}=0\right\}$ plane, we may assume that $a_{\alpha \beta}$ is diagonal. Noting that $\int_{\partial \mathrm{B}_{e}} x_{n}^{4} x_{\gamma} x_{\delta}=\int_{\partial \mathrm{B}_{e}} x_{n}^{2} x_{\alpha}^{2} x_{\gamma} x_{\delta}=0$ for $\gamma \neq \delta$, we have

$$
\begin{aligned}
\int_{\partial \mathbf{B}_{e}} q \bar{q}= & \int_{\partial \mathbf{B}_{e}}\left(\frac{a \bar{a}}{9} x_{n}^{6}+\sum_{\alpha}\left\{-\left(\frac{\bar{a}}{3} a_{\alpha \alpha}+\frac{a}{3} \bar{a}_{\alpha \alpha}\right) x_{n}^{4} x_{\alpha}^{2}+a_{\alpha \alpha} \bar{a}_{\alpha \alpha} x_{n}^{2} x_{\alpha}^{4}\right\}\right. \\
& \left.+\sum_{\alpha \neq \gamma} a_{\alpha \alpha} \bar{a}_{\gamma \gamma} x_{n}^{2} x_{\alpha}^{2} x_{\gamma}^{2}\right)
\end{aligned}
$$

We observe that

$$
\begin{aligned}
\int_{\partial \mathrm{B}_{1}} x_{i}^{4} & =\frac{1}{4} \int_{\partial \mathrm{B}_{1}} \partial_{v}\left(x_{i}^{4}\right)=\frac{1}{4} \int_{\mathrm{B}_{1}} \Delta\left(x_{i}^{4}\right)=3 \int_{\mathrm{B}_{1}} x_{i}^{2}=\frac{3}{n+2} \int_{\partial \mathrm{B}_{1}} x_{i}^{2} \\
& =\frac{3}{n(n+2)}\left|\partial \mathrm{B}_{1}\right| .
\end{aligned}
$$

Similarly,

$$
\begin{aligned}
\int_{\partial \mathbf{B}_{1}} x_{i}^{6} & =\frac{1}{6} \int_{\mathbf{B}_{1}} \Delta\left(x_{i}^{6}\right)=5 \int_{\mathbf{B}_{1}} x_{i}^{4}=\frac{5}{n+4} \int_{\partial \mathrm{B}_{1}} x_{i}^{4}=\frac{15}{n(n+2)(n+4)}\left|\partial \mathbf{B}_{1}\right|, \\
\int_{\partial \mathrm{B}_{1}} x_{i}^{2} x_{j}^{2} & =\frac{1}{4} \int_{\mathrm{B}_{1}} \Delta\left(x_{i}^{2} x_{j}^{2}\right)=\frac{2}{4(n+2)} \int_{\partial \mathrm{B}_{1}} 2 x_{i}^{2}=\frac{1}{n(n+2)}\left|\partial \mathbf{B}_{1}\right|, \\
\int_{\partial \mathrm{B}_{1}} x_{i}^{4} x_{j}^{2} & =\frac{1}{6} \int_{\mathbf{B}_{1}} \Delta\left(x_{i}^{4} x_{j}^{2}\right)=\frac{1}{6(n+4)} \int_{\partial \mathrm{B}_{1}}\left(12 x_{i}^{2} x_{j}^{2}+2 x_{i}^{4}\right) \\
& =\frac{3}{n(n+2)(n+4)}\left|\partial \mathbf{B}_{1}\right|,
\end{aligned}
$$

and

$$
\int_{\partial \mathbf{B}_{1}} x_{i}^{2} x_{j}^{2} x_{k}^{2}=\frac{1}{6} \int_{\mathbf{B}_{1}} \Delta\left(x_{i}^{2} x_{j}^{2} x_{k}^{2}\right)=\frac{3}{6(n+4)} \int_{\partial \mathbf{B}_{1}} 2 x_{i}^{2} x_{j}^{2}=\frac{1}{n(n+2)(n+4)}\left|\partial \mathbf{B}_{1}\right| .
$$

Thus, calling $c_{n}:=\frac{\left|\partial \mathrm{B}_{1}\right|}{n(n+2)(n+4)}$ and using that $\sum_{\alpha} a_{\alpha \alpha}=a$ and $\sum_{\alpha} \bar{a}_{\alpha \alpha}=\bar{a}$, we obtain

$$
\int_{\partial \mathbf{B}_{e}} q \bar{q}=\varrho^{n+5}\left(\frac{a \bar{a}}{9} 15 c_{n}-\frac{a \bar{a}}{3} 6 c_{n}+\sum_{\alpha} a_{\alpha \alpha} \bar{a}_{\alpha \alpha} 3 c_{n}+\sum_{\alpha \neq \gamma} a_{\alpha \alpha} \bar{a}_{\gamma \gamma} c_{n}\right) .
$$

Finally, since $\sum_{\alpha} \sum_{\gamma} a_{\alpha \alpha} \bar{a}_{\gamma \gamma}=\left(\sum_{\alpha} a_{\alpha \alpha}\right)\left(\sum_{\gamma} \bar{a}_{\gamma \gamma}\right)=a \bar{a}$ and recalling that $a_{\alpha \beta}$ is diagonal, we get

$$
\int_{\partial \mathrm{B}_{e}} q \bar{q}=2 \varrho^{n+5} c_{n} \sum_{\alpha} a_{\alpha \alpha} \bar{a}_{\alpha \alpha}=2 c_{n} \varrho^{n+5}\left(2 \operatorname{trace}\left(\left(a_{\alpha \beta}\right) \cdot\left(\bar{a}_{\gamma \delta}\right)\right)+\frac{2}{3} a \bar{a}\right),
$$

as claimed.
We can now finally prove the following fundamental result, which implies that $\Sigma_{n-1}^{3 r d} \backslash \Sigma_{n-1}^{>3}=\emptyset:$

Proposition 5.8. - Let $0 \in \Sigma_{n-1}^{3 r d}$, and set $w:=u-p_{2}-p_{3}$. Then $\phi\left(0^{+}, w\right)>3$.
Proof. - Without loss of generality, we can assume that $\left\{p_{2}=0\right\}=\left\{x_{n}=0\right\}$.
Suppose by contradiction that $\phi\left(0^{+}, w\right)=3$. Then we know that the accumulation points of $\tilde{w}_{r}$ as $r \downarrow 0$ must be 3-homogeneous even solutions of the Signorini problem, that is, of the form $q_{\mathrm{A}}$ for some symmetric matrix $\mathrm{A} \geq 0$ (see (5.4)). Note that, by construction, $\left\|q_{\mathrm{A}}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1$ and thus the matrix A must have at least one positive eigenvalue.

Let us define the quantity

$$
\begin{equation*}
\Psi(r):=\max \left\{\psi(r ; \mathrm{A}):\left\|q_{\mathrm{A}}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1\right\} \tag{5.6}
\end{equation*}
$$

where $\psi$ is given by (5.5). Let $\mathrm{A}_{r}^{*}$ be the matrix for which the previous maximum is attained. Then, as a consequence of Lemma 5.6, we have

$$
\frac{d}{d r} \Psi(r)=\theta(r) \Psi(r)-\frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{A}_{r}^{*}}+\mathrm{O}\left(r^{3} / h(r)\right), \quad \text { for a.e. } r>0 .
$$

On the other hand, if we define $\Phi(r):=\psi(r, \mathrm{Id})$, then

$$
\begin{equation*}
\frac{d}{d r} \Phi(r)=\theta(r) \Phi(r)-\frac{1}{r} \int_{\mathbf{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}}+\mathrm{O}\left(r^{3} / h(r)\right) \tag{5.7}
\end{equation*}
$$

We now claim that

$$
\Psi(r) \asymp \Phi(r) \asymp \frac{\Psi(r)}{\Phi(r)} \asymp 1 \quad \text { as } r \downarrow 0
$$

where $\mathrm{X} \asymp \mathrm{Y}$ is a short notation for $\mathrm{X} \leq \mathrm{C}(n) \mathrm{Y}$ and $\mathrm{Y} \leq \mathrm{C}(n) \mathrm{X}$. Indeed, the accumulation points of $\tilde{w}_{r}\left(\right.$ as $r \downarrow 0$ and in the $\mathrm{C}_{\mathrm{loc}}^{0}\left(\mathbf{R}^{n}\right)$ topology) are of the form $q_{\mathrm{A}}$ (and have unit norm) and thus for every $r>0$ we have $w_{r}-q_{\mathrm{A}_{r}}=o(1)$ for some $\mathrm{A}_{r}$. Hence, by definition of $\Psi$,

$$
\begin{aligned}
\Psi(r) \geq \psi\left(r ; \mathrm{A}_{r}\right) & =\int_{\partial \mathrm{B}_{1}} \tilde{w}_{r} q_{\mathrm{A}_{r}}-2 \int_{\partial \mathrm{B}_{1 / 2}} \tilde{w}_{r} q_{\mathrm{A}_{r}}=\int_{\partial \mathrm{B}_{1}} q_{\mathrm{A}_{r}}^{2}-2 \int_{\partial \mathrm{B}_{1 / 2}} q_{\mathrm{A}_{r}}^{2}+o(1) \\
& =\left(1-2^{-n-4}\right) \int_{\partial \mathbf{B}_{1}} q_{\mathrm{A}_{r}}^{2}+o(1) \geq c(n)>0 .
\end{aligned}
$$

Note that the above computation shows also that $\psi\left(r ; \mathrm{A}_{r}^{*}\right)=\left(1-2^{-n-4}\right) \int_{\partial \mathrm{B}_{1}} q_{\mathrm{A}_{r}} q_{\mathrm{A}_{r}^{*}}+o(1)$ (as $r \downarrow 0$ ). Thus since by definition of $\mathrm{A}_{r}^{*}$ we have $\psi\left(r ; \mathrm{A}_{r}^{*}\right) \geq \psi\left(r ; \mathrm{A}_{r}\right)$ it follows

$$
\int_{\partial \mathrm{B}_{1}} q_{\mathrm{A}_{*}^{*}} q_{\mathrm{A}_{r}} \geq \int_{\partial \mathrm{B}_{1}} q_{\mathrm{A}_{r}}^{2}+o(1)
$$

Since $\int_{\partial \mathrm{B}_{1}} q_{\mathrm{A}_{r}^{*}}^{2}=\int_{\partial \mathrm{B}_{1}} q_{\mathrm{A}_{r}}^{2}=1$ is follows that $q_{\mathrm{A}_{t}^{*}}=q_{\mathrm{A}_{t}}+o(1)$ and hence

$$
\mathrm{A}_{r}^{*}=\mathrm{A}_{r}+o(1) \quad \text { as } r \downarrow 0 .
$$

Similarly, using Lemma 5.7,

$$
\begin{aligned}
\Phi(r)= & \int_{\partial \mathbf{B}_{1}} \tilde{w}_{r} q_{\mathrm{Id}}-2 \int_{\partial \mathbf{B}_{1 / 2}} \tilde{w}_{r} q_{\mathrm{Id}}=\int_{\partial \mathbf{B}_{1}} q_{\mathrm{A}_{r}} q_{\mathrm{Id}}-2 \int_{\partial \mathbf{B}_{1 / 2}} q_{\mathrm{A}_{r}} q_{\mathrm{Id}}+o(1) \\
= & \frac{\left(1-2^{-n-4}\right) 4\left|\partial \mathbf{B}_{1}\right|}{n(n+2)(n+4)}\left\{\operatorname{trace}\left(\mathrm{A}_{r}\right)+\frac{1}{3} \operatorname{trace}\left(\mathrm{~A}_{r}\right)(n-1)\right\} \\
& +o(1) \geq c(n)>0 .
\end{aligned}
$$

Since $\Psi(r)$ and $\Phi(r)$ are bounded by above, the claim follows.
Now notice that, using the expressions for $\frac{d}{d r} \Psi$ and $\frac{d}{d r} \Phi$, we find

$$
\begin{aligned}
\frac{d}{d r}\left(\frac{\Psi(r)}{\Phi(r)}\right)= & \frac{1}{r} \frac{-\Phi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}}+\Psi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{A}_{r}^{*}}}{\Phi(r)^{2}} \\
& +\mathrm{O}\left(r^{3} / h(r)\right) .
\end{aligned}
$$

We claim that, given $\varepsilon>0$, for $r$ sufficiently small it holds

$$
\begin{equation*}
\left|\int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{A}_{r}^{*}}\right| \leq\left.\varepsilon\right|_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}} \mid+\mathrm{Cr}^{4} / h(r) . \tag{5.8}
\end{equation*}
$$

Indeed, it follows by Lemma 5.5 that, for any $\eta>0$, if $r>0$ is sufficiently small so that $\left\|\tilde{w}_{r}-q_{\mathrm{A}_{r}^{*}}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{2}\right)} \leq \delta(n, \eta)$ then (here we use the notation $\mathrm{B}_{r}^{\prime}:=\mathrm{B}_{r} \cap\left\{x_{n}=0\right\}$ )

$$
\begin{aligned}
-\int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} w_{r} \Delta q_{\mathrm{A}_{r}^{*}} & =2 \int_{\mathrm{B}^{\prime} \backslash \backslash \mathrm{B}_{1 / 2}^{\prime}} u\left(r x^{\prime}, 0\right)\left(x^{\prime} \cdot \mathrm{A}_{r}^{*} x^{\prime}\right) d x^{\prime} \\
& \leq 2 \eta \int_{\left(\mathrm{B}_{1}^{\prime} \backslash \mathrm{B}_{1 / 2}^{\prime}\right)} \int_{\left\{x^{\prime} \cdot \mathrm{A} x^{\prime} \leq \eta\right\}} u\left(r x^{\prime}, 0\right) d x^{\prime}+\int_{\left(\mathrm{B}_{1}^{\prime} \backslash \mathrm{B}_{1 / 2}^{\prime}\right) \cap\left\{x^{\prime} \cdot \mathrm{A} x^{\prime} \geq \eta\right\}} \mathrm{Cr} d x^{\prime} \\
& \leq 2 \eta \int_{\mathrm{B}^{\prime} \backslash \backslash \mathbf{B}_{1 / 2}^{\prime}} u\left(r x^{\prime}, 0\right)+\mathrm{C} r^{4}
\end{aligned}
$$

(here we used that $w_{r} \equiv u(r \cdot)$ on $\left\{x_{n}=0\right\}$ ), while

$$
-\int_{\mathbf{B}_{1} \backslash \mathbf{B}_{1 / 2}} w_{r} \Delta q_{\mathrm{Id}}=c_{n} \int_{\mathbf{B}_{1}^{\prime} \backslash \mathbf{B}_{1 / 2}^{\prime}} u\left(r x^{\prime}, 0\right)\left|x^{\prime}\right|^{2} d x^{\prime} \geq c_{n} \int_{\mathbf{B}_{1}^{\prime} \backslash \mathbf{B}_{1 / 2}^{\prime}} u\left(r x^{\prime}, 0\right) d x^{\prime},
$$

where $c_{n}>0$. Dividing by $h(r)$, we obtain

$$
0 \leq-\int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{A}_{r}^{*}} \leq-4 \eta \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}}+\mathrm{Cr}^{4} / h(r),
$$

and thus (5.8) follows.
Hence, thanks to (5.8), we have that

$$
\begin{aligned}
\frac{d}{d r}\left(\frac{\Psi(r)}{\Phi(r)}\right)= & \frac{1}{r} \frac{-\Phi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}}+\Psi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{A}_{r}^{*}}}{\Phi(r)^{2}} \\
& +\mathrm{O}\left(r^{3} / h(r)\right) \\
= & -\frac{a(r)}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}}+\mathrm{O}\left(r^{3} / h(r)\right), \quad a(r) \asymp 1 .
\end{aligned}
$$

Choosing $r_{0}$ so that $\mathrm{C}^{-1} \leq a(r) \leq \mathrm{C}$ over $\left[0, r_{0}\right]$, we can integrating the above ODE over $\left[\hat{r}, r_{0}\right]$ for any $\hat{r}>0$. Then, since the integrals of $\frac{d}{d r}\left(\frac{\Psi(r)}{\Phi(r)}\right)$ and $r^{3} / h(r)$ are both uniformly bounded independently of $\hat{r}$, so must be the integral of the negative term $\frac{a(r)}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}}$. Hence, this proves that

$$
\int_{0}^{r_{0}}\left|\frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta q_{\mathrm{Id}}\right| d r<\infty
$$

Since $\Phi(r) \asymp 1$ and $\theta(r)=\frac{d}{d r} \log \left(h(r) / r^{3}\right)$, it follows from (5.7) that

$$
\frac{d}{d r} \log \Phi(r)=\frac{d}{d r} \log \left(h(r) / r^{3}\right)+g(r), \quad \text { with } g \in \mathrm{~L}^{1}\left(\left[0, r_{0}\right]\right)
$$

Integrating over $\left[\hat{r}, r_{0}\right.$ ] and using again that $\Phi(r) \asymp 1$, we deduce that $\log \left(h(\hat{r}) / \hat{r}^{3}\right)$ is uniformly bounded as $\hat{r} \rightarrow 0$, therefore $h(r) \asymp r^{3}$. However, since $0 \in \Sigma_{n-1}^{3 r d}$ we know that $h(r)=o\left(r^{3}\right)$, contradiction.

## 6. Symmetry properties of blow-ups for 1-parameter family of solutions

As explained in the introduction, to establish generic regularity results, we shall consider 1-parameter monotone family of solutions. For this, we shall use the parameter $t$ (over which solutions are indexed) as a second variable for our solution $u$ (one may think of $t$ as a "time" variable, although there is no equation in $t$ ).

So, let $u: \overline{\mathrm{B}_{1}} \times[-1,1] \rightarrow \mathbf{R}, u \geq 0$, be a monotone 1-parameter family of solutions of the obstacle problem, namely

$$
\begin{equation*}
\Delta u(\cdot, t)=\chi_{\{u(\cdot t)>0\}} \quad \text { and } \quad 0 \leq u(\cdot, t) \leq u\left(\cdot, t^{\prime}\right) \quad \text { in } \mathrm{B}_{1}, \quad \text { for }-1 \leq t \leq t^{\prime} \leq 1 \tag{6.1}
\end{equation*}
$$

We will assume in addition that $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ (this continuity property in $t$ follows by the maximum principle whenever $\left.u \in \mathrm{C}^{0}\left(\partial \mathrm{~B}_{1} \times[-1,1]\right)\right)$.

Note that by, the regularity results for the obstacle problem, $u(\cdot, t)$ is of class $\mathrm{C}^{1,1}$ inside $\mathbf{B}_{1}$ for each $t \in(-1,1)$. Moreover, for each fixed $t \in(-1,1)$, we can apply the results of the previous sections, and define the different blow-ups at singular points.

So, following the previous sections, we say that $\left(x_{0}, t_{0}\right)$ is a singular point of $u$ if $x_{0}$ is a singular point of $u\left(\cdot, t_{0}\right)$. Given a singular free boundary point $\left(x_{0}, t_{0}\right)$, we denote

$$
p_{2, x_{0}, t_{0}}(x):=\lim _{r \rightarrow 0} r^{-2} u\left(x_{\circ}+r x, t_{\circ}\right)
$$

Note that $p_{2, x_{0}, t_{0}}$ is a convex 2 -homogeneous polynomials with $\Delta p_{2, x_{0}, t_{0}}=1$. When $\left(x_{0}, t_{\mathrm{o}}\right)=(0,0)$, we simplify the notation to $p_{2}$.

From now on, using the notation introduced in the previous sections, we set:

$$
\begin{aligned}
\boldsymbol{\Sigma} & :=\left\{\left(x_{\circ}, t_{0}\right) \text { singular points in } \mathrm{B}_{1} \times[-1,1]\right\}, \\
\boldsymbol{\Sigma}_{m} & :=\left\{\left(x_{\circ}, t_{0}\right): x_{\circ} \in \Sigma_{m} \text { for } u\left(\cdot, t_{0}\right)\right\}, \quad 0 \leq m \leq n-1, \\
\boldsymbol{\Sigma}_{m}^{a} & :=\left\{\left(x_{\circ}, t_{0}\right): x_{\circ} \in \Sigma_{m}^{a} \text { for } u\left(\cdot, t_{0}\right)\right\}, \quad 0 \leq m \leq n-2, \\
\boldsymbol{\Sigma}_{n-1}^{<3} & :=\left\{\left(x_{0}, t_{\circ}\right): x_{\circ} \in \Sigma_{n-1}^{<3} \text { for } u\left(\cdot, t_{\circ}\right)\right\}, \\
\boldsymbol{\Sigma}_{n-1}^{\geq 3} & :=\left\{\left(x_{\circ}, t_{\circ}\right): x_{\circ} \in \Sigma_{n-1}^{\geq 3} \text { for } u\left(\cdot, t_{\circ}\right)\right\}, \\
\boldsymbol{\Sigma}_{n-1}^{3 r d} & :=\left\{\left(x_{\circ}, t_{\circ}\right): x_{\circ} \in \Sigma_{n-1}^{3 r d} \text { for } u\left(\cdot, t_{\circ}\right)\right\},
\end{aligned}
$$

$$
\begin{aligned}
\boldsymbol{\Sigma}_{n-1}^{>3} & :=\left\{\left(x_{\circ}, t_{0}\right): x_{0} \in \Sigma_{n-1}^{>3} \text { for } u\left(\cdot, t_{0}\right)\right\}, \\
\boldsymbol{\Sigma}_{n-1}^{4 h} & :=\left\{\left(x_{0}, t_{0}\right): x_{0} \in \Sigma_{n-1}^{41 h} \text { for } u\left(\cdot, t_{0}\right)\right\}, \\
\boldsymbol{\Sigma}_{n-1}^{>4} & :=\left\{\left(x_{0}, t_{0}\right): x_{\circ} \in \Sigma_{n-1}^{>4} \text { for } u\left(\cdot, t_{\circ}\right)\right\}, \\
\boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta} & :=\left\{\left(x_{0}, t_{0}\right): x_{0} \in \Sigma_{n-1}^{\geq 5-\zeta} \text { for } u\left(\cdot, t_{0}\right)\right\}, \quad \zeta \in(0,1) .
\end{aligned}
$$

Recall that $\Sigma_{m}, \Sigma_{m}^{a}, \Sigma_{n-1}^{<3}$, and $\Sigma_{m}^{\geq 3}$ were defined in (3.13)-(3.16), while $\Sigma_{m}^{3 r d}, \Sigma_{n-1}^{>3}, \Sigma_{m}^{\geq 4}$, $\Sigma_{n-1}^{4 t h}, \Sigma_{n-1}^{>4}$, and $\Sigma_{n-1}^{\geq 5-\zeta}$ were defined in Definitions 3.11, 4.4, 4.10, 4.13, respectively.

Remark 6.1. - Note that, as a consequence of Proposition 5.8, $\boldsymbol{\Sigma}_{n-1}^{3 r d}=\boldsymbol{\Sigma}_{n-1}^{>3}$.
For $\left(x_{0}, t_{0}\right) \in \boldsymbol{\Sigma}_{m}^{3 r d}$ we define

$$
\begin{equation*}
p_{3, x_{0}, t_{0}}(x):=\lim _{r \rightarrow 0} r^{-3}\left(u\left(x_{\circ}+r x, t_{\circ}\right)-p_{2, x_{0}, t_{0}}(r x)\right), \tag{6.3}
\end{equation*}
$$

and for $\left(x_{\circ}, t_{\circ}\right) \in \boldsymbol{\Sigma}_{m}^{4 t h}$ we define $\mathscr{P}_{x_{0}, t_{0}}$ as the fourth order Ansatz of $u\left(x_{\circ}+\cdot, t_{\circ}\right)$ at 0 (cf. (4.10)), and

$$
\begin{equation*}
p_{4, x_{0}, t_{0}}(x):=\lim _{r \rightarrow 0} r^{-4}\left(u\left(x_{\circ}+r x, t_{0}\right)-\mathscr{P}_{x_{0}, t_{0}}(r x)\right) \tag{6.4}
\end{equation*}
$$

We begin with a simple lemma.
Lemma 6.2. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1). Then:
(a) The singular set is closed-more precisely $\boldsymbol{\Sigma} \cap \overline{\mathrm{B}}_{\varrho} \times[-1,1]$ is closed for any $\varrho<1$. Moreover,

$$
\boldsymbol{\Sigma} \cap \overline{\mathbf{B}}_{\varrho} \times[-1,1] \ni\left(x_{k}, t_{k}\right) \rightarrow\left(x_{\infty}, t_{\infty}\right) \quad \Rightarrow \quad p_{2, x_{k}, t_{k}} \rightarrow p_{2, x_{\infty}, t_{\infty}}
$$

(b) The frequency function

$$
\boldsymbol{\Sigma} \ni\left(x_{\circ}, t_{\mathrm{o}}\right) \mapsto \phi\left(0^{+}, u\left(x_{\circ}+\cdot, t_{\circ}\right)-p_{2, x_{0}, t_{0}}\right)
$$

is upper semi-continuous.
(c) If $\left(x_{0}, t_{1}\right)$ and $\left(x_{0}, t_{2}\right)$ belong both to $\boldsymbol{\Sigma}$ and $t_{1}<t_{2}$, then there exists $r>0$ such that $u(x, t)$ is independent of $t$ for all $(x, t) \in \mathbf{B}_{r}\left(x_{0}\right) \times\left[t_{1}, t_{2}\right]$.

Proof. - (a) We first show that if $\left(x_{k}, t_{k}\right)$ are singular points and $\left(x_{k}, t_{k}\right) \rightarrow\left(x_{\infty}, t_{\infty}\right)$ then the limit point is also singular. Indeed, by Lemma 3.1 we have

$$
\left\|u\left(x_{k}+\cdot, t_{k}\right)-p_{2, x_{k} t_{k}}\right\|_{L^{\infty}\left(\mathbf{B}_{r}\right)} \leq r^{2} \omega(r) \quad \forall r>0 .
$$

Hence, since $u\left(x_{k}+\cdot, t_{k}\right) \rightarrow u\left(x_{\infty}+\cdot, t_{\infty}\right)$ in $\mathrm{C}^{0}$ as $k \rightarrow \infty$ and (after taking a subsequence) $p_{2, x_{k} t_{k}} \rightarrow \mathrm{P}$ for some convex 2-homogeneous polynomials with $\Delta \mathrm{P}=1$, we obtain

$$
\begin{equation*}
\left\|u\left(x_{\infty}+\cdot, t_{\infty}\right)-\mathrm{P}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{r}\right)} \leq r^{2} \omega(r) \quad \forall r>0 . \tag{6.5}
\end{equation*}
$$

Thus $\left(x_{\infty}, t_{\infty}\right) \in \boldsymbol{\Sigma}$ and $p_{2, x_{\infty} t_{\infty}}=\mathrm{P}$. A posteriori, we deduce that for any other subsequence it must be $p_{2, x_{k} t_{k}} \rightarrow p_{2, x_{\infty} t_{\infty}}$ since there is only one P for which (6.5) holds, namely, $p_{2, x_{\infty} t_{\infty}}$.
(b) The upper semicontinuity follows from the fact that the map $r \mapsto \phi\left(r, u\left(x_{\circ}+\right.\right.$ $\left.\left.\cdot, t_{\mathrm{o}}\right)-p_{2, x_{0} t_{0}}\right)$ is increasing, and that for $r>0$ fixed the map $\left(x_{0}, t_{\mathrm{o}}\right) \mapsto \phi\left(r, u\left(x_{0}+\cdot, t_{\mathrm{o}}\right)-\right.$ $\left.p_{2, x_{0} t_{0}}\right)$ is continuous on $\boldsymbol{\Sigma}$-using (a) and the fact that $u\left(x_{0}+\cdot, t_{0}\right)$ satisfies uniform $\mathrm{C}^{1,1}$ estimates.
(c) As in (a), we have, for $i=1,2$,

$$
\left\|u\left(x_{\circ}+\cdot, t_{i}\right)-p_{2, x_{0}, t_{i}}\right\|_{L^{\infty}\left(\mathbf{B}_{r}\right)} \leq r^{2} \omega(r) \quad \forall r>0 .
$$

Since $u\left(x_{\circ}+\cdot, t_{1}\right) \leq u\left(x_{\circ}+\cdot, t_{2}\right)$ then it must be $p_{2, x_{0}, t_{1}} \equiv p_{2, x_{0}, t_{2}}=$ : P. Also, after a change of coordinates, we can assume that $\{\mathrm{P}=0\} \subset\left\{x_{n}=0\right\}$.

Take $r>0$ small enough, and set $v:=u\left(x_{\circ}+r \cdot, t_{2}\right)-u\left(x_{\circ}+r \cdot, t_{1}\right) \geq 0$. Then

$$
\Delta v=0 \quad \text { in }\left\{u\left(x_{\circ}+r \cdot, t_{1}\right)>0\right\} .
$$

Also, as a consequence of Lemma 3.1, given $\varepsilon>0$, for $r>0$ small enough we have

$$
\mathcal{C}_{\varepsilon}:=\left\{y: \operatorname{dist}\left(\frac{y}{|y|},\left\{x_{n}=0\right\}\right)>\varepsilon\right\} \subset\left\{u\left(x_{\circ}+r \cdot, t_{1}\right)>0\right\} .
$$

Consider now the first eigenfunction of

$$
-\Delta_{\mathbf{S}^{n-1}} \Psi=k_{\varepsilon} \Psi \quad \text { in } \mathbf{S}^{n-1} \cap \mathcal{C}_{2 \varepsilon}, \quad \Psi=0 \quad \text { in } \mathbf{S}^{n-1} \cap \partial \mathcal{C}_{2 \varepsilon}
$$

Then, setting $\psi(x):=|x|^{\lambda_{\varepsilon}} \Psi(x /|x|)$ with $k_{\varepsilon}=\left(n-2+\lambda_{\varepsilon}\right) \lambda_{\varepsilon}$, we have that $\psi$ is a positive $\lambda_{\varepsilon}$-homogeneous harmonic function in $\mathcal{C}_{2 \varepsilon}$ which vanishes on the boundary. Note that as $\varepsilon \downarrow 0$ we have $\mathbf{S}^{n-1} \cap \partial \mathcal{C}_{2 \varepsilon} \downarrow\left\{x_{n}=0\right\}$ and $\lambda_{0}=1$ (this corresponds to the solution $\left.\left|x_{n}\right|\right)$. Thus, by continuity, for $\varepsilon>0$ small enough, the function $\hat{\psi}(x):=|x|^{3 / 2} \Psi(x /|x|)$ is subharmonic and vanishes on $\partial \mathcal{C}_{2 \varepsilon}$. Hence using $\hat{\psi}$ as lower barrier and the standard Harnack inequality on $v$, we obtain that if $v>0$ somewhere then $v \geq c \hat{\psi}(x)$ in $\mathrm{B}_{1}$ for some $c>0$. This implies

$$
u\left(x_{\circ}+r \cdot, t_{2}\right) \geq c \hat{\psi}(x),
$$

which is impossible since $u\left(x_{\circ}+\cdot, t_{2}\right)=\mathrm{P}+o\left(|x|^{2}\right)=\mathrm{O}\left(|x|^{2}\right)$, while $\hat{\psi}$ is positive in some cone and 3/2-homogeneous. This proves that $u\left(\cdot, t_{1}\right) \equiv u\left(\cdot, t_{2}\right)$ inside $\mathbf{B}_{r}\left(x_{0}\right)$, which implies the result.

We now prove some relations between $p_{2}$ and singular points close to $(0,0)$.

Lemma 6.3. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathbf{B}_{1}} \times[-1,1]\right)$ solve $(6.1)$, let $\left(x_{k}, t_{k}\right) \in \boldsymbol{\Sigma},(0,0) \in \boldsymbol{\Sigma}$, and assume that $x_{k} \rightarrow 0$. Set $p_{2, k}:=p_{2, x_{k}, t_{k}}$. Then $p_{2, k} \rightarrow p_{2}$ and we have

$$
\begin{aligned}
& \left\|p_{2, k}-p_{2}\left(\frac{x_{k}}{\left|x_{k}\right|}+\cdot\right)\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{1}\right)} \leq \mathrm{C} \omega\left(2\left|x_{k}\right|\right) \text { and } \\
& \left\|p_{2, k}-p_{2}\right\|_{\mathrm{L}^{\infty}\left(\mathbf{B}_{1}\right)} \leq \mathrm{C} \omega\left(2\left|x_{k}\right|\right) .
\end{aligned}
$$

In addition,

$$
\operatorname{dist}\left(\frac{x_{k}}{\left|x_{k}\right|},\left\{p_{2}=0\right\}\right) \rightarrow 0 \quad \text { as } k \rightarrow \infty
$$

Proof. - We observe first that $p_{2, x_{k}, t_{k}} \rightarrow p_{2}$. Indeed, if $t_{k} \rightarrow t_{\infty}$ then (up to a subsequence) by Lemma 6.2 we have $p_{2, x_{k}, t_{k}} \rightarrow p_{2,0, t_{\infty}}$ and $p_{2,0, t_{\infty}} \equiv p_{2}$, as desired.

Now, set $r_{k}:=\left|x_{k}\right|$. By Lemma 3.1 we have

$$
\begin{aligned}
& \left\|r_{k}^{-2} u\left(x_{k}+r_{k} x, t_{k}\right)-p_{2, k}(x)\right\|_{L^{\infty}\left(\mathrm{B}_{2}\right)} \leq 4 \omega\left(2 r_{k}\right) \quad \text { and } \\
& \left\|r_{k}^{-2} u\left(r_{k} x, 0\right)-p_{2}(x)\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{2}\right)} \leq 4 \omega\left(2 r_{k}\right)
\end{aligned}
$$

Thus, defining $y_{k}:=x_{k} /\left|x_{k}\right|$, for all $x \in \mathrm{~B}_{2}$ we have the following: if $t_{k} \leq 0$ then

$$
\begin{aligned}
-4 \omega\left(2 r_{k}\right)+p_{2, k}(x) & \leq r_{k}^{-2} u\left(x_{k}+r_{k} x, t_{k}\right) \leq r_{k}^{-2} u\left(x_{k}+r_{k} x, 0\right) \\
& \leq 4 \omega\left(2 r_{k}\right)+p_{2}\left(y_{k}+x\right),
\end{aligned}
$$

while if $t_{k} \geq 0$ then

$$
\begin{aligned}
4 \omega\left(2 r_{k}\right)+p_{2, k}(x) & \geq r_{k}^{-2} u\left(x_{k}+r_{k} x, t_{k}\right) \geq r_{k}^{-2} u\left(x_{k}+r_{k} x, 0\right) \\
& \geq-4 \omega\left(2 r_{k}\right)+p_{2}\left(y_{k}+x\right)
\end{aligned}
$$

In both cases, since $p_{2, k}$ and $p_{2}$ are nonnegative 2-homogeneous polynomials vanishing at 0 and with Laplacian 1 , then $p_{2, k}-p_{2}\left(y_{k}+\cdot\right)$ is a harmonic quadratic polynomial which vanishes at some point of the segment joining 0 and $y_{k}$, where $y_{k}:=x_{k} /\left|x_{k}\right|$. Moreover, $\left|p_{2, k}-p_{2}\left(y_{k}+\cdot\right)\right|$ is bounded from above by $8 \omega\left(2 r_{k}\right)$ in $\mathbf{B}_{2}$. Using the mean value formula and the fact that all norms are comparable on polynomials, we obtain

$$
\left\|p_{2, k}-p_{2}\left(y_{k}+\cdot\right)\right\|_{\mathrm{L}^{\infty}\left(\mathbf{B}_{1}\right)} \leq \mathrm{C} \|_{p_{2, k}-p_{2}\left(y_{k}+\cdot\right) \|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \leq \mathrm{C} \omega\left(2 r_{k}\right) . . . ~}
$$

By orthogonality of spherical harmonics with different homogeneities (or by a direct computation) we then obtain

$$
\begin{aligned}
\left\|p_{2, k}-p_{2}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}^{2}+\left\|p_{2}-p_{2}\left(y_{k}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}^{2} & =\left\|p_{2, k}-p_{2}\left(y_{k}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}^{2} \\
& \leq \mathrm{C} \omega\left(2 r_{k}\right)^{2}
\end{aligned}
$$

In particular $\left\|p_{2}-p_{2}\left(y_{k}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \rightarrow 0$, and therefore $\operatorname{dist}\left(y_{k},\left\{p_{2}=0\right\}\right) \rightarrow 0$.

We prove next two key lemmas that will allow us to perform some dimension reduction arguments needed to control the spatial projection (i.e., $\left.\pi_{1}:(x, t) \mapsto x\right)$ of some "bad" subsets of $\boldsymbol{\Sigma} \subset \mathrm{B}_{1} \times[-1,1]$. Note that the spatial version of these first two lemmas (i.e., when considering $u\left(\cdot, t_{\mathrm{o}}\right)$ with $t_{\mathrm{o}}$ fixed) was proven in [FS19]. Here we need stronger results valid for a one-parameter monotone family of solutions to the obstacle problem. To our best knowledge, this is the first dimension reduction argument applicable to a one-parameter family of solutions to an elliptic equation, and it will involve several new and non-standard techniques.

We recall that, given $w: \mathbf{R}^{n} \rightarrow \mathbf{R}$, the rescaled functions $w_{r}$ and $\tilde{w}_{r}$ have been defined in (2.1).

The first lemma concerns the intermediate strata of the singular set $\boldsymbol{\Sigma}_{m}$ with $0 \leq$ $m \leq n-2$.

Lemma 6.4. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathbf{B}_{1}} \times[-1,1]\right)$ solve $(6.1)$, let $(0,0) \in \boldsymbol{\Sigma}_{m}$ with $0 \leq m \leq n-2$, and assume that $u(\cdot, 0) \not \equiv p_{2}$. Let $\left(x_{k}, t_{k}\right) \in \boldsymbol{\Sigma}_{m}$ satisfy $\left|x_{k}\right| \leq r_{k}$ with $r_{k} \downarrow 0$, and suppose that
(6.6)

$$
\tilde{w}_{r_{k}} \rightharpoonup q \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right) \quad \text { for } w:=u-p_{2} \quad \text { and } \quad y_{k}:=\frac{x_{k}}{r_{k}} \rightarrow y_{\infty} .
$$

Then $y_{\infty} \in\left\{p_{2}=0\right\}$ and $q\left(y_{\infty}\right)=0$.
Proof. - Let us define

$$
w_{k}:=u\left(x_{k}+r_{k} \cdot, t_{k}\right)-p_{2}\left(r_{k} \cdot\right)=w_{k}^{(1)}+w_{k}^{(2)}+w_{k}^{(3)},
$$

where

$$
\begin{aligned}
& w_{k}^{(1)}:=u\left(x_{k}+r_{k} \cdot, t_{k}\right)-u\left(x_{k}+r_{k} \cdot, 0\right), \\
& w_{k}^{(2)}:=u\left(x_{k}+r_{k} \cdot, 0\right)-p_{2}\left(x_{k}+r_{k} \cdot\right), \\
& w_{k}^{(3)}:=p_{2}\left(x_{k}+r_{k} \cdot\right)-p_{2}\left(r_{k} \cdot\right) .
\end{aligned}
$$

We divide the proof into three steps.

- Step 1. We prove that

$$
\tilde{w}_{k}:=\frac{w_{k}}{\left\|w_{k}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightharpoonup \mathrm{Q} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)
$$

for some harmonic function Q with polynomial growth.
Indeed, since $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$, by the monotonicity of $\phi$ there exist $r_{\circ}>0$ and $k_{\circ} \in \mathbf{N}$ such that, for $\mathrm{M}:=\phi\left(0^{+}, u(\cdot, 0)-p_{2}\right)+1$, we have

$$
\begin{equation*}
\phi\left(r, u\left(x_{k}+\cdot, t_{k}\right)-p_{2}\right) \leq \mathrm{M} \quad \forall r \in\left(0, r_{\circ}\right), \forall k \geq k_{\circ}, \tag{6.7}
\end{equation*}
$$

or equivalently

$$
\begin{equation*}
\phi\left(r, w_{k}\right) \leq \mathrm{M} \quad \forall r \in\left(0, r_{\circ} / r_{k}\right), \forall k \geq k_{\circ} . \tag{6.8}
\end{equation*}
$$

Then, applying Lemma 3.6 to $w_{k}$, we obtain the following polynomial growth control for $\tilde{w}_{k}:$

$$
\begin{equation*}
\mathrm{H}\left(\mathrm{R}, \tilde{w}_{k}\right) \leq \mathrm{CR}^{2 \mathrm{M}+1} \mathrm{H}\left(1, \tilde{w}_{k}\right)=\mathrm{CR}^{2 \mathrm{M}+1} \quad \forall \mathrm{R} \in\left[1, r_{\circ} / r_{k}\right), \forall k \geq k_{\circ} \tag{6.9}
\end{equation*}
$$

Note that (6.8) is equivalent to $\phi\left(r_{k}, \tilde{w}_{k}\right) \leq \mathrm{M}$, which combined with (6.9) implies that

$$
\begin{equation*}
\left\|\tilde{w}_{k}\right\|_{\mathrm{W}^{1,2}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R}) . \tag{6.10}
\end{equation*}
$$

This gives compactness of the sequence $\tilde{w}_{k}$ and hence (up to a subsequence)

$$
\tilde{w}_{k} \rightharpoonup \mathrm{Q} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)
$$

for some $\mathrm{Q}: \mathbf{R}^{n} \rightarrow \mathbf{R}$. Let us prove next that Q is harmonic.
Indeed, on the one hand we have

$$
\begin{equation*}
\Delta w_{k}=-r_{k}^{2} \chi_{\left\{u\left(x_{k}+r_{k} \cdot, t_{k}\right)=0\right\}} \leq 0 \quad \text { in } \mathrm{B}_{\frac{1}{2 r_{k}}} . \tag{6.11}
\end{equation*}
$$

On the other hand, Lemmas 3.1 and 6.3 imply that, for $\mathrm{R} \geq 1$,

$$
\begin{aligned}
x \in \mathrm{~B}_{\mathrm{R}} \cap\left\{u\left(x_{k}+r_{k} x, t_{k}\right)=0\right\} & \Rightarrow p_{2, x_{k}, t_{k}}(x) \leq \mathrm{R}^{2} \omega\left(\mathrm{R} r_{k}\right) \\
& \Rightarrow p_{2}(x) \leq \mathrm{CR}^{2} \omega\left(\mathrm{R} r_{k}\right) ;
\end{aligned}
$$

thus, since $p_{2}$ grows quadratically away from its zero set,

$$
\begin{equation*}
\mathrm{B}_{\mathrm{R}} \cap\left\{u\left(x_{k}+r_{k} \cdot, t_{k}\right)=0\right\} \subset\left\{y: \operatorname{dist}\left(y,\left\{p_{2}=0\right\}\right) \leq \mathrm{CR} \sqrt{\omega\left(\mathrm{R} r_{k}\right)}\right\} . \tag{6.12}
\end{equation*}
$$

Note that, for any fixed $\mathrm{R} \geq 1$, we have $\mathrm{CR} \sqrt{\omega\left(\mathrm{R} r_{k}\right)} \downarrow 0$ as $k \rightarrow \infty$. We have thus shown

$$
\sup \left\{\operatorname{dist}\left(x,\left\{p_{2}=0\right\}\right): x \in \mathrm{~B}_{\mathrm{R}} \cap\left\{u\left(x_{k}+r_{k} \cdot, t_{k}\right)=0\right\}\right\} \downarrow 0 \quad \text { as } k \rightarrow \infty .
$$

Therefore, the weak limit of the sequence of nonpositive measures $\Delta \tilde{w}_{k}$ will be supported on $\left\{p_{2}=0\right\}$. But then, recalling (6.10), we have shown that $Q$ is a locally $W^{1,2}$ function whose Laplacian is supported in linear space of dimension $m=\operatorname{dim}\left(\left\{p_{2}=0\right\}\right) \leq n-2$ and thus of zero harmonic capacity. It follows ${ }^{5}$ that $\mathbf{Q}$ must be harmonic.

Moreover, since $x_{k}$ is a singular point, Lemma 3.6 yields

$$
\mathrm{H}\left(\rho, w_{k}\right) \leq \rho^{4} \mathrm{H}\left(1, w_{k}\right) \quad \text { for all } \rho \in(0,1)
$$

and thus in the limit we find

$$
\begin{equation*}
\mathrm{H}(\rho, \mathrm{Q})^{1 / 2} \leq \rho^{2} \quad \text { for all } \rho \in(0,1) \tag{6.13}
\end{equation*}
$$

[^5]- Step 2. We now want to prove that
(6.14)

$$
\frac{w_{k}^{(2)}}{\left\|w_{k}^{(2)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightharpoonup \frac{q\left(y_{\infty}+\cdot\right)}{\left\|q\left(y_{\infty}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)
$$

(with $q$ defined in (6.6)), and

$$
\begin{equation*}
\frac{w_{k}^{(3)}}{\left\|w_{k}^{(3)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightarrow \nabla p_{2} \cdot \boldsymbol{e} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right) \tag{6.15}
\end{equation*}
$$

for some (nonzero) $\boldsymbol{e} \in\left\{p_{2}=0\right\}^{\perp}$.
Note that, since $y_{\infty} \in\left\{p_{2}=0\right\}$ (by Lemma 6.3),

$$
\frac{w_{k}^{(2)}}{\left\|w_{k}^{(2)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}}=\frac{w_{r_{k}}\left(y_{k}+\cdot\right)}{\left\|w_{r_{k}}\left(y_{k}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}}=\tilde{w}_{r_{k}}\left(y_{k}+\cdot\right) \frac{\left\|\tilde{w}_{r_{k}}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}}{\left\|\tilde{w}_{r_{k}}\left(y_{k}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} .
$$

Thus, noticing that $\left\|\tilde{w}_{r_{k}}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \rightarrow\|q\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}$ and $\left\|\tilde{w}_{r_{k}}\left(y_{k}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \rightarrow\left\|q\left(y_{\infty}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}$, since $q$ is a nonzero quadratic harmonic polynomial (see Proposition 3.9) both limits are nonzero and universally bounded. Thus (6.14) follows.

To prove (6.15), we set $\varepsilon_{k}:=\left\|p_{2}\left(y_{k}+\cdot\right)-p_{2}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \rightarrow 0$. Then, if $y_{k}^{*}$ denotes the projection of $y_{k}$ onto $\left\{p_{2}=0\right\}$, we have $p_{2}\left(y_{k}^{*}+\cdot\right) \equiv p_{2}$ and $y_{k}^{*}-y_{k} \rightarrow y_{\infty}^{*}-y_{\infty}=0$. Thus, up to taking a further subsequence, we obtain

$$
\begin{aligned}
\lim _{k} \frac{w_{k}^{(3)}}{\left\|w_{k}^{(3)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} & =\lim _{k} \frac{p_{2}\left(y_{k}+\cdot\right)-p_{2}}{\varepsilon_{k}}=\lim _{k} \frac{p_{2}\left(y_{k}-y_{k}^{*}+\cdot\right)-p_{2}}{\varepsilon_{k}} \\
& =c \nabla p_{2} \cdot \lim _{k} \frac{y_{k}-y_{k}^{*}}{\left|y_{k}-y_{k}^{*}\right|}=\nabla p_{2} \cdot \boldsymbol{e}
\end{aligned}
$$

for some nonzero $\boldsymbol{e} \in\left\{p_{2}=0\right\}^{\perp}$. Note that the limit in $k$ exists (up to subsequence) and is nonzero, since $\frac{w_{k}^{(3)}}{\left\|w_{k}^{3}\right\|_{L^{2}\left(\partial B_{1}\right)}}$ is a sequence of linear functions with unit $\mathrm{L}^{2}$ norm.

- Step 3. We finally prove that $q\left(y_{\infty}\right)=0$.

Let us consider

$$
\hat{\varepsilon}_{k}:=\sum_{i=1,2,3}\left\|w_{k}^{(i)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \quad \text { and } \quad \hat{w}_{k}:=\frac{w_{k}}{\hat{\varepsilon}_{k}} .
$$

By Step 1 we have
(6.16)

$$
\hat{w}_{k} \rightarrow \hat{\mathbf{Q}}=a \mathbf{Q} \quad \text { for some } a \in[0,1]
$$

Moreover, by Step 2,

$$
\hat{\mathbf{Q}}^{(2)}:=\lim _{k} w_{k}^{(2)} / \hat{\varepsilon}_{k}=b q\left(y_{\infty}+\cdot\right), \quad \hat{\mathbf{Q}}^{(3)}:=\lim _{k} w_{k}^{(3)} / \hat{\varepsilon}_{k}=c \nabla p_{2} \cdot \boldsymbol{e},
$$

for some constant $b, c \geq 0$. (Above, the convergences are weak in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$.)

Then, it is well defined

$$
\hat{Q}^{(1)}:=\lim _{k} w_{k}^{(1)} / \hat{\varepsilon}_{k}=\lim _{k} w_{k} / \hat{\varepsilon}_{k}-\lim _{k} w_{k}^{(2)} / \hat{\varepsilon}_{k}-\lim _{k} w_{k}^{(3)} / \hat{\varepsilon}_{k},
$$

and we observe that $\mathbf{Q}^{(1)}$ is either nonpositive or nonnegative (since $w_{k}^{(1)}=u\left(x_{k}+r_{k} \cdot, t_{k}\right)$ $u\left(x_{k}+r_{k} \cdot, 0\right)$ is so, depending on the sign of $\left.t_{k}\right)$. Moreover, since $\hat{\mathbf{Q}} \hat{\mathbf{Q}}^{(2)}$, and $\hat{\mathbf{Q}}^{(3)}$ are harmonic, so is $\hat{\mathbf{Q}}^{(1)}$ and thus it must be constant by Liouville Theorem. Hence, we have

$$
\hat{\mathbf{Q}}=\mathrm{C}+b q\left(y_{\infty}+\cdot\right)+c \nabla p_{2} \cdot \boldsymbol{e} .
$$

Note now that, by definition of $\hat{\varepsilon}_{k}$, we have $\sum_{i=1,2,3}\left\|\hat{\mathrm{Q}}^{(i)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1$. Moreover, since the homogeneity of $q$ at the origin is at least two, the three functions $\hat{Q}^{(i)}$ are linearly independent and hence their sum $\hat{Q}$ cannot be zero (equivalently, in (6.16) it must be $a>0)$. Note also that it must be $b>0$ since (6.13) implies that $\hat{Q}$ is at least quadratic and hence it can not be equal to the constant $\hat{\mathbf{Q}}^{(1)}$ plus the linear function, $\hat{\mathbf{Q}}^{(3)}$. Finally, (6.13) implies $\nabla \mathrm{Q}(0)=0$.

But then, since $y_{\infty} \in\left\{p_{2}=0\right\}$, and $q$ is a homogeneous polynomial of degree $\phi(q, 1)$,

$$
\begin{aligned}
0 & =y_{\infty} \cdot \nabla \hat{Q}(0)=y_{\infty} \cdot \nabla \hat{\mathbf{Q}}^{(1)}(0)+b y_{\infty} \cdot \nabla q\left(y_{\infty}\right)+c y_{\infty} \cdot \nabla\left(\nabla p_{2} \cdot \boldsymbol{e}\right)(0) \\
& =0+b \phi(q, 1) q\left(y_{\infty}\right)+0
\end{aligned}
$$

which proves that $q\left(y_{\infty}\right)=0$.
The next lemma concerns the maximal stratum $\boldsymbol{\Sigma}_{n-1}$. This case is more involved, since blow-ups are not necessarily harmonic functions as in the previous lemma. In particular, in this situation we will need to assume that the frequency is continuous along the sequence that we consider.

Lemma 6.5. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathbf{B}_{1}} \times[-1,1]\right)$ solve (6.1), let $(0,0) \in \boldsymbol{\Sigma}_{n-1}$, and assume that $u(\cdot, 0) \not \equiv p_{2}$. Let $\left(x_{k}, t_{k}\right) \in \boldsymbol{\Sigma}_{n-1}$ satisfy $\left|x_{k}\right| \leq r_{k}$ with $r_{k} \downarrow 0$, assume that (6.6) holds, and that $\lambda_{k}^{2 n d} \rightarrow \lambda^{2 n d}$, where

$$
\lambda_{k}^{2 n d}:=\phi\left(0^{+}, u\left(x_{k}+\cdot, t_{k}\right)-p_{2, x_{k}, t_{k}}\right) \quad \text { and } \quad \lambda^{2 n d}:=\phi\left(0^{+}, u-p_{2}\right) .
$$

Then $y_{\infty} \in\left\{p_{2}=0\right\}$ and $q^{\text {even }}$ is translation invariant in the direction $y_{\infty}$. (Here $q^{\text {even }}$ denotes the even symmetrisation of $q$ with respect to the hyperplane $\left\{p_{2}=0\right\}$.)

Proof. - Let us define

$$
w_{k}:=u\left(x_{k}+r_{k} \cdot, t_{k}\right)-p_{2, x_{k}, t_{k}}\left(r_{k} \cdot\right)=w_{k}^{(1)}+w_{k}^{(2)}+w_{k}^{(3)},
$$

where

$$
\begin{aligned}
& w_{k}^{(1)}:=u\left(x_{k}+r_{k} \cdot, t_{k}\right)-u\left(x_{k}+r_{k} \cdot, 0\right), \\
& w_{k}^{(2)}:=u\left(x_{k}+r_{k} \cdot, 0\right)-p_{2}\left(x_{k}+r_{k} \cdot\right), \\
& w_{k}^{(3)}:=p_{2}\left(x_{k}+r_{k} \cdot\right)-p_{2, x_{k}, t_{k}}\left(r_{k} \cdot\right) .
\end{aligned}
$$

We divide the proof into three steps.

- Step 1. Exactly as in Lemma 6.4,

$$
\tilde{w}_{k}:=\frac{w_{k}}{\left\|w_{k}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightharpoonup \mathrm{Q} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)
$$

for some $\mathrm{Q} \in \mathrm{W}_{\mathrm{loc}}^{1.2}\left(\mathbf{R}^{n}\right)$ with polynomial growth. We claim that Q is a $\lambda^{2 n d}$-homogeneous solution of the Signorini problem (3.12).

Indeed, by the upper-semicontinuity property in Lemma 6.2(b) and the assumption $\lambda_{k}^{2 n d} \rightarrow \lambda^{2 n d}$, given $\delta>0$ there exist $r_{\delta}>0$ and $k_{\delta}$ such that

$$
\begin{equation*}
\phi\left(r, u\left(x_{k}+\cdot, t_{k}\right)-p_{2, x_{k}, t_{k}}\right) \in\left(\lambda^{2 n d}-\delta, \lambda^{2 n d}+\delta\right) \quad \forall r \in\left(0, r_{\delta}\right), \forall k \geq k_{\delta}, \tag{6.17}
\end{equation*}
$$

or equivalently
(6.18)

$$
\phi\left(r, w_{k}\right) \in\left(\lambda^{2 n d}-\delta, \lambda^{2 n d}+\delta\right) \quad \forall r \in\left(0, r_{\delta} / r_{k}\right), \forall k \geq k_{\delta} .
$$

Then, applying Lemma 3.6 to $w_{k}$ we obtain the following polynomial growth control for $\tilde{w}_{k}:$

$$
\begin{equation*}
\mathrm{H}\left(\mathrm{R}, \tilde{w}_{k}\right) \leq \mathrm{C}_{\delta} \mathrm{R}^{2 \lambda^{2 n d}+3 \delta} \quad \forall \mathrm{R} \in\left[1, r_{\delta} / r_{k}\right), \quad \forall k \geq k_{0} \tag{6.19}
\end{equation*}
$$

and the decay estimate

$$
\begin{equation*}
\mathrm{H}\left(\varrho, \tilde{w}_{k}\right) \leq \mathrm{C} \varrho^{2\left(\lambda^{2 n d}-\delta\right)} \quad \forall \varrho \in(0,1], \forall k \geq k_{\circ} \tag{6.20}
\end{equation*}
$$

In addition, the Lipschitz estimate in Lemma 3.8 gives

$$
\left\|\tilde{w}_{k}\right\|_{\mathrm{Lip}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R})
$$

Hence $\tilde{w}_{k} \rightarrow \mathrm{Q}$ in $\mathrm{C}_{\mathrm{loc}}^{0}\left(\mathbf{R}^{n}\right)$ (up to a further subsequence).
Note that, using (6.11) and (6.12) in our context, one deduces that $\Delta \mathrm{Q}$ is a nonpositive measure supported on $\left\{p_{2}=0\right\}$. Moreover, since $w_{k}\left(y_{k}+\cdot\right)=u\left(x_{k}+r_{k} \cdot\right)-p_{2, x_{k} k}\left(r_{k} \cdot\right)$, it follows that $\tilde{w}_{k}\left(y_{k}+\cdot\right) \geq 0$ on $\left\{p_{2, x_{k} k_{k}}=0\right\}$ and thus, by uniform convergence, $\mathbf{Q} \geq 0$ on $\left\{p_{2}=0\right\}$.

On the other hand (6.11) and the fact that $\tilde{w}_{k}\left(y_{k}+\cdot\right) \leq 0$ on $\left\{u\left(x_{k}+r_{k} \cdot, t_{k}\right)=0\right\}$ imply that $\tilde{w}_{k} \Delta \tilde{w}_{k} \geq 0$, and since $\Delta \tilde{w}_{k} \rightharpoonup \Delta \mathrm{Q}$ weakly as measures and $\tilde{w}_{k} \rightarrow \mathrm{Q}$ in $\mathrm{C}^{0}$, we obtain $Q \Delta Q \geq 0$ in $\mathbf{R}^{n}$. But since $\Delta \mathrm{Q} \leq$ is nonpositive and supported on $\left\{p_{2}=0\right\}$
where $Q \geq 0$, it must be $Q \Delta Q \leq 0$. This implies that $Q$ is a solution of the Signorini problem (3.12).

Finally, taking the limit in (6.19) and (6.20) we obtain that, for any given $\delta>0$,
(6.21)

$$
\mathrm{H}(\mathrm{R}, \mathrm{Q}) \leq \mathrm{C}_{\delta} \mathrm{R}^{2 \lambda^{2 n d}+3 \delta} \quad \forall \mathrm{R} \in[1, \infty)
$$

and

$$
\begin{equation*}
\mathrm{H}(\varrho, \mathrm{Q}) \leq \mathrm{C}^{2\left(\lambda^{2 n d}-\delta\right)} \quad \forall \varrho \in(0,1] \tag{6.22}
\end{equation*}
$$

Since $\delta>0$ is arbitrary and Q is a global solution of Signorini, it follows by Lemma A. 3 that

$$
\lambda^{2 n d} \leq \phi\left(0^{+}, \mathrm{Q}\right) \leq \phi(+\infty, \mathrm{Q}) \leq \lambda^{2 n d}
$$

Hence $\phi(r, \mathrm{D})=\lambda^{2 n d}$ for all $r>0$, from which (using Lemma A. 3 again) it follows that Q is a $\lambda^{2 n d}$-homogeneous.

- Step 2. We now want to prove that

$$
\begin{equation*}
\frac{w_{k}^{(2)}}{\left\|w_{k}^{(2)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightharpoonup \frac{q\left(y_{\infty}+\cdot\right)}{\left\|q\left(y_{\infty}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right) \tag{6.23}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{k} \frac{w_{k}^{(3)}}{\left\|w_{k}^{3}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightarrow(\boldsymbol{e} \cdot x)+\left(\boldsymbol{e}^{\prime} \cdot x\right)(\boldsymbol{e} \cdot x) \not \equiv 0 \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right) \tag{6.24}
\end{equation*}
$$

for some $\boldsymbol{e} \in\left\{p_{2}=0\right\}^{\perp}$ and $\boldsymbol{e}^{\prime} \in\left\{p_{2}=0\right\}$.
Indeed, the proof of (6.23) is identical to the one of (6.14) in the proof of Lemma 6.4.

To show (6.24), denote $\varepsilon_{k}:=\left\|p_{2}\left(y_{k}+\cdot\right)-p_{2, x_{k}, t_{k}}\right\| \rightarrow 0$. Recall that (by Lemma 6.3) we have $y_{\infty} \in\left\{p_{2}=0\right\}$ and hence, if $y_{k}^{*}$ denotes the projection of $y_{k}$ onto $\left\{p_{2}=0\right\}$, then $p_{2}\left(y_{k}^{*}+\cdot\right) \equiv p_{2}$ and $y_{k}^{*}-y_{k} \rightarrow y_{\infty}^{*}-y_{\infty}=0$. Thus, up to taking a further subsequence, if $\left\{p_{2}=0\right\}=\{\hat{\boldsymbol{e}} \cdot x=0\}$ and $\left\{p_{2, x_{k}, t_{k}}=0\right\}=\left\{\hat{\boldsymbol{e}}_{k} \cdot x=0\right\}$ with $\hat{\boldsymbol{e}}, \hat{\boldsymbol{e}}_{k} \in \mathbf{S}^{n-1}$, then

$$
\begin{aligned}
\lim _{k} \frac{w_{k}^{(3)}}{\left\|w_{k}^{(3)}\right\|_{L^{2}\left(\partial \mathrm{~B}_{1}\right)}}= & \lim _{k} \frac{p_{2}\left(y_{k}+\cdot\right)-p_{2, x_{k}, t_{k}}}{\varepsilon_{k}}=\lim _{k} \frac{p_{2}\left(y_{k}-y_{k}^{*}+\cdot\right)-p_{2}}{\varepsilon_{k}} \\
& +\lim _{k} \frac{p_{2}-p_{2, x_{k}, t_{k}}}{\varepsilon_{k}} \\
= & c_{1} \nabla p_{2} \cdot \lim _{k} \frac{y_{k}-y_{k}^{*}}{y_{k}-y_{k}^{*} \mid}+c_{2} \lim _{k} \frac{(\hat{\boldsymbol{e}} \cdot x)^{2}-\left(\hat{\boldsymbol{e}}_{k} \cdot x\right)^{2}}{2\left|\hat{\boldsymbol{e}}-\hat{\boldsymbol{e}}_{k}\right|} \\
= & (\boldsymbol{e} \cdot x)+\left(\boldsymbol{e}^{\prime} \cdot x\right)(\boldsymbol{e} \cdot x),
\end{aligned}
$$

where $\boldsymbol{e} \in\left\{p_{2}=0\right\}^{\perp}$ and $\boldsymbol{e}^{\prime} \in\left\{p_{2}=0\right\}$. Note that the previous limit in $k$ must exist (up to subsequence) and will be nonzero, since $w_{k}^{(3)} /\left\|w_{k}^{(3)}\right\|_{L^{2}\left(\partial \mathrm{~B}_{1}\right)}$ is a sequence of quadratic polynomials with unit $\mathrm{L}^{2}$ norm.

- Step 3. We finally prove that $q$ is translation invariant in the direction $y_{\infty}$. Consider

$$
\hat{\varepsilon}_{k}:=\sum_{i=1,2,3}\left\|w_{k}^{(i)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \quad \text { and } \quad \hat{w}_{k}:=\frac{w_{k}}{\hat{\varepsilon}_{k}} .
$$

By Step 1 we have

$$
\hat{w}_{k} \rightarrow \hat{\mathbf{Q}}=a \mathbf{Q} \quad \text { for some } a \in[0,1] .
$$

Moreover, by Step 2

$$
\hat{\mathbf{Q}}^{(2)}:=\lim _{k} w_{k}^{(2)} / \hat{\varepsilon}_{k}=b q\left(y_{\infty}+\cdot\right)
$$

and, after choosing some appropriate coordinate frame (so that, in particular, $\left\{p_{2}=0\right\}=$ $\left\{x_{n}=0\right\}$ ),

$$
\hat{\mathbf{Q}}^{(3)}:=\lim _{k} w_{k}^{(3)} / \hat{\varepsilon}_{k}=c_{1} x_{n}+c_{2} x_{n} x_{n-1}
$$

for some $b, c \geq 0$. (Above, the convergences are weak in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$.)
Then, it is well defined

$$
\hat{\mathbf{Q}}^{(1)}:=\lim _{k} w_{k}^{(1)} / \hat{\varepsilon}_{k}=\lim _{k} w_{k} / \hat{\varepsilon}_{k}-\lim _{k} w_{k}^{(2)} / \hat{\varepsilon}_{k}-\lim _{k} w_{k}^{(3)} / \hat{\varepsilon}_{k},
$$

and we observe that $Q^{(1)}$ is either nonpositive or nonnegative (since the functions $w_{k}^{(1)}$ are so). Hence, we have

$$
\hat{\mathbf{Q}}=\hat{\mathbf{Q}}^{(1)}+b q\left(y_{\infty}+\cdot\right)+c_{1} x_{n}+c_{2} x_{n} x_{n-1} .
$$

Note now that, by definition of $\hat{\varepsilon}_{k}$, we have $\sum_{i=1,2,3}\left\|\hat{\mathbf{Q}}^{(i)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1$. Moreover, since the homogeneity of $q$ at the origin is at least $2+\alpha_{\circ}$ (see Proposition 3.9), the three functions $\hat{\mathrm{Q}}^{(i)}$ are linearly independent ${ }^{6}$ and thus their sum $\hat{\mathrm{Q}}$ cannot be zero.

Let us show next that $b>0$ and that $\hat{\mathbb{Q}} \equiv b q$. Indeed, since both $q$ and $\hat{\mathbb{Q}}$ are $\lambda^{2 n d}$-homogeneous with $\lambda^{2 n d} \geq 2+\alpha_{o}$, if $\mathbf{Q}^{(1)} \geq 0$ (resp. $\leq$ ) then

$$
\begin{aligned}
\hat{\mathrm{Q}} & =\lim _{\mathrm{R} \rightarrow \infty} \frac{\hat{\mathrm{Q}}(\mathrm{R} \cdot)}{\mathrm{R}^{\lambda^{2 n d}}} \\
& =\lim _{\mathrm{R} \rightarrow \infty} \frac{\mathrm{Q}^{(1)}(\mathrm{R} \cdot)+b q\left(y_{\infty}+\mathrm{R} \cdot\right)+\mathrm{Q}^{(3)}(\mathrm{R} \cdot)}{\mathrm{R}^{\lambda^{2 n d}}} \geq b q \quad(\text { resp. } \leq),
\end{aligned}
$$

[^6]where we used that $\mathbf{Q}^{(3)}$ is 2-homogeneous. Hence, $\hat{\mathbf{Q}}$ and $b q$ are two ordered solutions of Signorini with homogeneities greater than 1 at the origin and thus they must be equal by Lemma A. 4 .

Therefore, we have shown that

$$
\begin{equation*}
\hat{\mathbf{Q}}=\hat{\mathbf{Q}}^{(1)}+b q\left(y_{\infty}+\cdot\right)+x_{n}\left(c_{1} x_{n-1}+c_{2}\right)=b q . \tag{6.25}
\end{equation*}
$$

In particular, since $\hat{\mathbf{Q}}$ has unit $\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)$ norm this implies that $b>0$.
Now, taking the even parts, if $\mathbf{Q}^{(1)} \geq 0\left(\right.$ resp. if $\left.Q^{(1)} \leq 0\right)$ we obtain

$$
\begin{equation*}
b q^{\text {even }}\left(y_{\infty}+\cdot\right) \leq b q^{\text {even }} \quad(\text { resp. } \geq) . \tag{6.26}
\end{equation*}
$$

Hence it follows by homogeneity that, for all $s>0$,

$$
b s^{-\lambda^{2 n d} d} q^{\text {even }}\left(s y_{\infty}+x\right) \leq b s^{-\lambda^{2 n d}} q^{\text {even }}(x) \quad(\text { resp. } \geq) .
$$

Therefore, since $b>0$,

$$
q\left(s y_{\infty}+x\right) \leq q(x) \quad(\text { resp. } \geq),
$$

and thus

$$
y_{\infty} \cdot \nabla q^{\text {even }} \leq 0 \quad(\text { resp. } \geq) .
$$

In summary we obtain that $\psi:=y_{\infty} \cdot \nabla q^{\text {even }}$ has constant sign. But then $\psi$ restricted to the sphere $\mathbf{S}^{n-1}$ must be a multiple of the first even eigenfunction (since all other eigenfunctions change sign) of

$$
\begin{cases}-\Delta_{\mathbf{s}^{n-1}} \psi=k \psi & \text { in } \mathbf{S}^{n-1} \backslash \mathcal{Z} \\ \psi=0 & \text { on } \mathbf{S}^{n-1} \cap \mathcal{Z}\end{cases}
$$

where $\mathcal{Z}:=\left\{x_{n}=0\right\} \cap\{q=0\}$ and $k:=\left(n-2+\lambda^{2 n d}\right) \lambda^{2 n d}$. Note $\mathcal{Z} \subset\left\{x_{n}=0\right\}$, and the two extremal cases $\mathcal{Z}=\emptyset$ and $\mathcal{Z}=\left\{x_{n}=0\right\}$ correspond respectively to the eigenfunctions 1 and $\left|x_{n}\right|$ (restricted to the sphere), which have homogeneity 0 and 1 respectively. As a consequence of the monotonicity property of the eigenvalues with respect to the domain, for every $\mathcal{Z}$ we will have $(n-2+0) 0 \leq k=\left(n-2+\lambda^{2 n d}\right) \lambda^{2 n d} \leq(n-2+1) 1$. This leads to $\lambda^{2 n d} \leq 2$; a contradiction. Therefore, the only possibility is that $\psi=y_{\infty} \cdot \nabla q^{e v e n} \equiv 0$. In other words $q^{\text {even }}$ is translation invariant in the direction $y_{\infty}$.

The next result will imply that the projection $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}\right)$ (recall that $\left.\pi_{1}(x, t)=x\right)$ is contained in a countable union of $(n-2)$-dimensional Lipschitz manifolds, i.e., it is $(n-2)$-rectifiable. This will be crucial in our proof of Theorem 1.1.

Lemma 6.6. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathbf{B}_{1}} \times[-1,1]\right)$ solve (6.1), and let $(0,0) \in \boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}$. Then there exists a $(n-2)$-dimensional linear subspace L such that the following holds: for any $\varepsilon>0$ there exists $\varrho_{\varepsilon}>0$ such that

$$
\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \cap \mathrm{B}_{r} \subset \mathrm{~L}+\mathrm{B}_{\varepsilon r} \quad \text { for all } r \in\left(0, \varrho_{\varepsilon}\right),
$$

where $\mathrm{L}+\mathrm{B}_{\varepsilon r}:=\left\{z=(x+y): x \in \mathrm{~L}, y \in \mathrm{~B}_{\varepsilon r}\right\}$ denotes the sum of sets.
Proof. - Let $w:=u(\cdot, 0)-p_{2}$, and recall that $w_{r}(x)=w(r x)$ and $\tilde{w}_{r}=$ $w_{r} /\left\|w_{r}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}$. Recall also that, by Proposition 5.4, the following limit exists

$$
\tilde{q}:=\lim _{r \downarrow 0} r^{-3} w(r \cdot),
$$

and (after choosing suitable coordinate system) the even part of $\tilde{q}$ is of the form

$$
\begin{equation*}
\tilde{q}^{e v e n}(x)=b\left|x_{n}\right|^{3}-3\left|x_{n}\right|\left(\sum_{\alpha=1}^{n-1} b_{\alpha} x_{\alpha}^{2}\right), \tag{6.27}
\end{equation*}
$$

where $b>0, b_{\alpha} \geq 0$, and $b=\sum_{\alpha=1}^{n-1} b_{\alpha}$; see Lemma 5.2. Relabelling if necessary the indices, we may assume that $b_{1} \leq b_{2} \leq \cdots \leq b_{n-1}$. In particular we must have $b_{n-1}>0$.

Define L to be the ( $n-2$ )-dimensional subspace $\left\{x_{n}=x_{n-1}=0\right\}$ in this system of coordinates. We claim that, for any sequence $\left(x_{k}, t_{k}\right) \in \boldsymbol{\Sigma}_{n-1}^{\geq 3}$ such that $x_{k} \rightarrow 0$, we have

$$
\operatorname{dist}\left(\frac{x_{k}}{\left|x_{k}\right|}, L\right) \rightarrow 0 .
$$

Note that the lemma follows immediately from this claim. To prove the claim we observe that

$$
\lambda_{k}^{2 n d}:=\phi\left(0^{+}, u\left(x_{k}+\cdot, t_{k}\right)-p_{2, x_{k}, k_{k}}\right) \geq 3 \quad \text { and } \quad \lambda^{2 n d}:=\phi\left(0^{+}, u-p_{2}\right)=3 .
$$

Thus, since the frequency is upper-semicontinuous, $\lambda_{k}^{2 n d} \rightarrow \lambda^{2 n d}=3$. This allows us to apply Lemma 6.5 with $r_{k}:=\left|x_{k}\right|$ and deduce that, if $y_{\infty}$ is an accumulation point of $\left\{x_{k} /\left|x_{k}\right|\right\}$, then the even part of $q=\frac{\tilde{q}}{\|\tilde{q}\|_{L^{2}\left(\partial B_{1}\right)}}$ is translation invariant in the direction $y_{\infty}$. Thus $\tilde{q}^{\text {even }}$ has the same invariance. But then, recalling (6.27) and $b_{n-1}>0$, we find that $y_{\infty} \in\left\{x_{n}=x_{n-1}=0\right\}=\mathrm{L}$.

We next need the following Lipschitz estimate.
Lemma 6.7. - Let $u: \mathrm{B}_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{3 r d} \backslash \Sigma_{n-1}^{\geq 4}$. Set $w:=$ $u-p_{2}-\mathrm{P}$, where P is a 3-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$, and let $w_{r}$ and $\tilde{w}_{r}$ be as in (2.1). Assume that, for some $r_{\circ}>0, \gamma \in(3,4), \delta_{\circ}>0$, and $h_{\circ}>0$, we have

$$
\begin{equation*}
\phi^{\gamma}\left(r, u-p_{2}-\mathrm{P}\right) \leq \gamma-\delta_{\circ} \quad \forall r \in\left(0, r_{\mathrm{o}}\right) \quad \text { and } \quad \mathrm{H}\left(r_{\mathrm{o}}, u-p_{2}-\mathrm{P}\right) \geq h_{\mathrm{o}} . \tag{6.28}
\end{equation*}
$$

Then there exist positive constants $\varrho_{\circ}, \eta_{\circ}$, and C , depending only on $n, \gamma, \delta_{\circ}, r_{\mathrm{o}}, h_{\mathrm{o}}$, and $\|\mathrm{P}\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$, such that for any given $\mathrm{R} \geq 1$ and for all $r \in\left(0, \frac{\varrho_{\circ}}{10 \mathrm{R}}\right)$ we have

$$
\begin{equation*}
\left\|\tilde{w}_{r}\right\|_{\mathrm{Lip}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{CR}^{3} \quad \text { and } \quad \tilde{w}_{r} \Delta \tilde{w}_{r} \geq-\mathrm{C} r^{\eta_{0}} \mathrm{R}^{4} \Delta \tilde{w}_{r} \quad \text { in } \mathrm{B}_{\mathrm{R}} \tag{6.29}
\end{equation*}
$$

Proof. - With no loss of generality we can assume that $\left\{p_{2}=0\right\}=\left\{x_{n}=0\right\}$.
Since P is some 3-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$, for any unit vector $\boldsymbol{e}$ tangential to $\left\{p_{2}=0\right\}$ we have $\left|\partial_{e e} \mathrm{P}\right| \leq \mathrm{C}\left|x_{n}\right| \leq \mathrm{Cr}^{2}$ in $\mathrm{B}_{r} \cap\{u=0\}$ (cf. (4.16)). Thus, arguing as in the proof of Lemma 4.7 (see Step 3), we get

$$
\begin{equation*}
\inf _{\mathbf{B}_{r}} r^{2} \partial_{e e} w \geq-\mathrm{C}(\mathrm{P})\left(\|w(r \cdot)\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5}\right)}+r^{4}\right) . \tag{6.30}
\end{equation*}
$$

Also, since $0 \in \Sigma_{n-1}^{3 r d} \backslash \Sigma_{n-1}^{\geq 4}$, we can apply Lemmas 4.1 and 4.9 to deduce that $\phi\left(0^{+}, u-\mathscr{P}\right)$ exists and is less that 4 (cf. proof of Proposition 4.12(a)).

We now note that, as a consequence of (6.28), Lemmas 4.3 and 4.1 yield that, for any $\delta>0, r>0$, and $\varrho \in\left(r, r_{0}\right]$,

$$
\frac{\mathrm{H}(\varrho, w)+\rho^{2 \gamma}}{\mathrm{H}(r, w)+r^{2 \gamma}} \leq \mathrm{C}_{\delta}(\varrho / r)^{2(\gamma-\delta)+\delta} .
$$

In particular, for $\delta=4-\gamma$ and $\varrho=r_{\circ}$ we obtain

$$
\begin{equation*}
\mathrm{H}\left(1, w_{r}\right)=\mathrm{H}(r, w)=\frac{\mathrm{H}\left(r_{\mathrm{o}}, w\right)+r_{\mathrm{o}}^{2 \gamma}}{\mathrm{C}_{\delta}}\left(r / r_{\mathrm{o}}\right)^{2(\gamma-\delta)+\delta}-r^{2 \gamma} \geq c_{1} r^{2 \gamma-\delta}, \tag{6.31}
\end{equation*}
$$

provided that $r \in\left(0, r_{1}\right)$, where $c_{1}>0$ and $r_{1} \in\left(0, r_{\mathrm{o}}\right)$ is sufficiently small. Also, for $r \in$ ( $0, r_{1}$ ) and $\varrho=\mathrm{R} r \leq r_{\text {o }}$ we get

$$
\mathrm{H}(\mathrm{R} r, w) \leq \mathrm{C}_{\delta} \mathrm{R}^{2 \gamma-\delta}\left(\mathrm{H}(r, w)+r^{2 \gamma}\right) \leq \mathrm{CR}^{8} \mathrm{H}(r, w),
$$

where $\mathrm{C}=\mathrm{C}_{\delta}\left(1+1 / c_{1}\right)$ depends only on $n, \gamma, \delta$, and $h_{\circ}$ Thus, scaling (6.30), for $r \in$ ( $0, \frac{r_{1}}{10 \mathrm{R}}$ ) we obtain

$$
\begin{aligned}
(2 \mathrm{R})^{-2} \inf _{\mathrm{B}_{2 \mathrm{R}}} \partial_{e e} w_{r} & \geq-\mathrm{C}(\mathrm{P})\left(\|w(3 \mathrm{R} r \cdot)\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{5}\right)}+(2 \mathrm{R} r)^{4}\right) \\
& \geq-\mathrm{CR}^{4}\left(\mathrm{H}(r, w)^{1 / 2}+r^{4}\right) \geq-\mathrm{CH}\left(1, w_{r}\right)^{1 / 2},
\end{aligned}
$$

where C depends only on $n, \mathrm{R}, \gamma, \delta$ and $h_{\circ}$.
Hence, given $\mathrm{R} \geq 1$, for all $r \in\left(0, \frac{r_{o}}{10 \mathrm{R}}\right)$ we have $\partial_{e e} \tilde{w}_{r} \geq-\mathrm{CR}^{2}$ in $\mathrm{B}_{2 \mathrm{R}}$. Therefore, as in the proof of Lemma 3.8, we obtain $\left|\nabla \tilde{w}_{r}\right| \leq \mathrm{CR}^{3}$ in $\mathrm{B}_{\mathrm{R}}$, where C depends only on $n, \gamma, \delta$, and $h_{\circ}$. This proves the first part of (6.29).

For the second part, notice that $\left|u-p_{2}-\mathrm{P}\right|=\left|\frac{1}{2}\left(x_{n}\right)^{2}+\mathrm{P}\right| \leq \mathrm{C}|x|^{4}$ inside $\{u=0\}-$ here we used that $\left|x_{n}\right| \leq \mathrm{C}|x|^{2}$ in $\{u=0\}$ and that P is a cubic polynomial divisible by $x_{n}$.

Combining this bound with (6.31) and the fact that $\Delta \tilde{w}_{r}=0$ inside $\left\{u_{r}>0\right\}$, we find (choosing for instance $\eta_{0}:=4-\gamma$ )

$$
\tilde{w}_{r} \Delta \tilde{w}_{r} \geq-\frac{\mathrm{C}|r x|^{4}}{\mathrm{H}\left(1, w_{r}\right)^{1 / 2}} \Delta \tilde{w}_{r} \geq-\frac{\mathrm{C}(r \mathrm{R})^{4}}{c r^{4-\eta_{0}}} \Delta \tilde{w}_{r}=-\mathrm{C} r^{\eta_{0}} \mathrm{R}^{4} \Delta \tilde{w}_{r} \quad \text { in } \mathrm{B}_{r},
$$

which proves (6.29).
The following result will be needed in order to bound the Hausdorff dimension of the projection $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right)$. Although the argument is very similar to the one used in the proof of Lemma 6.5, we repeat the proof in detail since there are differences that require a detailed analysis. Recall that $p_{3}=p_{3,0,0}$ is defined in (6.3).

Lemma 6.8. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1), let $(0,0)$ and $\left(x_{k}, t_{k}\right)$ belong to $\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}$, and suppose that $\left|x_{k}\right| \leq r_{k} \downarrow 0$. Assume in addition that

$$
\begin{equation*}
\tilde{w}_{r_{k}} \rightharpoonup q \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right) \text { for } w:=u-p_{2}-p_{3} \quad \text { and } \quad y_{k}:=\frac{x_{k}}{r_{k}} \rightarrow y_{\infty}, \tag{6.32}
\end{equation*}
$$

and that $\lambda_{k}^{3 r d} \rightarrow \lambda^{3 r d}$, where

$$
\begin{aligned}
& \lambda_{k}^{3 r d}:=\phi\left(0^{+}, u\left(x_{k}+\cdot, t_{k}\right)-p_{2, x_{k}, l_{k}}-p_{3, x_{k}, t_{k}}\right) \quad \text { and } \\
& \lambda^{3 r d}:=\phi\left(0^{+}, u-p_{2}-p_{3}\right) .
\end{aligned}
$$

Then $y_{\infty} \in\left\{p_{2}=0\right\}$, and $q$ is translation invariant in the direction $y_{\infty}$.
Proof. - The fact that $y_{\infty} \in\left\{p_{2}=0\right\}$ follows from Lemma 6.3.
Since $(0,0) \in \boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}$, as in the proof of Lemma 6.7 the limit $\lim _{r \downarrow 0} \phi(r$, $\left.u(\cdot, 0)-p_{2}-p_{3}\right)$ exists and belongs to (3, 4), that is

$$
\lambda^{3 r d}:=\phi\left(0^{+}, u(\cdot, 0)-p_{2}-p_{3}\right) \in(3,4) .
$$

Similarly, the limits defining $\lambda_{k}^{3 r d}$ exist, and by assumption, we have

$$
\begin{equation*}
\lambda_{k}^{3 r d}:=\phi\left(0^{+}, u\left(x_{k}+\cdot, t_{k}\right)-p_{2, x_{k}, t_{k}}-p_{3, x_{k}, t_{k}}\right) \rightarrow \lambda^{3 r d} . \tag{6.33}
\end{equation*}
$$

We define

$$
\mathfrak{p}:=p_{2}+p_{3} \quad \text { and } \quad \mathfrak{p}_{k}:=p_{2, x_{k}, l_{k}}+p_{3, x_{k}, t_{k}}
$$

and consider

$$
\begin{aligned}
w_{k} & :=u\left(x_{k}+r_{k} \cdot t_{k}\right)-\mathfrak{p}_{k}\left(r_{k} \cdot\right)=w_{k}^{(1)}+w_{k}^{(2)}+w_{k}^{(3)}, \\
w_{k}^{(1)} & :=u\left(x_{k}+r_{k} \cdot, t_{k}\right)-u\left(x_{k}+r_{k} \cdot, 0\right),
\end{aligned}
$$

$$
\begin{equation*}
w_{k}^{(2)}:=u\left(x_{k}+r_{k} \cdot, 0\right)-\mathfrak{p}\left(x_{k}+r_{k} \cdot\right) \tag{6.34}
\end{equation*}
$$

$$
w_{k}^{(3)}:=\mathfrak{p}\left(x_{k}+r_{k} \cdot\right)-\mathfrak{p}_{k}\left(r_{k} \cdot\right)
$$

Recall that $y_{k}:=x_{k} / r_{k}$ and define
(6.35)

$$
\tilde{w}_{k}:=\frac{w_{k}}{\left\|w_{k}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} .
$$

- Step 1. Throughout the proof we fix $\gamma \in\left(\lambda^{3 r d}, 4\right)$. Thanks to Lemma 4.3, for any given $\delta>0$ we have

$$
\begin{equation*}
\left|\phi^{\gamma}\left(r, u\left(x_{k}+\cdot, t_{k}\right)-p_{2, x_{k}, t_{k}}-p_{3, x_{k}, t_{k}}\right)-\lambda^{3 r d}\right| \leq \delta \quad \forall r \in\left(0, r_{\delta}\right), \forall k \geq k_{\delta} . \tag{6.36}
\end{equation*}
$$

Hence, we may fix positive constants $\delta_{\circ}$ and $r_{\circ}$ such that, for $k \geq k_{\circ}$ large enough, we have

$$
\begin{equation*}
\phi^{\gamma}\left(r, u\left(x_{k}+\cdot, t_{k}\right)-\mathfrak{p}_{k}\right) \leq \gamma-3 \delta_{\circ} \quad \forall r \in\left(0, r_{\circ}\right), \tag{6.37}
\end{equation*}
$$

and Lemma 6.7—applied to the function $u\left(x_{k}+\cdot, t_{k}\right)$ and with $r=r_{k}$-yields

$$
\begin{equation*}
\left\|\tilde{w}_{k}\right\|_{\mathrm{Lip}\left(\mathrm{~B}_{\mathrm{R}}\right)} \leq \mathrm{C}(\mathrm{R}) \quad \text { in } \mathrm{B}_{\mathrm{R}} \tag{6.38}
\end{equation*}
$$

and $\tilde{w}_{k} \Delta \tilde{w}_{k} \geq-\mathrm{C}(\mathrm{R}) r_{k}^{\eta_{0}} \Delta \tilde{w}_{k}$, where $\eta_{\circ}>0$ and $\mathrm{C}(\mathrm{R})$ are independent of $k$. Then, similarly to the proof of Lemma 6.5 , the (locally uniformly bounded) nonpositive measures $\Delta \tilde{w}_{k}$ converge weakly to $\Delta \mathrm{Q} \leq 0$, and since $r_{k}^{\eta_{0}} \Delta \tilde{w}_{k} \rightharpoonup 0$ and $\tilde{w}_{k} \rightarrow \mathrm{Q}$ locally uniformly, we have $\tilde{w}_{k} \Delta \tilde{w}_{k} \rightarrow \mathbf{Q} \Delta \mathbf{Q} \geq 0$. Furthermore, since $w_{k}=u\left(x_{k}+r_{k} \cdot, t_{k}\right) \geq 0$ on $\left\{p_{2, x_{k}, t_{k}}=0\right\}$ and $p_{2, x_{k}, t_{k}} \rightarrow p_{2}$, we obtain that $Q \geq 0$ on $\left\{p_{2}=0\right\}$. Therefore, we proved that $Q$ is a solution of the Signorini problem (3.12). Finally, arguing as in the proof of Lemma 6.5, it follows by (6.36) that the function Q is $\lambda^{3 r d}$-homogeneous.

Note that, by the same reasoning, also $q$ is a $\lambda^{3 r d}$-homogeneous of the Signorini problem (3.12).

- Step 2. Recall that $y_{\infty} \in\left\{p_{2}=0\right\}$. In addition by Proposition 4.12(a) we have

$$
\frac{w_{k}^{(2)}}{\left\|w_{k}^{(2)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightharpoonup \frac{q\left(y_{\infty}+\cdot\right)}{\left\|q\left(y_{\infty}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right),
$$

and, by construction, $w_{k}^{(3)}$ is a cubic harmonic polynomial.
We claim that, for each $k$, there exists a point $\bar{y}_{k}$ in the segment $\overline{0 y_{k}}$ such that

$$
\begin{equation*}
\left|w_{k}^{(3)}\left(\bar{y}_{k}\right)\right| \leq \mathrm{C} r_{k}^{4} . \tag{6.39}
\end{equation*}
$$

Indeed, note that $p_{2}+p_{3} \geq-\frac{p_{3}^{2}}{2 p_{2}}$ and thus we have $\mathfrak{p}\left(r_{k} \cdot\right) \geq-\mathrm{C} r_{k}^{4}$ and $\mathfrak{p}_{k}\left(r_{k} \cdot\right) \geq-\mathrm{C} r_{k}^{4}$ in $\mathrm{B}_{1}$. Hence, since $\mathfrak{p}(0)=\mathfrak{p}_{k}(0)=0$,

$$
\begin{aligned}
& w_{k}^{(3)}(0)=\mathfrak{p}(0)-\mathfrak{p}_{k}\left(-r_{k} y_{k}\right) \geq-\mathrm{C} r_{k}^{4} \quad \text { and } \\
& w_{k}^{(3)}\left(y_{k}\right)=\mathfrak{p}\left(r_{k} y_{k}\right)-\mathfrak{p}_{k}(0) \leq \mathrm{C} r_{k}^{4},
\end{aligned}
$$

so (6.39) follows.

- Step 3. Let us consider

$$
\hat{\varepsilon}_{k}:=\sum_{i=1,2,3}\left\|w_{k}^{(i)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \quad \text { and } \quad \hat{w}_{k}:=\frac{w_{k}}{\hat{\varepsilon}_{k}} .
$$

Recalling that $\phi^{\gamma}\left(0^{+}, u(\cdot, 0)-\mathfrak{p}\right)=\lambda^{3 r d}<\gamma<4$, it follows by Lemma 4.1 that, for any given $\delta>0$,

$$
\hat{\varepsilon}_{k} \geq\left\|w_{k}^{(2)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=\left\|\left(u-p_{2}-p_{3}\right)\left(r_{k}\left(y_{k}+\cdot\right)\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \gg r_{k}^{\lambda^{3 n d}+\delta} \quad \text { as } k \rightarrow \infty
$$

Thus, by Step 1, we have

$$
\hat{w}_{k} \rightarrow \hat{\mathrm{Q}}=a \mathrm{Q} \quad \text { for some } a \in[0,1] .
$$

Moreover, by Step 2,

$$
\begin{aligned}
& \hat{\mathbf{Q}}^{(2)}:=\lim _{k} w_{k}^{(2)} / \hat{\varepsilon}_{k}=b q\left(y_{\infty}+\cdot\right) \quad \text { and } \\
& \hat{\mathbf{Q}}^{(3)}:=\lim _{k} w_{k}^{(3)} / \hat{\varepsilon}_{k}=[\text { degree } 3 \text { harmonic polynomial }]
\end{aligned}
$$

for some $b \geq 0$. (Above, the convergences are weak in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$.) Thus, it is well defined

$$
\hat{\mathbf{Q}}^{(1)}:=\lim _{k} w_{k}^{(1)} / \hat{\varepsilon}_{k}=\lim _{k} w_{k} / \hat{\varepsilon}_{k}-\lim _{k} w_{k}^{(2)} / \hat{\varepsilon}_{k}-\lim _{k} w_{k}^{(3)} / \hat{\varepsilon}_{k},
$$

and we observe that $\mathrm{Q}^{(1)}$ is either nonpositive or nonnegative (since so is $w_{k}^{(1)}$ ). Hence, we have

$$
\hat{\mathbf{Q}}=\hat{\mathbf{Q}}^{(1)}+b q\left(y_{\infty}+\cdot\right)+\hat{\mathbf{Q}}^{(3)} .
$$

Moreover, it follows by (6.39) that the polynomial $\hat{\mathbf{Q}}^{(3)}$ vanishes at some point $\bar{y}$ in the segment $\overline{0 y_{\infty}}$. Hence, since $\hat{\mathbf{Q}}^{(3)}$ is harmonic, we see that it cannot have constant sign (unless it is identically zero).

Note now that, by definition of $\hat{\varepsilon}_{k}$, we have $\sum_{i}\left\|\hat{\mathbf{Q}}^{(i)}\right\|_{L^{2}\left(\partial \mathrm{~B}_{1}\right)}=1$. Hence, since $q$ is a $\lambda^{3 r d}$-homogeneous solution of Signorini with $\lambda^{3 r d}>3, \hat{\mathbf{Q}}^{(1)}$ has constant sign, and $\hat{\mathbf{Q}}^{(3)}$ is a cubic harmonic polynomial that does not have constant sign, we deduce that the three functions $\hat{\mathbf{Q}}^{(i)}$ are linearly independent and their sum $\hat{Q}$ cannot be zero.

We show next that $b>0$ and that $\hat{Q} \equiv b q$. Indeed, since both $q$ and $\hat{Q}$ are $\lambda^{3 r d}-$ homogeneous, if $\hat{\mathbf{Q}}^{(1)} \geq 0($ resp. $\leq)$ then

$$
\begin{aligned}
\hat{\mathrm{Q}} & =\lim _{\mathrm{R} \rightarrow \infty} \frac{\hat{\mathrm{Q}}(\mathrm{R} \cdot)}{\mathrm{R}^{\lambda^{3 n d}}} \\
& \left.=\lim _{\mathrm{R} \rightarrow \infty} \frac{\hat{\mathrm{Q}}^{(1)}(\mathrm{R} \cdot)+b q\left(y_{\infty}+\mathrm{R} \cdot\right)+\hat{\mathrm{Q}}^{(3)}(\mathrm{R} \cdot)}{\mathrm{R}^{\lambda^{3 n d}}} \geq b q \quad \text { (resp. } \leq\right) .
\end{aligned}
$$

But then $\hat{\mathbf{Q}}$ and $b q$ are two solution ordered solutions of Signorini with homogeneities $>1$ at the origin, and thus they must be equal by Lemma A.4.

Therefore, we have shown

$$
b\left(q-q\left(y_{\infty}+\cdot\right)\right)=\hat{\mathbf{Q}}^{(1)}+\hat{\mathbf{Q}}^{(3)}
$$

Now, using homogeneity, we obtain that for all $s>0$

$$
\hat{\mathbf{Q}}^{(1)}\left(s^{-1} x\right)+b s^{-\lambda^{3 n d}} q\left(s y_{\infty}+x\right)+\hat{\mathbf{Q}}^{(3)}\left(s^{-1} x\right)=b s^{-\lambda^{3 n d}} q(x) .
$$

If $\mathrm{Q}^{(1)} \geq 0$ (resp. if $\mathrm{Q}^{(1)} \leq 0$ ), we obtain

$$
\begin{equation*}
b \frac{q\left(s y_{\infty}+x\right)-q(x)}{s} \leq s^{\lambda^{3 n d}-1} \hat{\mathbf{Q}}^{(3)}\left(s^{-1} x\right) \quad(\text { resp. } \geq) \tag{6.40}
\end{equation*}
$$

Note that, since $q$ is a solution of (3.12) (and so it is Lipschitz continuous, see for instance [ACS08]), the absolute value of the left hand side of (6.40) is bounded as $s \downarrow 0$. Hence, since $\lambda^{3 r d} \in(3,4)$, the cubic coefficients of $\hat{\mathbf{Q}}^{(3)}$ (recall that $\hat{\mathbf{Q}}^{(3)}$ is a cubic harmonic polynomial) must vanish as otherwise the right hand side would be unbounded. Thus, the cubic coefficients of $\hat{\mathbf{Q}}^{(3)}$ vanish and therefore right hand side converges to zero.

Thus, since $b>0$, we have shown that

$$
y_{\infty} \cdot \nabla q \leq 0 \quad(\text { resp. } \geq 0)
$$

Hence, reasoning as in Step 3 of the proof of Lemma 6.5, we obtain that $\psi:=y_{\infty} \cdot \nabla q$ restricted to $\mathbf{S}^{n-1}$ must be a multiple of the first eigenfunction of a certain elliptic problem, and this easily leads to a contradiction because the homogeneity of $q$ is greater than 2.

Our next goal is to prove a variant of Lemma 6.8 for points in $\boldsymbol{\Sigma}_{n-1}^{>4} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta}$. For that, we need the following Lipschitz estimate.

Lemma 6.9. -Let $u: B_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{>4} \backslash \Sigma_{n-1}^{\geq 5-\zeta}$. Set $w:=$ $u-\mathscr{P}-\mathrm{P}$, where P is some 4-homogeneous harmonic polynomial vanishing on $\left\{p_{2}=0\right\}$. Assume that, for some $r_{\circ}>0, \gamma \in(4,5), \delta>0$, and $h_{\circ}>0$,

$$
\begin{equation*}
\phi^{\gamma}(r, u-\mathscr{P}-\mathrm{P}) \leq \gamma-\delta_{\circ} \quad \forall r \in\left(0, r_{\circ}\right) \quad \text { and } \quad \mathrm{H}\left(r_{\circ}, u-\mathscr{P}-\mathrm{P}\right) \geq h_{\circ} \tag{6.41}
\end{equation*}
$$

Then there exist positive constants $\varrho_{0}, \eta_{0}$, and C , depending only on $n, \gamma, \delta_{\circ}, r_{\mathrm{o}}$, and $h_{\mathrm{o}}$, such that for any given $\mathrm{R} \geq 1$ and for all $r \in\left(0, \frac{\varrho_{\circ}}{10 \mathrm{R}}\right)$ we have

$$
\begin{equation*}
\left\|\tilde{w}_{r}\right\|_{\mathrm{Lip}\left(\mathrm{BR}_{\mathrm{R}}\right)} \leq \mathrm{CR}^{4} \quad \text { and } \quad \tilde{w}_{r} \Delta \tilde{w}_{r} \geq-\mathrm{C} r^{\eta_{0}} \mathrm{R}^{5} \Delta \tilde{w}_{r} \quad \text { in } \mathrm{B}_{\mathrm{R}} . \tag{6.42}
\end{equation*}
$$

Proof. - The proof is analogous to the one of Lemma 6.7, using Lemma 4.7 instead of (6.30) and Lemma 4.9 instead of Lemma 4.3.

Recalling that $p_{4}=p_{4,0,0}$ is defined in (6.4), we now prove the following:
Lemma 6.10. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1), let $(0,0)$ and $\left(x_{k}, t_{k}\right)$ belong to $\boldsymbol{\Sigma}_{n-1}^{>4} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta}$ for some $\zeta \in(0,1)$, and suppose that $\left|x_{k}\right| \leq r_{k} \downarrow 0$. Assume in addition that

$$
\tilde{w}_{r_{k}} \rightharpoonup q \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right) \quad \text { for } w:=u-\mathscr{P}-p_{4} \quad \text { and } \quad y_{k}:=\frac{x_{k}}{r_{k}} \rightarrow y_{\infty}
$$

and that $\lambda_{k}^{4 \text { th }} \rightarrow \lambda^{4 t h}$, where

$$
\begin{aligned}
\lambda_{k}^{4 t h} & :=\phi\left(0^{+}, u\left(x_{k}+\cdot, t_{k}\right)-\mathscr{P}_{x_{k}, t_{k}}-p_{4, x_{k}, t_{k}}\right) \quad \text { and } \\
\lambda^{4 t h} & :=\phi\left(0^{+}, u-\mathscr{P}-p_{4}\right) .
\end{aligned}
$$

Then $y_{\infty} \in\left\{p_{2}=0\right\}$, and $q$ is translation invariant in the direction $y_{\infty}$.
Proof. - The proof is very similar to that of Lemma 6.8, with some appropriate modifications. As before, the fact that $y_{\infty} \in\left\{p_{2}=0\right\}$ follows from Lemma 6.3.

Also, since $(0,0) \in \boldsymbol{\Sigma}_{n-1}^{>4} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta}$, as in the proof of Lemma 6.7 the limit $\lim _{r \downarrow 0} \phi\left(r, u(\cdot, 0)-\mathscr{P}-p_{4}\right)$ exists and belongs to $(4,5-\zeta)$, that is

$$
\lambda^{4 t h}:=\phi\left(0^{+}, u(\cdot, 0)-\mathscr{P}-p_{4}\right) \in(4,5-\zeta) .
$$

Similarly, the limits defining $\lambda_{k}^{4 \text { th }}$ exist, and by assumption we have

$$
\lambda_{k}^{4 t h}:=\phi\left(0^{+}, u\left(x_{k}+\cdot, t_{k}\right)-\mathscr{P}_{x_{k}, t_{k}}-p_{4, x_{k}, t_{k}}\right) \rightarrow \lambda^{4+h} .
$$

We define

$$
\mathfrak{p}:=\mathscr{P}+p_{4} \quad \text { and } \quad \mathfrak{p}_{k}:=\mathscr{P}_{x_{k}, t_{k}}+p_{4, x_{k}, t_{k}},
$$

and consider $w_{k}:=u\left(x_{k}+r_{k} \cdot, t_{k}\right)-\mathfrak{p}_{k}\left(r_{k} \cdot\right)=w_{k}^{(1)}+w_{k}^{(2)}+w_{k}^{(3)}$ as in (6.34). Recall that $y_{k}:=x_{k} / r_{k}$ and define $\tilde{w}_{k}$ as in (6.35).

- Step 1. Here we argue as in Step 1 in the proof of Lemma 6.8. More precisely, using Lemma 4.3 in place of Lemma 6.7 , by the very same argument we deduce that $\tilde{w}_{k}$ converges locally uniformly to $\mathbf{Q}$, and that both $q$ and Q are $\lambda^{4 t h}$-homogeneous solutions of (3.12).
- Step 2. By Proposition 4.12(a), we have

$$
\frac{w_{k}^{(2)}}{\left\|w_{k}^{(2)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \rightharpoonup \frac{q\left(y_{\infty}+\cdot\right)}{\left\|q\left(y_{\infty}+\cdot\right)\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}} \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)
$$

and, by construction, $w_{k}^{(3)}$ is a quartic harmonic polynomial. In addition, arguing as in Step 2 of the proof of Lemma 6.8 we obtain that, for each $k$, there exists a point $\bar{y}_{k}$ in the segment $\overline{0 y_{k}}$ such that

$$
\begin{equation*}
\left|w_{k}^{(3)}\left(\bar{y}_{k}\right)\right| \leq \mathrm{Cr} r_{k}^{5} . \tag{6.43}
\end{equation*}
$$

- Step 3. Considering

$$
\hat{\varepsilon}_{k}:=\sum_{i=1,2,3}\left\|w_{k}^{(i)}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \quad \text { and } \quad \hat{w}_{k}:=\frac{w_{k}}{\hat{\varepsilon}_{k}}
$$

as in Step 3 of the proof of Lemma 6.8 we have

$$
\begin{aligned}
& \hat{w}_{k} \rightarrow \hat{\mathbf{Q}}=a \mathbf{Q}, \quad \hat{\mathbf{Q}}^{(2)}:=\lim _{k} w_{k}^{(2)} / \hat{\varepsilon}_{k}=b q\left(y_{\infty}+\cdot\right), \\
& \hat{\mathbf{Q}}^{(3)}:=\lim _{k} w_{k}^{(3)} / \hat{\varepsilon}_{k}=[\text { degree } 4 \text { harmonic pol. }],
\end{aligned}
$$

where $a \in[0,1], b \geq 0$, and all the convergences hold weakly in $\mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right)$. Hence

$$
\hat{\mathbf{Q}}=\hat{\mathbf{Q}}^{(1)}+b q\left(y_{\infty}+\cdot\right)+\hat{\mathbf{Q}}^{(3)}
$$

where $\hat{\mathbf{Q}}^{(1)}:=\lim _{k} w_{k}^{(1)} / \hat{\varepsilon}_{k}$ has constant sign. Since $q$ is a $\lambda^{4 t h}$-homogeneous solution of Signorini with $\lambda^{4 t h}>4, \hat{Q}^{(1)}$ has constant sign, and $\hat{Q}^{(3)}$ is a forth order harmonic polynomial that does not have constant sign (as a consequence of (6.43)), we deduce that the three functions $\hat{\mathbf{Q}}^{(i)}$ are linearly independent and their sum $\hat{\mathbf{Q}}$ cannot be zero.

Also, exactly as in Step 3 of the proof of Lemma $6.8, b>0$ and $\hat{\mathbf{Q}} \equiv b q$, therefore

$$
b\left(q-q\left(y_{\infty}+\cdot\right)\right)=\hat{\mathbf{Q}}^{(1)}+\hat{\mathbf{Q}}^{(3)}
$$

Now, using homogeneity, if $\mathrm{Q}^{(1)} \geq 0$ (resp. if $\mathrm{Q}^{(1)} \leq 0$ ) we obtain

$$
\frac{q\left(s y_{\infty}+x\right)-q(x, t)}{s} \leq s^{\lambda^{4 t h}-1} \hat{\mathbf{Q}}^{(3)}\left(s^{-1} x\right) \quad(\text { resp. } \geq),
$$

for all $s>0$. As in Step 3 of the proof of Lemma 6.8, this is possible only if the quartic coefficients of $\hat{\mathbf{Q}}^{(3)}$ vanishes, and letting $s \rightarrow 0$ we get

$$
y_{\infty} \cdot \nabla q \leq 0 \quad(\text { resp. } \geq 0) .
$$

Reasoning now as in Step 3 of the proof of Lemma 6.5 (see also Step 3 of the proof of Lemma 6.8), we obtain that $\psi:=y_{\infty} \cdot \nabla q$ restricted to $\mathbf{S}^{n-1}$ must be a multiple of the first eigenfunction of a certain elliptic problem, and this easily leads to a contradiction.

Before proving the last result of this section, we introduce a definition:
Definition 6.11. - We denote by $\mathcal{P}_{4, \geq}^{\text {even }}$ the set of 4-homogeneous harmonic polynomials $p=$ $p\left(x_{1}, \ldots, x_{n}\right)$, such that, for some $\boldsymbol{e} \in \mathbf{S}^{n-1}$, we have:

- $p$ is even with respect to $\{\boldsymbol{e} \cdot x=0\}$, that is, $p(x)=p(x-2(\boldsymbol{e} \cdot x) \boldsymbol{e})$;
- $p \geq 0$ on $\{\boldsymbol{e} \cdot x=0\}$;
- $\|p\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1$.

Given $p \in \mathcal{P}_{4, \geq}^{\text {even }}$, we denote $\mathbf{S}(p, \varepsilon) \subset \mathbf{R}^{n}$ the set

$$
\mathbf{S}(p, \varepsilon):=\{|\boldsymbol{e} \cdot x| \leq \varepsilon\} \cap\{p \leq \varepsilon\} \cap \overline{\mathbf{B}_{2}} .
$$

We now show the following result, which will be used later to bound the Hausdorff dimension of $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 4} \backslash \boldsymbol{\Sigma}_{n-1}^{4 \text { th }}\right)$.

Lemma 6.12. - Let $u: B_{1} \rightarrow[0, \infty)$ solve (3.1), and let $0 \in \Sigma_{n-1}^{\geq 4} \backslash \Sigma_{n-1}^{4 t h}$. Let $\mathcal{P}_{4, \geq}^{\text {even }}$ and $\mathbf{S}(p, \varepsilon)$ be as in Definition 6.11. Then, given $\varepsilon>0$, there exists $\varrho_{\varepsilon}>0$ such that, for all $r \in\left(0, \varrho_{\varepsilon}\right)$,

$$
\begin{equation*}
\{u=0\} \cap \overline{\mathbf{B}_{r}} \subset r \mathbf{S}\left(p_{r}, \varepsilon\right) \quad \text { for some } p_{r} \in \mathcal{P}_{4, \geq}^{\text {even }} . \tag{6.44}
\end{equation*}
$$

Proof. - Consider the set of "accumulation points" $\overline{\mathcal{Q}}$ defined as

$$
\overline{\mathcal{Q}}:=\left\{q: \exists r_{k} \downarrow 0 \text { s.t. } r_{k}^{-4}(u-\mathscr{P})\left(r_{k} \cdot\right) \rightarrow q\right\} .
$$

Note that, for all $\eta>0$, there exists $\varrho_{\eta}>0$ such that for any $r \in\left(0, \varrho_{\eta}\right)$ we have

$$
\begin{equation*}
\left\|u-\mathscr{P}-q_{r}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{r}\right)} \leq \eta r^{4} \quad \text { for some } q_{r} \in \overline{\mathcal{Q}} . \tag{6.45}
\end{equation*}
$$

Thanks to Proposition 4.12(a) and [GP09, Lemma 1.3.4], $\overline{\mathcal{Q}}$ is a closed set of 4homogeneous harmonic polynomials. Also, using Lemma 4.11 with $\overline{\mathrm{P}} \equiv 0$ and $\gamma \in(4,5)$ fixed, we see that $\|q\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \leq \mathrm{C}$ for all $q \in \overline{\mathcal{Q}}$. This implies that set $\overline{\mathcal{Q}}$ is compact.

Now, since by assumption $0 \in \Sigma_{n-1}^{\geq 4} \backslash \Sigma_{n-1}^{4 \text { th }}$, then $q^{\text {even }} \neq 0$ for all $q \in \overline{\mathcal{Q}}$ (recall Definition 4.10). Thus, by compactness of $\overline{\mathcal{Q}}$, we deduce that

$$
0<c_{\mathrm{o}}\left\|q^{\text {even }}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \leq\|q\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \leq \mathrm{C} \quad \forall q \in \overline{\mathcal{Q}} .
$$

Now, for $r>0$ and $q_{r}$ as in (6.45), we define

$$
p_{r}:=\frac{q_{r}^{\text {even }}}{\left\|q_{r}^{\text {even }}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}}
$$

and note that $p_{r} \in \mathcal{P}_{4, \geq}^{\text {even }}$. We claim that (6.44) holds true provided that $r \in\left(0, \rho_{\varepsilon}\right)$, with $\rho_{\varepsilon}>0$ small.

Indeed, assume with no loss of generality that $\left\{p_{2}=0\right\}=\left\{x_{n}=0\right\}$. Then (since $q_{r}$ solves (3.12)) every $p_{r}$ is a 4-homogeneous harmonic polynomial, even in the variable $x_{n}$, nonnegative on $\left\{x_{n}=0\right\}$, and with unit $\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)$ norm.

We recall that

$$
\begin{equation*}
\mathscr{P}(x) \geq\left(x_{n}+p_{3} / x_{n}+\mathrm{Q}\right)^{2}-\mathrm{C}|x|^{5} . \tag{6.46}
\end{equation*}
$$

Now, by definition of $\mathbf{S}(p, \varepsilon)$, it follows in particular that, fixed $\theta>0$,

$$
y \in \mathbf{B}_{2} \backslash \mathbf{S}\left(p_{r}, \varepsilon\right) \quad \Rightarrow \quad \text { either }\left(p_{r}(y)>\varepsilon \text { and }\left|\nu_{n}\right| \leq \theta \varepsilon\right) \text { or }\left(\nu_{n} \mid>\theta \varepsilon\right) .
$$

We now observe that, if $p_{r}(y)>\varepsilon$ and $\left|\nu_{n}\right| \leq \theta \varepsilon$, since $q_{r}^{\text {odd }}$ vanishes on $\left\{x_{n}=0\right\}$ we get

$$
q_{r}(y)=p_{r}(y)\left\|q^{\text {even }}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}+q_{r}^{\text {odd }}(y) \geq c_{0} \varepsilon-\mathrm{C}\left|y_{n}\right| \geq\left(c_{\circ}-\mathrm{C} \theta\right) \varepsilon \geq \frac{1}{2} c_{0} \varepsilon>0
$$

provided we choose $\theta:=\frac{c_{0}}{2 \mathrm{C}}$ small enough. Thus, recalling (6.45) and (6.46), if $r>0$ is sufficiently small (so that we can take $\eta \ll \varepsilon$ ) we get

$$
u(r y) \geq \mathscr{P}(r y)+q_{r}(r y)-\eta r^{4} \geq-\mathrm{C} r^{5}+\frac{1}{2} c_{\circ} \varepsilon r^{4}-\eta r^{4}>0
$$

On the other hand, if $\left|y_{n}\right|>\theta \varepsilon$, using again (6.46) we obtain, for $r>0$ sufficiently small,

$$
u(r y) \geq \mathscr{P}(r y)+q_{r}(r y)-\eta r^{4} \geq\left(\theta \varepsilon r-\mathrm{Cr} r^{2}\right)^{2}-\mathrm{C} r^{5}-\mathrm{C} r^{4}-\eta r^{4}>0 .
$$

Therefore, we have proven that

$$
y \in \mathrm{~B}_{2} \backslash \mathbf{S}\left(p_{r}, \varepsilon\right) \quad \Rightarrow \quad u(r y)>0,
$$

which gives (6.44).

## 7. Hausdorff measures and covering arguments

As already explained in the introduction, to prove our main results we will need some auxiliary results from geometric measure theory. Before stating them, we recall some classical definitions.

Given $\beta>0$ and $\delta \in(0, \infty]$, the Hausdorff premeasures $\mathcal{H}_{\delta}^{\beta}(\mathrm{E})$ of a set E are defined as follows: ${ }^{7}$

$$
\begin{equation*}
\mathcal{H}_{\delta}^{\beta}(\mathrm{E}):=\inf \left\{\sum_{i} \operatorname{diam}\left(\mathrm{E}_{i}\right)^{\beta}: \mathrm{E} \subset \bigcup_{i} \mathrm{E}_{i}, \operatorname{diam}\left(\mathrm{E}_{i}\right)<\delta\right\}, \tag{7.1}
\end{equation*}
$$

where the index $i$ goes through a finite or countable set. Then, one defines the $\beta$ dimensional Hausdorff measure of E as $\mathcal{H}^{\beta}(\mathrm{E}):=\lim _{\delta \rightarrow 0^{+}} \mathcal{H}_{\delta}^{\beta}(\mathrm{E})$.

The Hausdorff dimension can be defined in terms of $\mathcal{H}_{\infty}^{\beta}$ as follows:

$$
\begin{equation*}
\operatorname{dim}_{\mathcal{H}}(\mathrm{E}):=\inf \left\{\beta>0: \mathcal{H}_{\infty}^{\beta}(\mathrm{E})=0\right\} \tag{7.2}
\end{equation*}
$$

(this follows from the fact that $\mathcal{H}_{\infty}^{\beta}(\mathrm{E})=0$ if and only if $\mathcal{H}^{\beta}(\mathrm{E})=0$, see for instance [Sim83, Section 1.2]).

We now state (and prove, for completeness) a couple of standard results.

[^7]Lemma 7.1. - Let $\mathrm{E} \subset \mathbf{R}^{n}$, and $f: \mathrm{E} \rightarrow \mathbf{R}$. Define

$$
\mathrm{F}:=\left\{x \in \mathrm{E}: \exists x_{k} \rightarrow x, x_{k} \in \mathrm{E}, \text { s.t. } f\left(x_{k}\right) \rightarrow f(x)\right\} .
$$

Then $\mathrm{E} \backslash \mathrm{F}$ is at most countable.
Proof.- Let $\mathrm{G}:=\{(x, f(x)): x \in \mathrm{E}\} \subset \mathbf{R}^{n} \times \mathbf{R}$ be the graph of $f$. We note that $x \in \mathrm{E} \backslash \mathrm{F}$ if a only if $(x, f(x))$ is a isolated point of G . In particular $\mathrm{E} \backslash \mathrm{F}$ is the projection of a discrete (and hence countable) set.

From now on, by convention, whenever we say that a set E can be covered by a number $\mathrm{M}>0$ of balls that it is not necessarily an integer, we mean that it can be covered by $\lfloor M\rfloor$ balls, where $\lfloor M\rfloor$ denotes the integer part of $M$.

Lemma 7.2. - Let $\mathrm{B}_{r}(x) \subset \mathbf{R}^{n}$ be an open ball, and $\Pi$ be a m-dimensional plane. Let $\beta_{1}>m$. Then there exists $\hat{\varepsilon}=\hat{\varepsilon}\left(m, \beta_{1}\right)>0$ such that the following holds: Let $\mathrm{E} \subset \mathbf{R}^{n}$ satisfy

$$
\mathrm{E} \subset \mathrm{~B}_{r}(x) \cap\{y: \operatorname{dist}(y, \Pi) \leq \varepsilon r\}, \quad \text { for some } 0<\varepsilon \leq \hat{\varepsilon}, x \in \mathbf{R}^{n}, r>0 .
$$

Then E be covered with $\gamma^{-\beta_{1}}$ balls of radius $\gamma r$ centered at points of E , where $\gamma:=5 \varepsilon$.
Proof. - Up to a scaling and a translation, it suffices to prove the result when $r=1$ and $\mathrm{B}_{r}(x)$ is the unit ball $\mathrm{B}_{1}$ centered at the origin. Consider the $m$-dimensional set $\mathrm{B}_{1} \cap \Pi$, and given $\varepsilon>0$ small consider the covering of $\mathrm{E} \subset \mathrm{B}_{1} \cap\{y: \operatorname{dist}(y, \Pi) \leq \varepsilon\}$ given by the closed balls $\left\{\overline{\mathrm{B}_{\varepsilon}(x)}\right\}_{x \in \mathrm{E}}$. By Vitali Covering Lemma, there exists a disjoint family $\left\{\overline{\mathbf{B}_{\varepsilon}\left(x_{i}\right)}\right\}_{i \in \mathcal{I}}$ such that

$$
\bigcup_{i \in \mathcal{I}} \overline{\mathrm{~B}_{5 \varepsilon}\left(x_{i}\right)} \supset \bigcup_{x \in \mathbb{E}} \overline{\mathrm{~B}_{\varepsilon}(x)} \supset \mathrm{E} .
$$

Note that

$$
\overline{\mathbf{B}_{\varepsilon}\left(x_{i}\right)} \subset \mathcal{N}_{2 \varepsilon}(\Pi):=\left\{x \in \mathrm{~B}_{2}: \operatorname{dist}(x, \Pi) \leq 2 \varepsilon\right\} .
$$

Since $\mathcal{H}^{n}\left(\mathcal{N}_{2 \varepsilon}(\Pi)\right) \leq \mathrm{C}(n) \varepsilon^{n-m}$, denoting by $\omega_{n}$ the volume of the $n$-dimensional unit ball we have

$$
\omega_{n} \varepsilon^{n} \# \mathcal{I} \leq \sum_{i \in \mathcal{I}} \mathcal{H}^{n}\left(\mathrm{~B}_{\varepsilon}\left(x_{i}\right)\right) \leq \mathcal{H}^{n}\left(\mathcal{N}_{2 \varepsilon}(\Pi)\right)=\mathrm{C}(n) \varepsilon^{n-m},
$$

which proves that $\# \mathcal{I} \leq \mathrm{C}(n) \varepsilon^{-m}$. Set $\gamma:=5 \varepsilon$. Then, since $\beta_{1}>m$, choosing $\varepsilon$ sufficiently small we have $\mathrm{C}(n) \varepsilon^{-m}=\mathrm{C}(n) 5^{m} \gamma^{-m} \leq \gamma^{-\beta_{1}}$, proving that E can be covered by $\gamma^{-\beta_{1}}$ open balls of radius $\gamma$ centered at points of E , as desired.

The following Reifenberg-type result will be used later choosing as function $f$ the frequency function, and it will allow us to perform our dimension reduction arguments only at continuity points of the frequency.

Proposition 7.3. - Let $\mathrm{E} \subset \mathbf{R}^{n}$, and $f: \mathrm{E} \rightarrow \mathbf{R}$. Assume that, for any $\varepsilon>0$ and $x \in \mathrm{E}$ there exists $\varrho=\varrho(x, \varepsilon)>0$ such that, for all $r \in(0, \varrho)$, we have

$$
\mathrm{E} \cap \overline{\mathrm{~B}_{r}(x)} \cap f^{-1}([f(x)-\varrho, f(x)+\varrho]) \subset\left\{y: \operatorname{dist}\left(y, \Pi_{x, r}\right) \leq \varepsilon r\right\}
$$

for some m-dimensional plane $\Pi_{x, r}$ passing through $x$ (possibly depending on $r$ ). Then $\operatorname{dim}_{\mathcal{H}}(\mathrm{E}) \leq m$.
Proof. - We need to prove that, given $\beta>m$, we have $\mathcal{H}^{\beta}(\mathrm{E})=0$.
Let $\varepsilon>0$ be a small constant to be fixed later, and for any $k>1$ and $j \in \mathbf{Z}$ define

$$
\mathrm{E}_{k, j}:=\left\{x \in \mathrm{E}: \varrho(x, \varepsilon)>1 / k, f(x) \in\left[\frac{j}{2 k}, \frac{j+1}{2 k}\right)\right\} .
$$

Since $\mathrm{E}=\bigcup_{k, j} \mathrm{E}_{k, j}$, it suffices to prove that $\mathcal{H}^{\beta}\left(\mathrm{E}_{k, j}\right)=0$ for each $k, j$. So, we fix $k>1$ and $j \in \mathbf{Z}$. Similarly, it suffices to prove that for all $\mathrm{R} \geq 1$ we have $\mathcal{H}^{\beta}\left(\mathrm{E}_{k, j}^{\mathrm{R}}\right)=0$, where $\mathrm{E}_{k, j}^{\mathrm{R}}:=\mathrm{B}_{\mathrm{R}} \cap \mathrm{E}_{k, j}$.

By assumption, for every $x \in \mathrm{E}_{k, j}^{\mathrm{R}}$ and $r \in(0,1 / k]$, there exists a $m$-dimensional plane $\Pi_{x, r}$ such that

$$
\mathrm{E}_{k, j}^{\mathrm{R}} \cap \overline{\mathrm{~B}_{r}(x)} \subset\left\{y: \operatorname{dist}\left(y, \Pi_{x, r}\right) \leq \varepsilon r\right\} .
$$

So, we consider the covering $\left\{\overline{\mathrm{B}_{1 / k}(x)}\right\}_{x \in \mathrm{E}_{k, j}^{\mathrm{R}}}$, and since $\mathrm{E}_{k, j}^{\mathrm{R}} \subset \mathrm{B}_{\mathrm{R}}$ we extract a finite subcovering of closed balls $\mathrm{B}_{1}^{(0)}, \ldots, \mathrm{B}_{\mathrm{M}}^{(0)}$. (Indeed, by Vitali's lemma there is a covering $\left\{\overline{\mathrm{B}_{1 / k}\left(x_{\ell}\right)}\right\}$ such that the balls $\left\{\overline{\mathrm{B}_{1 /(5 k)}\left(x_{\ell}\right)}\right\}$ are disjoint, and hence there is a finite number of them.) Inside each ball $\mathrm{B}_{i}^{(0)}$ we have, by assumption,

$$
\mathrm{E}_{k, j}^{\mathrm{R}} \cap \mathrm{~B}_{i}^{(0)} \subset\left\{y: \operatorname{dist}\left(y, \Pi_{\mathrm{B}_{i}^{(0)}}\right) \leq \varepsilon / k\right\} .
$$

Choose $\beta_{1}:=\frac{m+\beta}{2} \in(m, \beta)$. Applying Lemma 7.2 with $\varepsilon=\hat{\varepsilon}\left(m, \beta_{1}\right)$ we deduce that, for each fixed $i, j, k, \mathbf{R}$, the set $\mathrm{E}_{k, j}^{\mathrm{R}} \cap \mathbf{B}_{i}^{(0)}$ can be covered with $\gamma^{-\beta_{1}}$ closed balls $\hat{\mathbf{B}}_{1}^{(1)}, \ldots, \hat{\mathbf{B}}_{\gamma^{-\beta_{1}}}^{(1)}$ of radius $\gamma / k$ centered at points of $\mathrm{E}_{k, j}^{\mathrm{R}} \cap \mathrm{B}_{i}^{(0)}$, where $\gamma=5 \varepsilon$. Using our assumption again, in each of these balls we have

$$
\mathrm{E}_{k, j}^{\mathrm{R}} \cap \hat{\mathbf{B}}_{\ell}^{(1)} \subset\left\{y: \operatorname{dist}\left(y, \Pi_{x_{\ell}^{(1)}}\right) \leq \varepsilon \gamma / k\right\},
$$

where $x_{\ell}^{(1)}$ is the centre of $\hat{\mathbf{B}}_{\ell}^{(1)}$. We then apply again Lemma 7.2 so that, for each $\ell \in$ $\left\{1, \ldots, \gamma^{-\beta_{1}}\right\}$, we can cover the set $\mathrm{E}_{k, j}^{\mathrm{R}} \cap \hat{\mathbf{B}}_{\ell}^{(1)}$ with $\gamma^{-\beta_{1}}$ closed balls of radius $\gamma^{2} / k$. This gives a new covering of $\mathbf{E}_{k, j}^{\mathrm{R}} \cap \mathrm{B}_{i}^{(0)}$ with $\gamma^{-2 \beta_{1}}$ closed balls $\hat{\mathbf{B}}_{1}^{(2)}, \ldots, \hat{\mathbf{B}}_{\gamma^{-2 \beta_{1}}}^{(2)}$ of radius $\gamma^{2} / k$
centered at points of $\mathrm{E}_{k, j}^{\mathrm{R}}$. Iterating this construction, we conclude that $\mathrm{E}_{k, j}^{\mathrm{R}} \cap \mathrm{B}_{i}^{(0)}$ can be covered by $\gamma^{-\mathrm{N} \beta_{1}}$ closed balls $\left\{\hat{\mathbf{B}}_{\ell}^{(\mathrm{N})}\right\}$ of radius $\gamma^{\mathrm{N}} / k$ for any $\mathrm{N} \geq 1$, which implies that

$$
\mathcal{H}_{\infty}^{\beta}\left(\mathrm{E}_{k, j}^{\mathrm{R}} \cap \mathrm{~B}_{i}\right) \leq \mathrm{C}_{n, m} \sum_{\ell} \operatorname{diam}\left(\hat{\mathrm{B}}_{\ell}^{(\mathrm{N})}\right)^{\beta} \leq \mathrm{C}_{n, m} \gamma^{-\mathrm{N} \beta_{1}}\left(\gamma^{\mathrm{N}} / k\right)^{\beta} \leq \mathrm{C} \gamma^{\mathrm{N}\left(\beta-\beta_{1}\right)}
$$

Since $\beta_{1} \in(m, \beta)$, letting $N \rightarrow \infty$ we conclude that

$$
\mathcal{H}_{\infty}^{\beta}\left(\mathrm{E}_{k, j}^{\mathrm{R}} \cap \mathrm{~B}_{i}^{(0)}\right)=0 \quad \text { for all } i, j, k, \mathrm{R},
$$

concluding the proof.
In our study of 4-homogeneous blow-ups, we will need a variant of the previous results involving zero sets of 4-homogeneous harmonic polynomials instead of hyperplanes (recall Definition 6.11).

Lemma 7.4. - Given $\beta_{1}>n-2$, there exists $\hat{\varepsilon}=\hat{\varepsilon}\left(n, \beta_{1}\right)>0$ such that the following holds: Let $\mathrm{E} \subset \mathbf{R}^{n}$ satisfy

$$
\mathrm{E} \subset \mathrm{~B}_{1} \cap \mathbf{S}(p, \varepsilon), \quad \text { for some } 0<\varepsilon \leq \hat{\varepsilon}, p \in \mathcal{P}_{4, \geq}^{\text {even }} .
$$

Then E can be covered with $\gamma^{-\beta_{1}}$ balls of radius $\gamma$ centered at points of E , for some $\gamma=\gamma\left(n, \beta_{1}\right) \in$ $(0,1)$.

Proof. - Given $t, \varepsilon>0$ small, consider the covering of $\mathrm{E} \subset \mathrm{B}_{1} \cap\{y: \operatorname{dist}(y, \mathbf{S}(p, \varepsilon))$ $\leq t\}$ given by $\left\{\overline{\mathrm{B}_{t}(x)}\right\}_{x \in \mathrm{E}}$. By Vitali's lemma, there exists a disjoint family $\left\{\overline{\mathrm{B}_{t}\left(x_{i}\right)}\right\}_{i \in \mathcal{I}}$ such that

$$
\bigcup_{i \in \mathcal{I}} \overline{\mathrm{~B}_{5 t}\left(x_{i}\right)} \supset \bigcup_{x \in \mathrm{E}} \overline{\mathrm{~B}_{t}(x)} \supset \mathrm{E} .
$$

Note that, since $x_{i} \in \mathrm{E}$,

$$
\overline{\mathbf{B}_{t}\left(x_{i}\right)} \subset \mathcal{N}_{2 t}(\mathbf{S}(p, \varepsilon)):=\left\{x \in \mathbf{B}_{2}: \operatorname{dist}(x, \mathbf{S}(p, \varepsilon)) \leq 2 t\right\} .
$$

We claim that there exists a dimensional constant $\mathrm{C}(n)$ such that, for any given $t \in(0,1)$, there is $\varepsilon_{t}>0$ such that

$$
\begin{equation*}
\mathcal{H}^{n}\left(\mathcal{N}_{2 t}(\mathbf{S}(p, \varepsilon))\right) \leq \mathrm{C}(n) t^{2} \quad \forall \varepsilon \in\left(0, \varepsilon_{t}\right) \tag{7.3}
\end{equation*}
$$

Indeed, if not, then for arbitrarily large $M$ there would exist some $t_{M} \in(0,1)$ and sequences $\varepsilon_{k} \downarrow 0$ and $p_{k} \in \mathcal{P}_{4, \geq}^{\text {even }}$ such that

$$
\begin{equation*}
\mathcal{H}^{n}\left(\mathcal{N}_{2 \mathrm{~m}_{\mathrm{M}}}\left(\mathbf{S}\left(p_{k}, \varepsilon_{k}\right)\right)\right) \geq \mathrm{M} t_{\mathrm{M}}^{2} \quad \forall k \geq 1 \tag{7.4}
\end{equation*}
$$

Now, given $p \in \mathrm{P}_{4, \geq}^{\text {even }}$ which is even with respect to the hyperplane $\{\boldsymbol{e} \cdot x=0\}$ and nonnegative on it, let us define

$$
\boldsymbol{z}(p):=\{p=0\} \cap\{\boldsymbol{e} \cdot x=0\} .
$$

Notice that, by definition of $\mathbf{S}(p, \varepsilon)$, for all $p \in \mathrm{P}_{4, \geq}^{\text {even }}$ we have

$$
\begin{equation*}
\mathbf{S}(p, \varepsilon) \downarrow \boldsymbol{z}(p) \quad \text { as } \varepsilon \downarrow 0 \tag{7.5}
\end{equation*}
$$

In addition, for all $x \in \boldsymbol{z}(p)$, we have that $\boldsymbol{e} \cdot \nabla p(x)=0$ (since $p$ is even with respect $\{\boldsymbol{e} \cdot x=0\}$ ). Furthermore, the tangential gradient vanishes at $x \in \boldsymbol{z}(p)$ (since $p \geq 0$ on $\{\boldsymbol{e} \cdot x=0\}$ and $p(x)=0)$. Hence, this proves that

$$
\begin{equation*}
\boldsymbol{z}(p) \subset\{p=|\nabla p|=0\} . \tag{7.6}
\end{equation*}
$$

Let $p_{k}$ be even with respect to $\boldsymbol{e}_{k} \in \mathbf{S}^{n-1}$, and assume without loss of generality that $p_{k} \rightarrow$ $p_{\infty} \in \mathcal{P}_{4, \geq}^{\text {even }}$ and that $\boldsymbol{e}_{k} \rightarrow \boldsymbol{e}_{\infty}$. Then it follows by (7.5) that, for all $\delta>0$, there exists $k_{\delta}$ such that

$$
\mathcal{N}_{2 \mathfrak{M}_{\mathrm{M}}}\left(\mathbf{S}\left(p_{k}, \varepsilon_{k}\right)\right) \subset\left\{x \in \mathrm{~B}_{2}: \operatorname{dist}\left(x, \boldsymbol{z}\left(p_{\infty}\right)\right) \leq 2 t_{\mathrm{M}}+\delta\right\}, \quad \forall k \geq k_{\delta} .
$$

Recalling (7.4), this implies that

$$
\mathcal{H}^{n}\left(\left\{x \in \mathrm{~B}_{2}: \operatorname{dist}\left(x, \boldsymbol{z}\left(p_{\infty}\right)\right) \leq 2 t_{\mathrm{M}}\right\}\right) \geq \mathrm{M} t_{\mathrm{M}}^{2}
$$

On the other hand, [NV17, Theorem 1.1] implies the existence of a dimensional constant $\mathrm{C}(n)$ such that

$$
\mathcal{H}^{n}\left(\left\{x \in \mathrm{~B}_{2}: \operatorname{dist}(x,\{u=|\nabla u|=0\}) \leq 2 t\right\}\right) \leq \mathrm{C}(n) t^{2} \quad \forall t \in(0,1)
$$

for every nonzero harmonic function $u$ in $\mathrm{B}_{4}$. Recalling (7.6), we obtain a contradiction by choosing $\mathrm{M}>\mathrm{C}(n)$.

Now, denoting by $\omega_{n}$ the volume of the $n$-dimensional unit ball, given $t \in(0,1)$, thanks to (7.3) we have

$$
\omega_{n} t^{n} \# \mathcal{I} \leq \sum_{i \in \mathcal{I}} \mathcal{H}^{n}\left(\mathrm{~B}_{t}\left(x_{i}\right)\right) \leq \mathcal{H}^{n}\left(\mathrm{~N}_{2 t}(\mathbf{S}(p, \varepsilon))=\mathrm{C}(n) t^{2} \quad \forall \varepsilon \in\left(0, \varepsilon_{t}\right),\right.
$$

which proves that $\# \mathcal{I} \leq \mathrm{C}(n) t^{2-n}$. Set $\gamma:=5 t$. Since $\beta_{1}>n-2$, choosing $t$ sufficiently small we have $\mathrm{C}(n) t^{2-n}=\mathrm{C}(n) 5^{n-2} \gamma^{2-n} \leq \gamma^{-\beta_{1}}$, proving that E can be covered by $\gamma^{-\beta_{1}}$ open balls of radius $\gamma$ centered at points of E whenever $\varepsilon<\hat{\varepsilon}:=\varepsilon_{t}$.

Proposition 7.5. - Let $\mathrm{E} \subset \mathbf{R}^{n}$ be a measurable set, and $\tau: \mathrm{E} \rightarrow \mathbf{R}$ a lower-semicontinuous function. Assume that, for any $\varepsilon>0$ and $x \in \mathrm{E}$, there exists $\varrho=\varrho(x, \varepsilon)>0$ such that, for all $r \in(0, \varrho)$, we have

$$
\mathrm{E} \cap \overline{\mathbf{B}_{r}(x)} \cap \tau^{-1}([\tau(x),+\infty)) \subset\left\{x+r y: y \in \mathbf{S}\left(p_{x, r}, \varepsilon\right)\right\}
$$

for some $p_{x, r} \in \mathcal{P}_{4, \geq}^{\text {even }}$. Then $\operatorname{dim}_{\mathcal{H}}(\mathrm{E}) \leq n-2$.

Proof. - Given $\beta>n-2$, we need to prove that $\mathcal{H}^{\beta}(\mathrm{E})=0$. Let $\varepsilon>0$ be a small constant to be fixed later, and for any $k>1$ define

$$
\mathrm{E}_{k}:=\{x \in \mathrm{E}: \rho(x) \geq 1 / k\}
$$

Since $\mathrm{E}=\bigcup_{k} \mathrm{E}_{k}$, it suffices to prove that $\mathcal{H}^{\beta}\left(\mathrm{E}_{k}\right)=0$ for each $k$. So, we fix $k>1$. Thanks to [Fed69, Corollary 2.10.23], it suffices to prove that $\mathcal{H}^{\beta}(\mathrm{K})=0$ for any compact set K contained inside $\mathrm{E}_{k}$.

We now claim the following: For each closed ball $\overline{\mathbf{B}_{r}(x)}$ centered at a point $x \in \mathrm{E}$ and of radius $r \leq 1 / k$, there exist $\bar{x} \in \mathrm{~K} \cap \overline{\mathrm{~B}_{r}(x)}$ and $p_{\bar{x}, 2 r} \in \mathcal{P}_{4, \geq}^{\text {even }}$ such that

$$
\mathrm{K} \cap \overline{\mathrm{~B}_{r}(x)} \subset \bar{x}+r \mathbf{S}\left(p_{\bar{x}, r}, \varepsilon\right)
$$

To prove this claim it suffices to observe that it is trivially true if $\mathrm{K} \cap \overline{\mathrm{B}_{r}(x)}$ is empty. Otherwise, the lower semicontinuous function $\tau$ attains its minimum at some point $\bar{x} \in$ $\mathrm{K} \cap \overline{\mathrm{B}_{r}(x)}$. Then, by the assumption of the lemma,

$$
\begin{aligned}
\mathrm{K} \cap \overline{\mathbf{B}_{r}(x)} & =\mathrm{K} \cap \overline{\mathbf{B}_{r}(x)} \cap \tau^{-1}([\tau(\bar{x}), \infty)) \subset \mathrm{E} \cap \overline{\mathrm{~B}_{2 r}(\bar{x})} \cap \tau^{-1}([\tau(\bar{x}), \infty)) \\
& \subset \bar{x}+r \mathbf{S}\left(p_{\bar{x}, r}, \varepsilon\right)
\end{aligned}
$$

which proves the claim.
Now, consider the covering $\left\{\overline{\mathrm{B}_{1 / k}(x)}\right\}_{x \in \mathrm{~K}}$, and extract a finite subcovering of closed balls $\mathbf{B}_{1}^{(0)}, \ldots, \mathbf{B}_{\mathrm{M}}^{(0)}$. Inside each ball $\mathrm{B}_{i}^{(0)}$ we can apply the claim to deduce that

$$
\mathrm{K} \cap \mathrm{~B}_{i}^{(0)} \subset \bar{x}_{i}+r \mathbf{S}\left(p_{\bar{x}_{i}, r}, \varepsilon\right) .
$$

Applying now Lemma 7.4 we deduce that, for each fixed $i$, the set $\mathrm{K} \cap \mathbf{B}_{i}^{(0)}$ can be covered with $\gamma^{-\beta_{1}}$ closed balls $\hat{\mathbf{B}}_{1}^{(1)}, \ldots, \hat{\mathbf{B}}_{\gamma^{-\beta_{1}}}^{(1)}$ of radius $\gamma / k$ centered at points of E . In each of these balls we now reapply the claim to deduce that

$$
\mathrm{K} \cap \hat{\mathbf{B}}_{\ell}^{(1)} \subset \bar{x}_{\ell}^{(1)}+\frac{\gamma}{k} \mathbf{S}\left(p_{\bar{x}_{\ell}^{(1)}, \frac{\gamma}{k}}, \varepsilon\right)
$$

Thus we can apply again Lemma 7.4 (rescaled) to cover, for each $\ell$, the set $\mathrm{K} \cap \hat{\mathbf{B}}_{\ell}^{(1)}$ with $\gamma^{-\beta_{1}}$ closed balls. In this way we obtain a new covering of $\mathrm{K} \cap \mathrm{B}_{i}^{(0)}$ by $\gamma^{-2 \beta_{1}}$ closed balls $\hat{\mathbf{B}}_{1}^{(2)}, \ldots, \hat{\mathbf{B}}_{\gamma^{-2 \beta_{1}}}^{(2)}$ of radius $\gamma^{2} / k$ centered at points of E . Iterating this construction, we conclude that $\mathrm{K} \cap \mathrm{B}_{i}^{(0)}$ can be covered by $\gamma^{-\mathrm{N} \beta_{1}}$ closed balls $\left\{\hat{\mathbf{B}}_{\ell}^{(\mathrm{N})}\right\}$ of radius $\gamma^{\mathrm{N}} / k$ for any $\mathrm{N} \geq 1$, which implies that

$$
\mathcal{H}_{\infty}^{\beta}\left(\mathrm{K} \cap \overline{\mathrm{~B}_{i}}\right) \leq \mathrm{C}_{n, m} \sum_{\ell} \operatorname{diam}\left(\hat{\mathbf{B}}_{\ell}^{(\mathrm{N})}\right)^{\beta} \leq \mathrm{C}_{n, m} \gamma^{-\mathrm{N} \beta_{1}}\left(\frac{\gamma^{\mathrm{N}}}{k}\right)^{\beta} \leq \mathrm{C} \gamma^{\mathrm{N}\left(\beta-\beta_{1}\right)}
$$

Since $\beta_{1} \in(n-2, \beta)$, letting $N \rightarrow \infty$ we conclude that

$$
\mathcal{H}_{\infty}^{\beta}\left(\mathrm{K} \cap \mathrm{~B}_{i}^{(0)}\right)=0 \quad \text { for all } i=1,2, \ldots, \mathrm{M}
$$

This proves that $\mathcal{H}_{\infty}^{\beta}(\mathrm{K})=0$ and therefore $\mathcal{H}^{\beta}(\mathrm{K})=0($ see [Sim83, Section 1]), concluding the proof.

We will also use the following basic result about Hausdorff measures. We refer to [Fed69, 2.10.19(2)] and [FS19, Lemma 3.5] for a proof of such result; see also [Whi97, Lemma 2.4].

Lemma 7.6. - Let $\mathrm{E} \subset \mathbf{R}^{n}$ be a set satisfying $\mathcal{H}_{\infty}^{\beta}(\mathrm{E})>0$ for some $\beta \in(0, n]$. Then:
(a) For $\mathcal{H}^{\beta}$-almost every point $x_{0} \in \mathrm{E}$, there is a sequence $r_{k} \downarrow 0$ such that

$$
\begin{equation*}
\lim _{k \rightarrow \infty} \frac{\mathcal{H}_{\infty}^{\beta}\left(\mathrm{E} \cap \mathrm{~B}_{r_{k}}\left(x_{\circ}\right)\right)}{r_{k}^{\beta}} \geq c_{n, \beta}>0 \tag{7.7}
\end{equation*}
$$

where $c_{n, \beta}$ is a constant depending only on $n$ and $\beta$. We call these points "density points".
(b) Assume that 0 is a "density point", let $r_{k} \downarrow 0$ be a sequence along which (7.7) holds, and define the "accumulation set" for E at 0 along $r_{k}$ as

$$
\mathcal{A}=\mathcal{A}_{\mathrm{E}}:=\left\{z \in \overline{\mathrm{~B}_{1}}: \exists(z \ell)_{\ell \geq 1},\left(k_{\ell}\right)_{\ell \geq 1} \text { s.t. } z \ell \in r_{k \ell}^{-1} \mathrm{E} \cap \mathrm{~B}_{1} \text { and } z_{\ell} \rightarrow z\right\} .
$$

Then $\mathcal{H}_{\infty}^{\beta}(\mathcal{A})>0$.
The last main result of this section is the following covering-type result that will play a crucial role in the understanding of the generic size of the singular set, and in particular in the proof of Theorem 1.1. Notice that, when $k=1, \beta$ is an integer, and $\pi_{1}(\mathrm{E})$ is $\beta$ rectifiable, then the result follows from the coarea formula; see also Eilenberg's inequality [BZ80, 13.3].

Proposition 7.7. - Let $\mathrm{E} \subset \mathbf{R}^{n} \times[-1,1]$, let $(x, t)$ denote a point in $\mathbf{R}^{n} \times[-1,1]$, and let $\pi_{1}:(x, t) \mapsto x$ and $\pi_{2}:(x, t) \mapsto t$ be the standard projections. Assume that for some $\beta \in(0, n]$ and $s>0$ we have:

- $\mathcal{H}^{\beta}\left(\pi_{1}(\mathrm{E})\right)<+\infty$;
- For any $\left(x_{0}, t_{0}\right) \in \mathrm{E}$ there exists a modulus of continuity $\omega_{x_{0}, t_{0}}: \mathbf{R}^{+} \rightarrow \mathbf{R}^{+}$such that

$$
\left\{(x, t) \in \mathbf{R}^{n} \times[-1,1]: t-t_{0}>\omega_{x_{0}, t_{0}}\left(\left|x-x_{0}\right|\right)\left|x-x_{0}\right|^{s}\right\} \cap \mathrm{E}=\emptyset .
$$

Then:
(a) If $\beta \leq s$, we have $\mathcal{H}^{\beta / s}\left(\pi_{2}(\mathrm{E})\right)=0$.
(b) If $\beta>s$, for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$ we have $\mathcal{H}^{\beta-s}\left(\mathrm{E} \cap \pi_{2}^{-1}(\{t\})\right)=0$.

Proof. - Fix $\varepsilon>0$ be arbitrarily small. We decompose $\mathrm{E}=\bigcup_{i \geq 1} \mathrm{E}_{i}$ as

$$
\mathrm{E}_{1}:=\left\{\left(x_{\circ}, t_{\circ}\right) \in \mathrm{E}: \omega_{x_{0}, t_{0}}(1)<\varepsilon\right\},
$$

$$
\mathrm{E}_{i}:=\left\{\left(x_{\mathrm{o}}, t_{\mathrm{o}}\right) \in \mathrm{E}: \omega_{x_{0}, t_{0}}\left(2^{-i+1}\right)<\varepsilon \leq \omega_{x_{0}, t_{0}}\left(2^{-i+2}\right)\right\}, \quad \text { for } i \geq 2 .
$$

Fix $i \geq 1$ and note that, if $\left(x_{1}, t_{1}\right)$ and $\left(x_{2}, t_{2}\right)$ belong to $\mathrm{E}_{i}$, then

$$
\begin{equation*}
\left\{(x, t) \in \mathrm{B}_{2^{-i+1}\left(x_{j}\right)} \times(-1,1): t-t_{j}>\varepsilon\left|x-x_{j}\right|^{s}\right\} \cap \mathrm{E}=\emptyset, \quad j=1,2 . \tag{7.8}
\end{equation*}
$$

Hence, by triangle inequality,

$$
\begin{equation*}
x_{1}, x_{2} \in \mathrm{E}_{i},\left|x_{1}-x_{2}\right| \leq 2^{-i} \quad \Rightarrow \quad\left|t_{1}-t_{2}\right| \leq \varepsilon\left|x_{1}-x_{2}\right|^{s} . \tag{7.9}
\end{equation*}
$$

In particular, since the sets $\left\{\mathrm{E}_{i}\right\}$ give a partition of E , it follows from (7.9) that the projection $\pi_{1}: \mathrm{E} \rightarrow \mathbf{R}^{n}$ is injective, and thus the sets $\pi_{1}\left(\mathrm{E}_{i}\right)$ are disjoint.

Now, by definition of $\mathcal{H}^{\beta}\left(\pi_{1}\left(\mathrm{E}_{i}\right)\right)$, there is countable collection of balls $\left\{\mathrm{B}_{\ell}\right\}$ such that $\pi_{1}\left(\mathrm{E}_{i}\right) \subset \bigcup_{\ell} \mathrm{B}_{\ell}$, with

$$
\begin{equation*}
\operatorname{diam}\left(\mathrm{B}_{\ell}\right)<2^{-i} \text { and } \sum_{\ell} \operatorname{diam}\left(\mathrm{B}_{\ell}\right)^{\beta}<\mathcal{H}^{\beta}\left(\pi_{1}\left(\mathrm{E}_{i}\right)\right)+2^{-i} \tag{7.10}
\end{equation*}
$$

Then, thanks to (7.9), we see that $\mathrm{E}_{i}$ can be covered by the family of cylinders

$$
\mathcal{F}_{i}:=\left\{\mathrm{C}_{\ell}:=\mathrm{B}_{\ell} \times\left(t_{\ell}-\varepsilon \operatorname{diam}\left(\mathrm{B}_{\ell}\right)^{s}, t_{\ell}+\varepsilon \operatorname{diam}\left(\mathrm{B}_{\ell}\right)^{s}\right)\right\}
$$

for some suitable $t_{\ell} \in(-1,1)$.
Let us show (a). Since $\left\{\pi_{2}\left(\mathrm{C}_{\ell}\right)\right\}$ is a covering of $\pi_{2}\left(\mathrm{E}_{i}\right)$ made of intervals of length $2 \varepsilon \operatorname{diam}\left(\mathrm{~B}_{\ell}\right)^{s}$, we have

$$
\mathcal{H}_{\infty}^{\beta / s}\left(\pi_{2}\left(\mathrm{E}_{i}\right)\right) \leq(2 \varepsilon)^{\beta / s} \sum_{\ell} \operatorname{diam}\left(\mathrm{B}_{\ell}\right)^{\beta} \leq(2 \varepsilon)^{\beta / s}\left(\mathcal{H}^{\beta}\left(\pi_{1}\left(\mathrm{E}_{i}\right)\right)+2^{-i}\right) .
$$

Summing over $i \geq 1$ we obtain

$$
\mathcal{H}_{\infty}^{\beta / s}\left(\pi_{2}(\mathrm{E})\right) \leq(2 \varepsilon)^{\beta / s}\left(\mathcal{H}^{\beta}\left(\pi_{1}(\mathrm{E})\right)+1\right)
$$

and (a) follows from the arbitrariness of $\varepsilon$.
To show (b), following the same notation as above, we define the function

$$
\begin{aligned}
\mathrm{N}_{i}(t, j)=\# & \left\{\mathrm{C}_{\ell} \in \mathcal{F}_{i}: \operatorname{diam}\left(\mathrm{B}_{\ell}\right) \in\left(2^{-j-1}, 2^{-j}\right)\right. \\
& \left.t \in\left(t_{\ell}-\varepsilon \operatorname{diam}\left(\mathrm{B}_{\ell}\right)^{s}, t_{\ell}+\varepsilon \operatorname{diam}\left(\mathrm{B}_{\ell}\right)^{s}\right)\right\}
\end{aligned}
$$

Let $\mathcal{I}_{i, j}$ denote the set of indices $\ell$ such that $\mathrm{C}_{\ell} \in \mathcal{F}_{i}$ and $\operatorname{diam}\left(\mathrm{B}_{\ell}\right) \in\left(2^{-j-1}, 2^{-j}\right)$. Then we can rewrite $\mathrm{N}_{i}(t, j)$ as

$$
\mathrm{N}_{i}(t, j)=\sum_{\ell \in \mathcal{I}_{i, j}} \chi_{\left(t_{\ell}-\varepsilon \operatorname{diam}\left(\mathbb{B}_{\ell}\right)^{s}, t_{\ell}+\varepsilon \operatorname{diam}\left(\mathrm{B}_{\ell}\right)^{s}\right)}(t) .
$$

Hence, integrating over $[-1,1]$ we get

$$
\int_{-1}^{1} \mathrm{~N}_{i}(t, j) d t \leq \sum_{\ell \in \mathcal{I}_{i, j}} 2 \varepsilon \operatorname{diam}\left(\mathrm{~B}_{\ell}\right)^{s} \leq 2 \varepsilon\left(2^{-j}\right)^{s} \# \mathcal{I}_{i, j},
$$

therefore, multiplying this estimate by $\left(2^{-j}\right)^{\beta-s}$ and summing over $j$, we obtain (recall (7.10))

$$
\begin{align*}
\int_{-1}^{1} \sum_{j}\left(2^{-j}\right)^{\beta-s} \mathrm{~N}_{i}(t, j) d t & =2 \varepsilon \sum_{j}\left(2^{-j}\right)^{\beta} \# \mathcal{I}_{i, j} \leq 2^{1+\beta} \varepsilon \sum_{\mathrm{C}_{\ell} \in \mathcal{F}_{i}} \operatorname{diam}\left(\mathrm{~B}_{\ell}\right)^{\beta}  \tag{7.11}\\
& \leq 2^{1+\beta} \varepsilon\left(\mathcal{H}^{\beta}\left(\pi_{1}\left(\mathrm{E}_{i}\right)\right)+2^{-i}\right)
\end{align*}
$$

We now consider the functions $f_{i, \varepsilon}(t):=\sum_{j}\left(2^{-j}\right)^{\beta-s} \mathrm{~N}_{i}(t, j)$ (note that the covering used to define $\mathrm{N}_{i}(t, j)$ depends on $\epsilon$ ), and $f_{\varepsilon}(t):=\sum_{i} f_{i, \varepsilon}(t)$. Then, summing (7.11) over $i$, we have

$$
\int_{-1}^{1} f_{\varepsilon}(t) d t \leq 2^{1+\beta} \varepsilon\left(\mathcal{H}^{\beta}\left(\pi_{1}(\mathrm{E})\right)+1\right)
$$

and it follows by Chebyshev inequality

$$
\begin{aligned}
& \mathcal{H}^{1}\left(\mathrm{X}^{\varepsilon}\right) \leq 2^{1+\beta} \varepsilon^{1 / 2}\left(\mathcal{H}^{\beta}\left(\pi_{1}(\mathrm{E})\right)+1\right) \\
& \quad \text { where } \mathrm{X}^{\varepsilon}:=\left\{t \in(-1,1): f_{\varepsilon}(t)>\varepsilon^{1 / 2}\right\} .
\end{aligned}
$$

Set $\mathrm{X}:=\cap_{\mathrm{M}=1}^{\infty} \mathrm{X}_{\mathrm{M}}$, where $\mathrm{X}_{\mathrm{M}}:=\bigcup_{k=\mathrm{M}}^{\infty} \mathrm{X}^{2^{-2 k}}$. Then

$$
\mathcal{H}^{1}\left(\mathrm{X}_{\mathrm{M}}\right) \leq \sum_{k=\mathrm{M}}^{\infty} 2^{1+\beta} 2^{-k}\left(\mathcal{H}^{\beta}\left(\pi_{1}(\mathrm{E})\right)+1\right) \leq 2^{1+\beta} 2^{1-\mathrm{M}}\left(\mathcal{H}^{\beta}\left(\pi_{1}(\mathrm{E})\right)+1\right)
$$

therefore $\mathcal{H}^{1}(\mathrm{X})=0$.
Also, for any $t \in[-1,1] \backslash \mathrm{X}$, there exists $\mathrm{M}_{t}$ such that $t \in[-1,1] \backslash \mathrm{X}_{\mathrm{M}} \subset[-1,1] \backslash$ $\mathrm{X}^{2^{-2 \mathrm{M}}}$ for any $\mathrm{M} \geq \mathrm{M}_{t}$. Therefore, considering the covering associated to $\varepsilon=2^{-2 \mathrm{M}}$, we get

$$
\begin{aligned}
& \mathcal{H}_{\infty}^{\beta-s}\left(\pi_{1}(\mathrm{E}) \cap \pi_{2}^{-1}(\{t\})\right) \\
& \quad \leq \sum_{i} \mathcal{H}_{\infty}^{\beta-s}\left(\pi_{1}\left(\mathrm{E}_{i}\right) \cap \pi_{2}^{-1}(\{t\})\right) \\
& \quad \leq \sum_{i} \sum_{j}\left(2^{-j}\right)^{\beta-s} \mathrm{~N}_{i}(t, j)=f_{2}-2 \mathrm{M}(t) \leq 2^{-\mathrm{M}} \quad \forall \mathrm{M} \geq \mathrm{M}_{t} .
\end{aligned}
$$

This proves that $\mathcal{H}_{\infty}^{\beta-s}\left(\pi_{1}(\mathrm{E}) \cap \pi_{2}^{-1}(\{t\})\right)=0$ for all $t \in[-1,1] \backslash \mathrm{X}$, as wanted.

As an immediate consequence of Proposition 7.7, we get:
Corollary 7.8. - Let $\mathrm{E} \subset \mathbf{R}^{n} \times[-1,1]$, let $(x, t)$ denote a point in $\mathbf{R}^{n} \times[-1,1]$, and let $\pi_{1}:(x, t) \mapsto x$ and $\pi_{2}:(x, t) \mapsto t$ be the standard projections. Assume that, for some $\beta \in(0, n]$ and $s>0$, we have:

- $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}(\mathrm{E})\right) \leq \beta$;
- For all $\left(x_{0}, t_{0}\right) \in \mathrm{E}$ and $\varepsilon>0$, there exists $\varrho=\varrho_{x_{0}, t_{0}, \varepsilon}>0$ such that

$$
\left\{(x, t) \in \mathrm{B}_{\varrho}\left(x_{0}\right) \times[-1,1]: t-t_{0}>\left|x-x_{0}\right|^{s-\varepsilon}\right\} \cap \mathrm{E}=\emptyset
$$

Then:
(a) If $\beta<s$, we have $\operatorname{dim}_{\mathcal{H}}\left(\pi_{2}(\mathrm{E})\right) \leq \beta / s$.
(b) If $\beta \geq s$, for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$ we have $\operatorname{dim}_{\mathcal{H}}\left(\mathrm{E} \cap \pi_{2}^{-1}(\{t\})\right) \leq \beta-s$.

## 8. Dimension reduction results

This section is concerned with bounding the Hausdorff dimension of the differences of the subsets of $\boldsymbol{\Sigma}_{n-1}$ defined in (6.2). Note that we have the chain of inclusions

$$
\begin{equation*}
\boldsymbol{\Sigma} \supset \boldsymbol{\Sigma}_{n-1} \supset \boldsymbol{\Sigma}_{n-1}^{\geq 3} \supset \boldsymbol{\Sigma}_{n-1}^{3 r d}=\boldsymbol{\Sigma}_{n-1}^{>3} \supset \boldsymbol{\Sigma}_{n-1}^{\geq 4} \supset \boldsymbol{\Sigma}_{n-1}^{4 t h}=\boldsymbol{\Sigma}_{n-1}^{>4} \supset \boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta} \tag{8.1}
\end{equation*}
$$

where the two equalities in such chain of inclusions follow from Propositions 4.15 and 5.8.
For $0 \leq m \leq n-2$, we simply consider the sets

$$
\boldsymbol{\Sigma}_{m} \supset \boldsymbol{\Sigma}_{m}^{\geq 3}=\boldsymbol{\Sigma}_{m}^{3 r d}, \quad 0 \leq m \leq n-2,
$$

as this suffices for our purposes. Recall that, by Proposition 3.9(a), we have $\boldsymbol{\Sigma}_{m} \backslash \boldsymbol{\Sigma}_{m}^{a}=$ $\boldsymbol{\Sigma}_{m}^{\geq 3}=\boldsymbol{\Sigma}_{m}^{3 r d}$.

Our goal is to show that $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta}\right)\right) \leq n-2$ for any $\zeta \in(0,1)$, where $\pi_{1}$ denotes the canonical projection $\pi_{1}:(x, t) \mapsto x$. For this, using the tools developed in the previous sections, in the next lemmas we bound the size all the differences between consecutive sets of the previous chain of inclusions.

Proposition 8.1. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1). Then:
(a) $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{m}^{a}\right)\right) \leq m-1$ for $1 \leq m \leq n-2\left(\pi_{1}\left(\boldsymbol{\Sigma}_{m}^{a}\right)\right.$ is discrete if $\left.m=1\right)$.
(b) $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{<3}\right)\right) \leq n-3\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{<3}\right)\right.$ is countable if $\left.n=3\right)$.
(c) For any $\varrho \in(0,1)$ :

- if $m \leq n-2$ then $\pi_{1}\left(\boldsymbol{\Sigma}_{m} \backslash \boldsymbol{\Sigma}_{m}^{a}\right) \cap \overline{\mathbf{B}_{\rho}}$ is covered by a $\mathrm{C}^{1,1}$ m-dimensional manifold;
$-\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \cap \overline{\mathrm{B}_{\rho}}$ is covered by a $\mathrm{C}^{1,1}(n-1)$-dimensional manifold.

Proof. - (a) We need to prove that, for any $\beta>m-1$, the set $\pi_{1}\left(\boldsymbol{\Sigma}_{m}^{a}\right)$ has zero $\mathcal{H}^{\beta}$ measure. Assume by contradiction that

$$
\mathcal{H}^{\beta}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{m}^{a}\right)\right)>0
$$

Then, by Lemma 7.6, there exists a point $\left(x_{0}, t_{0}\right) \in \boldsymbol{\Sigma}_{m}^{a}$-which we assume for simplicity to be $\left(x_{\circ}, t_{\mathrm{o}}\right)=(0,0)$-, a sequence $r_{k} \downarrow 0$, and a set $\mathcal{A} \subset \overline{\mathrm{B}_{1}}$ with $\mathcal{H}^{\beta}(\mathcal{A})>0$, such that for every point $y \in \mathcal{A}$ there is a sequence $\left(x_{k}, t_{k}\right)$ in $\boldsymbol{\Sigma}_{m}^{a}$ such that $x_{k} / r_{k} \rightarrow y$.

Let $w=u(\cdot, 0)-p_{2}, w_{r}=w(r \cdot), \tilde{w}_{r}=w_{r} / \mathrm{H}\left(1, w_{r}\right)^{1 / 2}$. Then, thanks to Proposition 3.9, up to extracting a subsequence we have

$$
\begin{equation*}
\tilde{w}_{r_{k}} \rightharpoonup q \quad \text { in } \mathrm{W}_{\mathrm{loc}}^{1,2}\left(\mathbf{R}^{n}\right) \tag{8.2}
\end{equation*}
$$

where $q$ is $\lambda^{2 n d}$-homogeneous harmonic function. By definition of $\boldsymbol{\Sigma}_{m}^{a}$ we have $\lambda^{2 n d}=2$, and thus $q$ is a quadratic harmonic polynomial satisfying (3.11).

Thanks to Lemma 6.4 we have $\mathcal{A} \subset\{q=0\} \cap\left\{p_{2}=0\right\}$. Therefore, since $\mathcal{H}^{\beta}(\mathcal{A})>0$, the polynomial $q$ vanishes in a subset of dimension $\beta>m-1$ of the $m$ dimensional linear space $\left\{p_{2}=0\right\}$. The only possibility is that $q \equiv 0$ on $\left\{p_{2}=0\right\}$, and then (3.11) implies $q \equiv 0$; a contradiction since $\mathrm{H}(1, q)=1$.

We note that in the case $m=1$ the same proof gives that $\boldsymbol{\Sigma}_{1}^{a}$ cannot have accumulation points, i.e., it must be a discrete set.
(b) We apply Proposition 7.3 to the set $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{<3}\right)$ with the function $f: \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{<3}\right) \rightarrow$ $[0, \infty)$ defined by

$$
\begin{aligned}
f\left(x_{\circ}\right) & :=\phi\left(0^{+}, u\left(\cdot, \tau\left(x_{\circ}\right)\right)-p_{2, x_{0}, \tau\left(x_{0}\right)}\right), \\
\quad \text { with } \tau\left(x_{\circ}\right) & :=\min \left\{t \in[-1,1]:\left(x_{\circ}, t\right) \in \boldsymbol{\Sigma}\right\} .
\end{aligned}
$$

Note that, by Lemma 6.2(c), we have $\phi\left(0^{+}, u(\cdot, t)-p_{2, x_{0}, t}\right)=f\left(x_{0}\right)$ for every $t$ such that $\left(x_{\circ}, t\right) \in \boldsymbol{\Sigma}$. Also, by Proposition 3.9 (b) and the definition of $\boldsymbol{\Sigma}_{n-1}^{<3}$, we have $f\left(x_{0}\right) \in[2+$ $\left.\alpha_{o}, 3\right)$.

To obtain the result, thanks to Proposition 7.3, it suffices to show the following property: for all $x_{\circ} \in \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{<3}\right)$ and for all $\varepsilon>0$ there exists $\varrho=\varrho\left(x_{0}, \varepsilon\right)>0$ such that

$$
\begin{aligned}
\mathrm{B}_{r}\left(x_{\circ}\right) & \cap \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{<3}\right) \cap f^{-1}\left(\left[f\left(x_{\circ}\right)-\varrho, f\left(x_{\circ}\right)+\varrho\right]\right) \\
& \subset\left\{y: \operatorname{dist}\left(y, \Pi_{x_{0}, r}\right) \leq \varepsilon r\right\} \quad \forall r \in(0, \varrho)
\end{aligned}
$$

where $\Pi_{x_{0}, r}$ is a $(n-3)$-dimensional plane passing through $x_{\circ}$.
With no loss of generality we can assume that $\left(x_{\mathrm{o}}, t_{\mathrm{o}}\right)=(0,0)$, and we prove this statement by contradiction. If such $\varrho>0$ did not exist for some $\varepsilon>0$, then we would have sequences $r_{k} \downarrow 0$ and $x_{k}^{(j)} \in \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{<3}\right) \cap \mathbf{B}_{r_{k}}, 1 \leq j \leq n-2$, such that

$$
\begin{aligned}
& y_{k}^{(j)}:=x_{k}^{(j)} / r_{k} \rightarrow y_{\infty}^{(j)} \in \overline{\mathrm{B}_{1}}, \quad \operatorname{dim}\left(\operatorname{span}\left(y_{\infty}^{(1)}, y_{\infty}^{(2)}, \ldots, y_{\infty}^{(n-2)}\right)\right)=n-2, \\
& \left|f\left(x_{k}^{(j)}\right)-f(0)\right| \downarrow 0 .
\end{aligned}
$$

Let $w=u(\cdot, 0)-p_{2}, w_{r}=w(r \cdot), \tilde{w}_{r}=w_{r} / \mathrm{H}\left(1, w_{r}\right)^{1 / 2}$. It follows by Proposition 3.9 that (8.2) holds, where $q$ is a $\lambda^{2 n d}$-homogeneous solution of the Signorini problem (3.12). Also, since we are supposing that $(0,0) \in \boldsymbol{\Sigma}_{n-1}^{<3}$, we have $\lambda^{2 n d} \in\left[2+\alpha_{\circ}, 3\right)$.

Applying then Lemma 6.5 to the sequences $\left(x_{k}^{(j)}, \tau\left(x_{k}^{(j)}\right)\right)$ we deduce that $q$ is translation invariant in the $n-2$ independent directions

$$
y_{\infty}^{(1)}, y_{\infty}^{(2)}, \ldots, y_{\infty}^{(n-2)} \in\left\{p_{2}=0\right\} .
$$

As a consequence $q$ is a two dimensional $\lambda^{2 n d}$-homogeneous solution of Signorini, with $\lambda^{2 n d} \in\left[2+\alpha_{\circ}, 3\right)$. However, it follows from Lemma A. 2 that 2D homogeneous solutions of Signorini have homogeneities $\{1,2,3,4, \ldots\} \cup\left\{1+\frac{1}{2}, 3+\frac{1}{2}, 5+\frac{1}{2}, 7,+\frac{1}{2}, \ldots\right\}$, impossible.

Note finally that, when $n=3$, the same argument (but using Lemma 7.1 in place of Proposition 7.3) implies that $\boldsymbol{\Sigma}_{n-1}^{<3}$ is at most countable.
(c) We prove the statement for the maximal stratum $\boldsymbol{\Sigma}_{n-1}^{\geq 3}$; the proof for $\boldsymbol{\Sigma}_{m} \backslash \boldsymbol{\Sigma}_{m}^{a}=$ $\Sigma_{m}^{\geq 3}$ is analogous.

Given $x_{\circ} \in \underline{\pi_{1}}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right)$, set $\mathrm{P}_{x_{0}}:=p_{2, x_{0}, \tau\left(x_{0}\right)}\left(\cdot-x_{0}\right)$. We claim that, for every pair $x_{0}, x \in \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \cap \overline{\mathrm{B}_{\varrho}}$, we have
(8.3)

$$
\left|\mathrm{D}^{k} \mathrm{P}_{x_{0}}(x)-\mathrm{D}^{k} \mathrm{P}_{x}(x)\right| \leq \mathrm{C}\left|x-x_{0}\right|^{3-k} \quad \text { for } k=0,1,2 .
$$

Indeed, note that for all $\hat{x} \in \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \cap \overline{\mathbf{B}_{\varrho}}$ we have $\phi\left(0^{+} u(\hat{x}+\cdot, \tau(\hat{x}))-p_{2, \hat{x}, \tau(\hat{x})}\right) \geq 3$. Thus, by Lemma 3.6,

$$
\left\|u(\hat{x}+\cdot, \tau(\hat{x}))-p_{2, \hat{x}, \tau(\hat{x})}\right\|_{L^{\infty}\left(\mathbf{B}_{r}\right)} \leq \mathrm{C}(n, \varrho) r^{3} \quad \forall r \in\left(0, \frac{1-\varrho}{2}\right)
$$

therefore, applying this bound both to $\hat{x}=x_{\circ}$ and $\hat{x}=x$, we get

$$
\begin{aligned}
& \left|u\left(\cdot, \tau\left(x_{0}\right)\right)-\mathrm{P}_{x_{0}}\right| \leq \mathrm{C}^{3} \quad \text { in } \mathrm{B}_{r}\left(x_{0}\right) \quad \text { and } \\
& \left|u(\cdot, \tau(x))-\mathrm{P}_{x}\right| \leq \mathrm{C}^{3} \quad \text { in } \mathrm{B}_{r}(x) .
\end{aligned}
$$

Choosing $r=2\left|x-x_{0}\right|$, and assuming without loss of generality that $\tau\left(x_{0}\right) \leq \tau(x)$, since $u\left(\cdot, \tau\left(x_{0}\right)\right) \leq u(\cdot, \tau(x))$ we obtain

$$
\mathrm{P}_{x_{0}}-\mathrm{P}_{x} \leq \mathrm{C}^{3}+u\left(\cdot, \tau\left(x_{\circ}\right)\right)-u(\cdot, \tau(x)) \leq \mathrm{Cr}^{3} \quad \text { in } \mathrm{B}_{r}\left(x_{\circ}\right) \cap \mathrm{B}_{r}(x) .
$$

Noticing that $P_{x_{0}}-P_{x}$ is a harmonic quadratic polynomial that vanishes at some point $\hat{x}$ in the segment joining $x_{\circ}$ to $x$, as a consequence of the above upper bound we easily deduce that

$$
\left\|\mathrm{P}_{x_{0}}-\mathrm{P}_{x}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{4 r}(\hat{x})\right)} \leq \mathrm{Cr} r^{3},
$$

and since the $L^{\infty}\left(B_{1}\right)$ and the $C^{3}\left(B_{1}\right)$ norm are equivalent on space of quadratic polynomials, (8.3) holds.

Then, applying Whitney's extension theorem (see [Fef09] or [FS19, Lemma 3.10]) we obtain a $\mathrm{C}^{2,1}$ function $\mathrm{F}: \mathrm{B}_{1} \rightarrow \mathbf{R}$ satisfying

$$
\mathrm{F}(x)=\mathrm{P}_{x_{0}}(x)+\mathrm{O}\left(\left|x_{0}-x\right|^{3}\right)
$$

for all $x_{\circ} \in \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \cap \overline{\mathbf{B}_{\varrho}}$. In particular $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \subset\{\nabla \mathrm{F}=0\}$ and $\mathrm{D}^{2} \mathrm{~F}\left(x_{\circ}\right)=$ $\mathrm{D}^{2} p_{2, x_{0}, \tau\left(x_{0}\right)}(0)$ has rank one (recall that $\left.\left(x_{0}, \tau\left(x_{0}\right)\right) \in \boldsymbol{\Sigma}_{n-1}\right)$. Hence, by the implicit function theorem, we find that $\{\nabla \mathrm{F}=0\}$ is a $\mathrm{C}^{1,1}(n-1)$-dimensional manifold in a neighborhood of $x_{0}$.

As a consequence of the previous result, we get the following:
Corollary 8.2. - Let $n=3$, let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1), and assume that $u\left(x, t^{\prime}\right)>u(x, t)$ whenever $t^{\prime}>t$ and $u(x, t)>0$. Then, for all but a countable set of singular points ( $x_{\circ}, t_{\circ}$ ), we have

$$
\left\|u\left(x_{\mathrm{o}}+\cdot, t_{\mathrm{o}}\right)-p_{2, x_{0}, t_{0}}\right\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{r}\right)} \leq \mathrm{C} r^{3} \quad \forall r \in\left(0, \frac{1-\left|x_{\mathrm{o}}\right|}{2}\right)
$$

where C depends only on $n$ and $1-\left|x_{0}\right|$.
Proof. - On the one hand, since $n=3$, Proposition 8.1 implies that $\boldsymbol{\Sigma}_{m} \backslash \boldsymbol{\Sigma}_{m}^{\geq 3}$ is a countable set for $m=0,1,2 \cdot{ }^{8}$ On the other hand, for $\left(x_{0}, t_{0}\right) \in \boldsymbol{\Sigma}_{m}^{\geq 3}$, setting $\rho=\frac{\boldsymbol{I}_{m}-\left|x_{0}\right|}{2}$ and applying Lemma 3.6 to the function $w=\rho^{-2} u\left(x_{0}+\rho \cdot, t_{0}\right)-p_{2, x_{0}, t_{0}}$ (note that then $\left.\phi\left(0^{+}, w\right) \geq 3\right)$ we obtain

$$
\left(\frac{\rho}{r}\right)^{6} \leq \frac{\mathrm{H}(w, \rho)}{\mathrm{H}(w, r)}
$$

Therefore, using Lemma 3.7, we obtain

$$
\|w\|_{\mathrm{L}^{\infty}\left(\mathrm{B}_{\mathrm{r}}\right)} \leq \mathrm{C}(n) \mathrm{H}(w, 2 r)^{1 / 2} \leq \mathrm{C}(n) \frac{\mathrm{H}(w, \rho)^{1 / 2}}{\rho^{3}} r^{3},
$$

as desired.
Proposition 8.3. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1). Then $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 n d}\right)$ is contained in a countable union of $(n-2)$-dimensional Lipschitz manifolds.

Proof. - For any $\left(x_{\circ}, t_{\circ}\right) \in \boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}$ we apply Lemma 6.6 to $u\left(x_{\circ}+\cdot, t_{\circ}+\cdot\right)$ to find a $(n-2)$-dimensional linear subspace $\mathrm{L}_{x_{0}, t_{0}}$ and $\varrho_{x_{0}, t_{0}}>0$ such that

$$
\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \cap \mathrm{B}_{r}\left(x_{\circ}\right) \subset x_{\circ}+\mathrm{L}_{x_{0}, t_{0}}+\mathrm{B}_{r} \quad \text { for all } r \in\left(0, \varrho_{x_{0}, t_{0}}\right) .
$$

[^8]Write $\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}=\bigcup_{j} \mathrm{E}_{j}$, where

$$
\mathrm{E}_{j}:=\left\{\left(x_{\circ}, t_{\circ}\right) \in \boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}: \varrho_{x_{0}, t_{o}}>1 / j\right\} .
$$

Note that, for any $\left(x_{\circ}, t_{\circ}\right) \in \mathrm{E}_{j}$, the set $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3}\right) \cap \mathrm{B}_{1 / j}\left(x_{\circ}\right)$ is contained inside the cone

$$
\left\{x \in \mathrm{~B}_{1 / j}\left(x_{\circ}\right): \operatorname{dist}\left(\frac{x-x_{\circ}}{\left|x-x_{0}\right|}, x_{\circ}+\mathrm{L}_{x_{0}, t_{0}}\right) \leq 1\right\},
$$

which implies (by a classical geometric argument) that the set $\pi_{1}\left(\mathrm{E}_{j}\right) \cap \mathrm{B}_{1 / 2}$ can be covered by a 1-Lipschitz $(n-2)$-dimensional manifold. The result follows by taking the union of these manifolds over all $j \in \mathbf{N}$.

Lemma 8.4. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1). Then:
(a) $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right)\right) \leq n-2$ (countable if $n=2$ ).
(b) $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>4} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta}\right)\right) \leq n-3$ (countable if $n=3$ ).

Proof. - (a) The proof is similar to the one of Proposition 8.1(b). Indeed, we apply Proposition 7.3 to the set $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right)$ with the function $f: \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right) \rightarrow[0, \infty)$ defined as

$$
\begin{align*}
f\left(x_{0}\right) & :=\phi\left(0^{+}, u\left(\cdot, \tau\left(x_{\circ}\right)\right)-\mathscr{P}_{x_{0}, \tau\left(x_{0}\right)}\right)  \tag{8.4}\\
\quad \text { where } \tau\left(x_{0}\right) & :=\min \left\{t \in[-1,1]:\left(x_{\circ}, t\right) \in \boldsymbol{\Sigma}\right\} .
\end{align*}
$$

By Lemma 6.2 (c) we have $\phi\left(0^{+}, u(\cdot, t)-\mathscr{P}_{x_{0}, t}\right)=f\left(x_{\mathrm{o}}\right)$ for every $t$ such that $\left(x_{\mathrm{o}}, t\right) \in \boldsymbol{\Sigma}$. Moreover, by definition of $\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}$, we have $f\left(x_{0}\right) \in(3,4)$. Then, thanks to Proposition 7.3, it is enough to show that for all $x_{0} \in \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right)$ and for all $\varepsilon>0$ there exist $\varrho=\varrho\left(x_{0}, \varepsilon\right)>0$, and a $(n-2)$-dimensional plane $\Pi_{x_{0}}$ passing through $x_{0}$, such that

$$
\begin{aligned}
\mathrm{B}_{r}\left(x_{\circ}\right) & \cap \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right) \cap f^{-1}\left(\left[f\left(x_{\circ}\right)-\varrho, f\left(x_{\circ}\right)+\varrho\right]\right) \\
& \subset\left\{y: \operatorname{dist}\left(y, \Pi_{x_{0}}\right) \leq \varepsilon r\right\} \quad \forall r \in(0, \varrho) .
\end{aligned}
$$

Assuming $\left(x_{0}, t_{0}\right)=(0,0)$ and arguing by contradiction, we find sequences $r_{k} \downarrow 0$ and $x_{k}^{(j)} \in \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right) \cap \mathrm{B}_{r_{k}}, 1 \leq j \leq n-1$, such that

$$
\begin{aligned}
& y_{k}^{(j)}:=x_{k}^{(j)} / r_{k} \rightarrow y_{\infty}^{(j)} \in \overline{\mathrm{B}_{1}}, \quad \operatorname{dim}\left(\operatorname{span}\left(y_{\infty}^{(1)}, y_{\infty}^{(2)}, \ldots, y_{\infty}^{(n-1)}\right)\right)=n-1, \\
& \left|f\left(x_{k}^{(j)}\right)-f(0)\right| \downarrow 0 .
\end{aligned}
$$

Setting $w=u(\cdot, 0)-\mathscr{P}, w_{r}=w(r \cdot), \tilde{w}_{r}=w_{r} / \mathrm{H}\left(1, w_{r}\right)^{1 / 2}$, it follows by Proposition 4.12(a) that (8.2) holds, where $q$ is a $\lambda^{3 r d}$-homogeneous solution of the Signorini problem (3.12) with $\lambda^{3 r d} \in(3,4)$ (recall that $(0,0) \in \boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}$ ). Also, applying Lemma 6.8
to the sequences $\left(x_{k}^{(j)}, \tau\left(x_{k}^{(j)}\right)\right)$, we deduce that $q$ is translation invariant in the $n-1$ independent directions

$$
y_{\infty}^{(1)}, y_{\infty}^{(2)}, \ldots, y_{\infty}^{(n-1)} \in\left\{p_{2}=0\right\} .
$$

Thus $q$ is a 1D $\lambda^{3 r d}$-homogeneous solution of Signorini, with $\lambda^{3 r d} \in(3,4)$, and this is impossible by Lemma A.1.

Finally, when $n=2$, the same argument (using Lemma 7.1 instead of Proposition 7.3) implies that $\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}$ is at most countable.
(b) The proof is completely analogous to the one of part (a), using Lemmas 6.10 and A. 2 instead of Lemmas 6.8 and A.1.

Remark 8.5. - Notice that the difference between parts (a) and (b) in the previous Lemma comes from the fact that there exist 2D solutions to the Signorini problem with homogeneity $3+\frac{1}{2} \in(3,4)$, while there is no such solution with homogeneity in the interval $(4,5)$. Hence, using the exact same proof as above, one can show that $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 7 / 2}\right)\right) \leq n-3$, where we define $\boldsymbol{\Sigma}_{n-1}^{\geq 7 / 2}$ as the set at which $\phi\left(0^{+}, u-\mathscr{P}\right) \geq 7 / 2$.

With the aid of Lemmas 7.4 and 7.5, we can next prove the following:
Lemma 8.6. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathrm{~B}_{1}} \times[-1,1]\right)$ solve (6.1). Then

$$
\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 4} \backslash \boldsymbol{\Sigma}_{n-1}^{4 t / h}\right)\right) \leq n-2 .
$$

Proof. - Define $\tau: \pi_{1}(\boldsymbol{\Sigma}) \rightarrow[-1,1]$ as in (8.4) and note that, by Lemma 6.2, it is lower semicontinuous.

Hence, thanks to Lemma 7.5, it suffices to prove that, for any given $\varepsilon>0$ and $\left(x_{0}, \tau\left(x_{0}\right)\right) \in \boldsymbol{\Sigma}_{n-1}^{\geq 4} \backslash \boldsymbol{\Sigma}_{n-1}^{4 \text { th }}$, there exists $\varrho=\varrho\left(x_{0}, \varepsilon\right)>0$ such that
$(\mathbf{8 . 5}) \quad \boldsymbol{\Sigma} \cap \overline{\mathbf{B}_{r}\left(x_{0}\right)} \times\left[\tau\left(x_{\circ}\right), 1\right) \subset\left\{x_{\circ}+r y: y \in \mathbf{S}\left(p_{x_{0}, r}, \varepsilon\right)\right\} \quad \forall r \in(0, \varrho)$,
for some $p_{x_{0}, r} \in \mathcal{P}_{4, \geq}^{\text {even }}$. This follows from Lemma 6.12 applied to $u\left(x_{\circ}+\cdot, \tau\left(x_{0}\right)\right)$, since by monotonicity

$$
\boldsymbol{\Sigma} \cap \pi_{2}^{-1}\left(\left[\tau\left(x_{0}\right), 1\right]\right) \subset\left\{u\left(x_{\circ}+\cdot, \tau\left(x_{\circ}\right)\right)=0\right\} .
$$

We can finally prove the following:
Theorem 8.7. - Let $u \in \mathrm{C}^{0}\left(\overline{\mathbf{B}_{1}} \times[-1,1]\right)$ solve (6.1). There exists a set $\boldsymbol{\Sigma}^{*} \subset \boldsymbol{\Sigma}_{n-1} \subset \boldsymbol{\Sigma}$, with $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma} \backslash \mathbf{\Sigma}^{*}\right)\right) \leq n-2$, such that for any given $\varepsilon>0$ the following holds:

$$
\begin{aligned}
& \left\|u\left(x_{\circ}+\cdot, t_{\circ}\right)-\mathscr{P}_{x_{\circ}, t_{0}}-p_{4, x_{o}, t_{0}}\right\|_{L^{\infty}\left(\mathbf{B}_{r}\right)} \leq \mathrm{C} r^{5-\varepsilon} \\
& \quad \forall r \in\left(0, \frac{1}{2}\right), \forall\left(x_{\circ}, t_{\circ}\right) \in\left(\boldsymbol{\Sigma}^{*} \cap \mathrm{~B}_{1 / 2}\right) \times(-1,1),
\end{aligned}
$$

where C depends only on $n$ and $\varepsilon$.

Proof. - Recall the chain of inclusions (8.1). We have:

- Proposition 8.1 (a) and (c) $\Rightarrow \operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma} \backslash \boldsymbol{\Sigma}_{n-1}\right) \quad\right) \leq n-2$,
- Proposition 8.1 (b) $\Rightarrow \operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 3}\right)\right) \leq n-3$,
- Proposition $8.3 \Rightarrow \operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 n d}\right)\right) \leq n-2$,
- Remark 6.1 $\quad \Rightarrow \quad \pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{3 r d} \backslash \boldsymbol{\Sigma}_{n-1}^{>3}\right)=\emptyset$,
- Lemma 8.4(a) $\Rightarrow \operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 4}\right)\right) \leq n-2$,
- Lemma 8.6 $\Rightarrow \operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 4} \backslash \boldsymbol{\Sigma}_{n-1}^{>4}\right)\right) \leq n-2$,
- Lemma 8.4(b) $\Rightarrow \operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>4} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 5-\zeta}\right)\right) \leq n-3$.

Thus, if we define

$$
\boldsymbol{\Sigma}^{*}:=\bigcap_{\varepsilon>0} \boldsymbol{\Sigma}_{n-1}^{\geq 5-\varepsilon},
$$

then $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma} \backslash \boldsymbol{\Sigma}^{*}\right)\right) \leq n-2$. Fix $\varepsilon>0$, and let $\left(x_{0}, t_{0}\right) \in\left(\boldsymbol{\Sigma}^{*} \cap \mathrm{~B}_{1 / 2}\right) \times(-1,1)$. By Lemmas 4.9 and 4.1 applied to $w:=u\left(x_{0}+\cdot, t_{0}\right)-\mathscr{P}_{x_{0}, t_{0}}-p_{4, x_{0}, t_{0}}$ we obtain

$$
c\left(\frac{1}{r}\right)^{2(5-\varepsilon)} \leq \frac{\mathrm{H}(1 / 2, w)+(1 / 2)^{2(5-\varepsilon)}}{\mathrm{H}(r, w)+r^{2(5-\varepsilon)}}
$$

therefore

$$
\begin{aligned}
\mathrm{H}(r, w)^{1 / 2} & \leq \mathrm{C}\left(\int_{\mathbf{B}_{1 / 2}}\left(u\left(x_{0}+\cdot, t_{0}\right)-\mathscr{P}_{x_{0}, t_{0}}-p_{4, x_{0}, t_{0}}\right)^{2}+(1 / 2)^{2(5-\varepsilon)}\right)^{1 / 2} r^{5-\varepsilon} \\
& \leq \mathrm{C}(n, \varepsilon) r^{5-\varepsilon} .
\end{aligned}
$$

Combining this bound with the Lipschitz estimate in Lemma 4.7, we easily conclude that

$$
\left\|u\left(x_{\circ}+\cdot, t_{\circ}\right)-\mathscr{P}_{x_{0}, t_{0}}-p_{4, x_{0}, t_{0}}\right\|_{L^{\infty}\left(\mathbf{B}_{r}\right)}=\|w\|_{\mathrm{L}^{\infty}\left(\mathbf{B}_{r}\right)} \leq \mathrm{C} r^{5-\varepsilon} \quad \forall 0<r<1 / 2,
$$

where C depends only on $n$ and $\varepsilon$.

## 9. Cleaning lemmas and proof of the main results

Recall that, in all the previous sections, we only assumed that $u(\cdot, t)$ was nondecreasing in $t$. Now, in order to conclude the proof of Theorem 1.1, we will assume the "uniform monotonicity" condition (1.2). Note that condition (1.2) rules out the existence of connected components of the complement of the contact set that remain unchanged for some interval of times.

The first bound involves a barrier argument that will play an important role.

Lemma 9.1. - Let $u: \mathrm{B}_{1} \times(-1,1) \rightarrow[0, \infty)$ solve $(6.1)$, with $(0,0) \in \boldsymbol{\Sigma}$ and $\left\{p_{2}=\right.$ $0\} \subset\left\{x_{n}=0\right\}$. Let $\mathfrak{p}$ be a polynomial satisfying $\Delta \mathfrak{p}=1$. Assume that, for some $\beta \geq 0$, we have

$$
|u(\cdot, 0)-\mathfrak{p}| \leq \mathrm{C} r^{\beta} \quad \text { in } \mathrm{B}_{r} \forall r \in\left(0, r_{\mathrm{o}}\right),
$$

and define

$$
\begin{aligned}
& \psi(x):=-\sum_{i=1}^{n-1} x_{i}^{2}+(n-1) x_{n}^{2}+\frac{1}{2}, \quad \psi^{r}(x):=\psi(x / r), \\
& \mathrm{D}_{r}:=\partial \mathbf{B}_{r} \cap\left\{\psi^{r}>0\right\}=\partial \mathbf{B}_{r} \cap\left\{\left|x_{n}\right|>\frac{r}{\sqrt{2 n}}\right\} .
\end{aligned}
$$

Then, for all $t \geq 0$ we have

$$
u(\cdot, t) \geq \mathfrak{p}+\frac{\min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)]}{\max _{\partial \mathrm{B}_{1}} \psi} \psi^{r}-\mathrm{C}^{\beta} \quad \text { in } \mathbf{B}_{r} \forall r \in\left(0, r_{\circ}\right) .
$$

Proof. - It follows by our assumption on $u$ that

$$
\begin{equation*}
u(\cdot, 0)-\mathfrak{p} \geq-\mathrm{C} r^{\beta} \quad \forall r \in\left(0, \frac{1}{2}\right) \tag{9.1}
\end{equation*}
$$

Set

$$
v:=\mathfrak{p}+\mathrm{M} \psi^{r}-\mathrm{C} r^{\beta}, \quad \text { with } \mathrm{M}:=\frac{\min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)]}{\max _{\partial \mathrm{B}_{1}} \psi} .
$$

We claim that $v \leq u(\cdot, t)$ on $\partial \mathbf{B}_{r}$. Indeed, since $t \geq 0$, it follows by (9.1) that

$$
v \leq \mathfrak{p}-\mathrm{C} r^{\beta} \leq u(\cdot, 0) \leq u(\cdot, t) \quad \text { on } \partial \mathbf{B}_{r} \cap\left\{\psi^{r} \leq 0\right\}
$$

On the other hand, since $\max _{\partial \mathrm{B}_{1}} \psi=\max _{\mathrm{D}_{r}} \psi^{r}$, we see that $\mathrm{M} \psi^{r} \leq \min _{\mathrm{D}_{r}}[u(\cdot, r)-$ $u(\cdot, 0)$ ] on $\partial \mathrm{B}_{r}$. Hence

$$
v=\mathfrak{p}+\mathrm{M} \psi^{r}-\mathrm{C} r^{\beta} \leq u(\cdot, 0)+\mathrm{M} \psi^{r} \leq u(\cdot, t) \quad \text { on } \mathrm{D}_{r}=\partial \mathrm{B}_{r} \cap\left\{\psi^{r}>0\right\}
$$

and the claim follows.
To conclude the proof it suffices to observe that, since $\psi^{r}$ is harmonic, we have $\Delta v=1 \geq \chi_{\{u(\cdot, t)>0\}}=\Delta u(\cdot, t)$. Thus, combining the claim with the maximum principle, we conclude that

$$
v \leq u(\cdot, t) \quad \text { in } \mathrm{B}_{r} .
$$

The second result gives us a bound on the speed at which $u$ increases in $t$ at singular points. Note that this speed is much better in the lower strata $\boldsymbol{\Sigma}_{m}$ with $m \leq n-2$ with respect to $\boldsymbol{\Sigma}_{n-1}$. This is one of the reasons why, in the previous sections, we needed to perform a very refined analysis at points in $\boldsymbol{\Sigma}_{n-1}$.

Lemma 9.2. - Let $u: \mathrm{B}_{1} \times(-1,1) \rightarrow[0, \infty)$ satisfy (6.1) and (1.2), with $(0,0) \in \boldsymbol{\Sigma}$ and $\left\{p_{2}=0\right\} \subset\left\{x_{n}=0\right\}$. Let $\mathrm{D}_{r}$ be defined as in Lemma 9.1.
(a) If $(0,0) \in \boldsymbol{\Sigma}_{m}$ with $m \leq n-2$, then for all $\varepsilon>0$ there exist $c_{\varepsilon}, \rho_{\varepsilon}>0$ such that

$$
\min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \geq c_{\varepsilon} r^{\varepsilon} t, \quad \forall r \in\left(0, \rho_{\varepsilon}\right) .
$$

(b) If $(0,0) \in \mathbf{\Sigma}_{n-1}$, there exists $c, \rho>0$ such that

$$
\min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \geq c r t, \quad \forall r \in(0, \rho)
$$

Proof. - (a) Note that, by the uniform convergence of $r^{-2} u(r \cdot, 0)$ to $p_{2}$, given $\delta>0$ there exists $r_{\delta}>0$ such that

$$
\{u(\cdot, 0)=0\} \cap \mathrm{B}_{r_{\delta}} \subset \mathcal{C}_{\delta}:=\left\{x \in \mathbf{R}^{n}: \operatorname{dist}\left(\frac{x}{|x|},\left\{p_{2}=0\right\}\right) \leq \delta\right\} .
$$

Denote by $\tilde{\mathcal{C}_{\delta}}:=\mathbf{R}^{n} \backslash \mathcal{C}_{\delta}$ the complementary cone, and let $\psi_{\delta}(x):=|x|^{\mu_{\delta}} \Psi_{\delta}(x /|x|)$, where $\Psi_{\delta} \geq 0$ is the first eigenfunction of the spherical Laplacian in $\tilde{\mathcal{C}}_{\delta} \cap \mathbf{S}^{n-1}$ with zero boundary conditions and $\mu_{\delta}$ is chosen so that the first eigenvalue is $\left.\mu_{\delta}\left(n-2+\mu_{\delta}\right)\right)$. In this way $\psi_{\delta}$ is a $\mu_{\delta}$-homogeneous harmonic function vanishing on the boundary of $\tilde{\mathcal{C}}_{\delta}$.

Since $\operatorname{dim}\left(\left\{p_{2}=0\right\}\right)=m \leq n-2$, the set $\left\{p_{2}=0\right\}$ has zero capacity and so $\psi_{\delta}$ converges to a positive constant as $\delta \downarrow 0$. Thus $\mu_{\delta} \downarrow 0$, and we can choose $\delta=\delta(\varepsilon)>0$ such that $\mu_{2 \delta}<\varepsilon$.

We now observe that, for $t \geq 0$, we have

$$
\{u(\cdot, t)>0\} \supset\{u(\cdot, 0)>0\} \supset \tilde{\mathcal{C}}_{\delta} \cap \mathrm{B}_{r_{\delta}} \supset \tilde{\mathcal{C}}_{2 \delta} \cap \mathrm{~B}_{r_{\delta}},
$$

and $v:=u(\cdot, t)-u(\cdot, 0)$ is nonnegative and harmonic in $\{u(\cdot, t)>0\}$. Note also that, by the maximum principle, every connected component of $\{u(\cdot, t)>0\}$ must have a part of its boundary on $\partial \mathrm{B}_{1}$, and thus (1.2) and the Harnack inequality (applied to a chain of balls) imply that

$$
v \geq c_{\delta} t \quad \text { in } \tilde{\mathcal{C}}_{2 \delta} \cap \partial \mathbf{B}_{r \delta}, \quad c_{\delta}>0 .
$$

Hence we can use the function

$$
v^{\prime}:=\frac{c_{\delta} \psi_{2 \delta}}{\left\|\psi_{2 \delta}\right\|_{L^{\infty}\left(\partial \mathrm{B}_{\left.r_{\delta}\right)}\right.}} t
$$

as lower barrier, and applying the maximum principle we obtain $v-v^{\prime} \geq 0$ inside the domain $\tilde{\mathcal{C}}_{2 \delta} \cap \mathrm{~B}_{r \delta}$. Since $\mathrm{D}_{r_{\delta}} \subset \tilde{\mathcal{C}}_{2 \delta} \cap \mathrm{~B}_{r \delta}$, this proves that

$$
\min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)]=\min _{\mathrm{D}_{r}} v \geq \min _{\mathrm{D}_{r}} v^{\prime}=c r^{\mu_{2 \delta}} t \geq c r^{\varepsilon} t \quad \forall r \in\left(0, r_{\delta}\right) .
$$

(b) After a rotation, we may assume that $\left\{p_{2}=0\right\}=\left\{x_{n}=0\right\}$. By Propositions 3.9 and 3.10, we have that $\{u(\cdot, 0)>0\} \supset\left\{\left|x_{n}\right| \leq \mathrm{C}\left|x^{\prime}\right|^{1+\alpha_{0}}\right\}$ in a neighborhood of the origin, where $x=\left(x^{\prime}, x_{n}\right)$ and $\alpha_{\circ}>0$. In particular, there exists a $\mathrm{C}^{1, \alpha_{o}}$ domain $\Omega$ contained inside $\{u(\cdot, 0)>0\}$ and satisfying $0 \in \partial \Omega$. By monotonicity of $u$ in $t$, the same domain $\Omega$ is contained in $\{u(\cdot, t)>0\}$ for $t>0$.

Hence, the function $v:=u(\cdot, t)-u(\cdot, 0)$ is positive and harmonic in $\Omega$, and by assumption (1.2) we have - as in the proof of (a) - that $v \geq c_{1} t>0$ in a small ball $\mathrm{B} \subset \subset \Omega$. Using Hopf's lemma in $\mathrm{C}^{1, \alpha}$ domains, we deduce that $\partial_{{x_{n}} v(0) \geq c_{2} t>0 \text {, and the result }}$ follows.

We can now prove the following key result:
Lemma 9.3. - Let $u: \mathrm{B}_{1} \times(-1,1) \rightarrow[0, \infty)$ satisfy (6.1) and (1.2), let $\alpha_{\circ}>0$ be given by Proposition 3.9, and let $\boldsymbol{\Sigma}^{*} \subset \boldsymbol{\Sigma}_{n-1}$ be given by Theorem 8.7.
(a) If $(0,0) \in \boldsymbol{\Sigma}_{m}^{a}$ and $m \leq n-2$, then for all $\varepsilon>0$ there exists $\varrho>0$ such that

$$
\left\{(x, t) \in \mathrm{B}_{\varrho} \times(0,1): t>|x|^{2-\varepsilon}\right\} \cap\{u=0\}=\emptyset
$$

(b) If $(0,0) \in \boldsymbol{\Sigma}_{m} \backslash \mathbf{\Sigma}_{m}^{a}, m \leq n-2$, then for all $\varepsilon>0$ there exists $\varrho>0$ such that

$$
\left\{(x, t) \in \mathrm{B}_{\varrho} \times(0,1): t>|x|^{3-\varepsilon}\right\} \cap\{u=0\}=\emptyset
$$

(c) If $(0,0) \in \mathbf{\Sigma}_{n-1}^{<3}$, then there exist $\mathrm{C}, \varrho>0$ such that

$$
\left\{(x, t) \in \mathrm{B}_{\varrho} \times(0,1): t>\mathrm{C}|x|^{1+\alpha_{0}}\right\} \cap\{u=0\}=\emptyset .
$$

(d) If $(0,0) \in \boldsymbol{\Sigma}_{n-1}^{>3}$, then there exist $\delta, \varrho>0$ such that

$$
\left\{(x, t) \in \mathrm{B}_{Q} \times(0,1): t>|x|^{2+\delta}\right\} \cap\{u=0\}=\emptyset .
$$

(e) If $(0,0) \in \mathbf{\Sigma}^{*}$, then for all $\varepsilon>0$ there exists $\varrho>0$ such that

$$
\left\{(x, t) \in \mathrm{B}_{\varrho} \times(0,1): t>|x|^{4-\varepsilon}\right\} \cap\{u=0\}=\emptyset
$$

Proof. - After a rotation, we may assume $\left\{p_{2}=0\right\} \subset\left\{x_{n}=0\right\}$. In all the following cases we will apply Lemma 9.1 and use that $\psi^{r} \geq \frac{1}{4}$ in $\mathrm{B}_{r / 2}$.
(a) By Lemma 9.2(a) we have, for any $\varepsilon>0$,

$$
\begin{equation*}
\min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \geq c_{\varepsilon} r^{\varepsilon / 2} t . \tag{9.2}
\end{equation*}
$$

Also, since $u$ is $\mathrm{C}^{1,1},|u(\cdot, 0)| \leq \mathrm{C}_{0} r^{2}$ in $\mathrm{B}_{r}$ for all $r \in(0,1 / 2)$. Thus, by Lemma 9.1 applied with $\mathfrak{p} \equiv 0$ and $\beta=2$,

$$
\begin{equation*}
u(\cdot, t) \geq c_{1} \min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \psi^{r}-\mathrm{C}_{0} r^{2} \quad \text { in } \mathrm{B}_{r}, \forall r \in(0,1 / 2) \tag{9.3}
\end{equation*}
$$

Since $\psi^{r} \geq \frac{1}{4}$ in $\mathrm{B}_{r / 2}$, thanks to (9.2) we deduce that

$$
u(\cdot, t)>0 \quad \text { in } \mathrm{B}_{r / 2} \text { for } t \geq(r / 2)^{2-\varepsilon},
$$

therefore

$$
\{u=0\} \cap\left\{t>|x|^{2-\varepsilon}\right\}=\emptyset .
$$

(b) Using again Lemma 9.2(a), it follows that (9.2) holds. Also, since $(0,0) \in \boldsymbol{\Sigma}_{m} \backslash$ $\boldsymbol{\Sigma}_{m}^{a}$, it follows from Proposition 3.9(a) that $\lambda^{2 n d} \geq 3$. Hence Lemmas 3.6 and 3.7 imply that $\left|u(\cdot, 0)-p_{2}\right| \leq \mathrm{C}_{0} r^{3}$ in $\mathrm{B}_{r}$, and therefore Lemma 9.1 applied with $\mathfrak{p} \equiv p_{2}$ and $\beta=3$ yields

$$
u(\cdot, t) \geq p_{2}+c_{1} \min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \psi^{r}-\mathrm{C}_{0} r^{3} \quad \text { in } \mathrm{B}_{r}, \forall r \in(0,1 / 2) .
$$

Since $p_{2} \geq 0$, one concludes as in the proof of (a).
(c) By Lemma 9.2(b) we have

$$
\begin{equation*}
\min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \geq c r t \tag{9.4}
\end{equation*}
$$

Since at the maximal stratum the frequency is at least $2+\alpha_{\circ}$ (see Proposition 3.9(b)), using Lemmas 3.6 and 3.7 we have $\left|u(\cdot, 0)-p_{2}\right| \leq \mathrm{C}_{0} r^{2+\alpha_{\circ}}$ in $\mathrm{B}_{r}$. Therefore, it follows from Lemma 9.1 applied with $\mathfrak{p} \equiv p_{2}$ and $\beta=2+\alpha_{o}$, that

$$
u(\cdot, t) \geq p_{2}+c_{1} \min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \psi^{r}-\mathrm{C}_{0} r^{2+\alpha_{o}} \quad \text { in } \mathrm{B}_{r}, \forall r \in(0,1 / 2) .
$$

Thus, since $\psi^{r} \geq \frac{1}{4}$ in $\mathrm{B}_{r / 2}$, thanks to (9.4) we obtain

$$
u(\cdot, t)>0 \quad \text { in } \mathrm{B}_{r / 2} \text { for } t \geq \mathrm{C}_{3} r^{1+\alpha_{o}} .
$$

(d) Again, (9.4) holds as a consequence of Lemma 9.2(b). Moreover, since $(0,0) \in$ $\boldsymbol{\Sigma}_{n-1}^{>3}$, thanks to Lemma 4.7 we deduce that $|u(\cdot, 0)-\mathscr{P}| \leq \mathrm{C}_{0} r^{3+2 \delta}$ in $\mathrm{B}_{r}$ for some $\delta>0$ (note that $\delta$ may depend on the point $(0,0)$ ). Therefore, Lemma 9.1 applied with $\mathfrak{p} \equiv \mathscr{P}$ and $\beta=3+2 \delta$ yields

$$
u(\cdot, t) \geq \mathscr{P}+c_{1} \min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \psi^{r}-\mathrm{C}_{0} r^{3+2 \delta} \quad \text { in } \mathrm{B}_{r}, \forall r \in(0,1 / 2)
$$

Recalling that $\mathscr{P} \geq-\overline{\mathrm{C}}|x|^{5}$, it follows from (9.4) that, for $t>(r / 2)^{2+\delta}$ and $r$ sufficiently small,

$$
\begin{aligned}
& u(\cdot, t) \geq \mathscr{P}+c_{3} r t-\mathrm{C}_{0} r^{3+\delta} \geq-\overline{\mathrm{C}} r^{5}+c_{3} r t-\mathrm{C}_{0} r^{3+2 \delta}>0 \\
& \quad \text { in } \mathrm{B}_{r / 2} \cap\{u(\cdot, 0)=0\} .
\end{aligned}
$$

Since $u(\cdot, t) \geq u(\cdot, 0)$, this proves the result.
(e) Again, (9.4) holds as a consequence of Lemma 9.2(b). Moreover, by Theorem 8.7, for every $\varepsilon>0$ we have $\left|u(\cdot, 0)-\mathscr{P}-p_{4}\right| \leq \mathrm{C}_{0} r^{5-\varepsilon / 2}$ in $\mathrm{B}_{r}$. Then, applying Lemma 9.1 with $\mathfrak{p} \equiv \mathscr{P}+p_{4}$ and $\beta=5-\varepsilon / 2$,

$$
\begin{aligned}
& u(\cdot, t) \geq \mathscr{P}+p_{4}+c_{1} \min _{\mathrm{D}_{r}}[u(\cdot, t)-u(\cdot, 0)] \psi^{r}-\mathrm{C}_{0} r^{5-\varepsilon / 2} \\
& \quad \text { in } \mathrm{B}_{r} \forall r \in(0,1 / 2) .
\end{aligned}
$$

Also,

$$
\mathscr{P}+p_{4} \geq-\overline{\mathrm{C}}|x|^{5} \quad \text { in }\{u(\cdot, 0)=0\} \subset\left\{x_{n} \leq \mathrm{C}\left|x^{\prime}\right|^{2}\right\} .
$$

Thus (9.4) yields, for $t>r^{4-\varepsilon}$ and $r$ small,

$$
\begin{aligned}
& u(\cdot, t) \geq \mathscr{P}+p_{4}+c_{3} r t-\mathrm{C}_{0} r^{5-\varepsilon / 2} \geq-\overline{\mathrm{C}} \mathrm{r}^{5}+c_{3} r t-\mathrm{C}_{0} r^{5-\varepsilon / 2}>0 \\
& \quad \text { in } \mathrm{B}_{r / 2} \cap\{u(\cdot, 0)=0\} .
\end{aligned}
$$

The set $\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}$ is treated separately in the following lemma. Since in this case the 3 rd order blow-up is not harmonic, the proof is more involved. In particular, instead of proving that there are no singular points in the "future" $t>0$, we show that they do not exists in the past.

Lemma 9.4. - Let $u: \mathrm{B}_{1} \times(-1,1) \rightarrow[0, \infty)$ satisfy (6.1) and (1.2), with $(0,0) \in$ $\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}$. Then

$$
\left\{(x, t) \in \mathbf{B}_{1} \times(-1,0): t<-\omega(|x|)|x|^{2}\right\} \cap \boldsymbol{\Sigma}_{n-1}^{\geq 3}=\emptyset,
$$

for some modulus a continuity $\omega:[0, \infty) \rightarrow[0, \infty)$.
Proof. - Let $w=u(\cdot 0)-p_{2}, w_{r}=w(r \cdot)$. Also, with no loss of generality we assume that $p_{2}=\frac{1}{2} x_{n}^{2}$. By Proposition 5.4 we have that
(9.5) $\left\|r^{-3} w_{r}-\tilde{q}\right\|_{L^{\infty}\left(\mathrm{B}_{4}\right)} \leq \delta(r) \downarrow 0 \quad$ for $\tilde{q}(x)=\left|x_{n}\right|\left(\frac{a}{3} x_{n}^{2}-x^{\prime} \cdot \mathrm{A} x^{\prime}\right)+x_{n}\left(\frac{b}{3} x_{n}^{2}-x^{\prime} \cdot \mathrm{B} x^{\prime}\right)$, where $x=\left(x^{\prime}, x_{n}\right)$, and $\mathrm{A} \in \mathbf{R}^{n-1} \times \mathbf{R}^{n-1}$ is symmetric, nonnegative definite, and has at least one positive eigenvalue.

Fix $\eta>0$, and assume by contradiction that there exists $r>0$ small and $t \leq-\eta r^{2}$ such that $u(\cdot, t)$ has a singular point in $\boldsymbol{\Sigma}_{n-1}^{\geq 3} \cap \mathbf{B}_{r}$. Under this assumption, we claim that

$$
\begin{equation*}
\left\{x_{n}+\frac{b}{3} x_{n}^{2}=x^{\prime} \cdot \mathrm{B} x^{\prime}\right\} \cap \mathrm{B}_{2 r} \subset\{u(\cdot, t / 2)=0\} \tag{9.6}
\end{equation*}
$$

where $b$ and B are given by (9.5).

Before proving the claim, we show that it leads to a contradiction. Indeed, thanks to Proposition 3.10, since we are assuming that $u(\cdot, t)$ has a singular point $x_{r} \in \Sigma_{n-1}^{\geq 3}$ with $\left|x_{r}\right| \leq r$, then for some $e_{r} \in \mathbf{S}^{n-1}$ we have

$$
\begin{equation*}
\{u(\cdot, t)=0\} \cap \mathrm{B}_{\rho} \subset\left\{x \in \mathrm{~B}_{\rho}:\left|e_{r} \cdot\left(x-x_{r}\right)\right| \leq \mathrm{C} \rho^{2}\right\}, \quad \text { for all } \rho \in[r, 1] . \tag{9.7}
\end{equation*}
$$

Note that the hypersurface $\left\{x_{n}+b x_{n}^{2}=x^{\prime} \cdot \mathbf{B} x^{\prime}\right\} \cap \mathbf{B}_{2 r}$ separates the ball $\mathbf{B}_{2 r}$ in two connected components $\mathbf{B}_{2 r}^{+}$and $\mathbf{B}_{2 r}^{-}$. Also, by monotonicity, (9.6) and (9.7) hold for $u\left(\cdot, t^{\prime}\right)$ for all $t^{\prime} \in[t, t / 2]$. Hence, if we define

$$
u^{+}\left(x, t^{\prime}\right):=\left\{\begin{array}{ll}
u\left(x, t^{\prime}\right) & \text { in } \mathbf{B}_{22}^{+}, \\
0 & \text { in } \mathbf{B}_{2 r}^{-},
\end{array} \quad u^{-}\left(x, t^{\prime}\right):= \begin{cases}0 & \text { in } \mathbf{B}_{2 r}^{+}, \\
u\left(x, t^{\prime}\right) & \text { in } \mathbf{B}_{2 r}^{-},\end{cases}\right.
$$

then both $u^{+}\left(\cdot, t^{\prime}\right)$ and $u^{-}\left(\cdot, t^{\prime}\right)$ are solutions to the obstacle problem with a thick contact set at 0 . Combining this information with (9.7), it follows by [Caf77] that the free boundaries of $u^{+}\left(\cdot, t^{\prime}\right)$ and $u^{-}\left(\cdot, t^{\prime}\right)$ are uniformly smooth hypersurface inside $\mathrm{B}_{3 r / 2}$, for all $t^{\prime} \in[t, t / 2]$. In addition, by strict monotonicity, these hypersurfaces are disjoint for any $t^{\prime}<t / 2$. Since the free boundary of $u\left(\cdot, t^{\prime}\right)$ is the union of these hypersurfaces, this proves in particular that the free boundary of $u(\cdot, t)$ has no singular points, a contradiction.

Thus, we are left with proving (9.6). First of all we note that, by Lemma 6.3, we have
(9.8) $\quad e_{r} \rightarrow \boldsymbol{e}_{n} \quad$ and $\quad\left(r^{-1} x_{r}\right) \cdot e_{r} \rightarrow 0 \quad$ as $r \downarrow 0$,
where $e_{r}$ is the unit vector appearing in (9.7). Furthermore, by the classical barrier argument used in proof of Hopf's Lemma (see for instance [Eva10, Chapter 6.4.2]), it follows from (9.7) that

$$
\begin{equation*}
u(\cdot, 0)-u(\cdot, t) \geq c_{1}|t|\left(\left|e_{r} \cdot\left(x-x_{r}\right)\right|-\mathrm{C}\left|x-x_{r}\right|^{2}\right)_{+} . \tag{9.9}
\end{equation*}
$$

Now, given $z^{\prime} \in \mathbf{B}_{2}^{\prime} \subset \mathbf{R}^{n-1}$ and $c \geq 0$, we define the function

$$
\phi_{z^{\prime}, c}(x):=\left(\frac{1}{2 r}-n\right)\left(x_{n}+\frac{b r}{3} x_{n}^{2}-r x^{\prime} \cdot \mathbf{B} x^{\prime}\right)^{2}+\left(x^{\prime}-z^{\prime}\right)^{2}+c .
$$

Note that $\phi_{z^{\prime}, c} \geq c \geq 0$ and

$$
\Delta \phi_{z^{\prime}, c}=\frac{1}{r}-2 n+\mathrm{O}(r)+2(n-1) \leq \frac{1}{r}, \quad \text { provided } 0<r \ll 1 .
$$

Also, since $\mathrm{A} \geq 0$ we have $\tilde{q}(x) \leq-x_{n}\left(x^{\prime} \cdot \mathrm{B} \chi^{\prime}\right)+\mathrm{C}\left|x_{n}\right|^{3}$, therefore (recall (9.5))

$$
r^{-3} u(r x, 0)-\frac{1}{2 r} x_{n}^{2}=r^{-3} w_{r}(x) \leq \tilde{q}(x)+\delta(r) \leq-x_{n}\left(x^{\prime} \cdot \mathrm{B} x^{\prime}\right)+\mathrm{C}\left|x_{n}\right|^{3}+\delta(r) .
$$

Thus, combining the bound above with (9.9), we get

$$
\begin{aligned}
& r^{-3} u(r x, t)-\frac{1}{2 r} x_{n}^{2} \\
& \quad \leq r^{-3}(u(r x, t)-u(r x, 0))-x_{n}\left(x^{\prime} \cdot \mathrm{B} x^{\prime}\right)+\mathrm{C}\left|x_{n}\right|^{3}+\delta(r) \\
& \quad \leq-r^{-3} c_{1}|t|\left(\left|e_{r} \cdot\left(r x-x_{r}\right)\right|-\mathrm{C}\left|r x-x_{r}\right|^{2}\right)_{+}-x_{n}\left(x^{\prime} \cdot \mathrm{B} x^{\prime}\right)+\mathrm{C}\left|x_{n}\right|^{3}+\delta(r) \\
& \quad \leq-c_{1} \eta\left(\left|e_{r} \cdot\left(x-\hat{x}_{r}\right)\right|-\mathrm{C} r\left|x-\hat{x}_{r}\right|^{2}\right)_{+}-x_{n}\left(x^{\prime} \cdot \mathrm{B} x^{\prime}\right)+\mathrm{C}\left|x_{n}\right|^{3}+\delta(r),
\end{aligned}
$$

where we used that $|t| \geq \eta r^{2}$ and we denoted $\hat{x}_{r}=: r^{-1} x_{r} \in \mathrm{~B}_{1}$.
Recalling (9.8), this implies that

$$
v(x):=r^{-3} u(r x, t) \leq \frac{1}{2 r} x_{n}^{2}-x_{n}\left(x^{\prime} \cdot \mathrm{B} x^{\prime}\right)-c_{1} \eta\left|x_{n}\right|+\mathrm{C}\left|x_{n}\right|^{3}+\theta(r),
$$

for some modulus of continuity $\theta(r)$. On the other hand, for any $c \geq 0$ we have

$$
\phi_{z^{\prime}, c}(x) \geq \frac{1}{2 r} x_{n}^{2}-n x_{n}^{2}-\mathrm{C}\left|x_{n}\right|^{3}-x_{n}\left(x^{\prime} \cdot \mathrm{B} x^{\prime}\right)+\left(x^{\prime}-z^{\prime}\right)^{2}+\mathrm{O}(r) .
$$

Let now $z:=\left(z^{\prime}, z_{n}\right)$ satisfy $z_{n}+r \frac{b}{3} z_{n}^{2}=r z^{\prime} \cdot \mathrm{B} z^{\prime}$, and consider a point $x \in \partial \mathrm{~B}_{s}(z)$, where $0<r \ll s \ll 1$. Then, since $\left|z_{n}\right|=\mathrm{O}(r)$, we have $\left(x^{\prime}-z^{\prime}\right)^{2}=s^{2}-x_{n}^{2}+\mathrm{O}(r)$, and therefore

$$
\phi_{z^{\prime}, c}(x) \geq \frac{1}{2 r} x_{n}^{2}-(n+1) x_{n}^{2}-\mathrm{C}\left|x_{n}\right|^{3}-x_{n}\left(x^{\prime} \cdot \mathrm{B} x^{\prime}\right)+s^{2}+\mathrm{O}(r) \quad \text { on } \partial \mathrm{B}_{s}(z) .
$$

Since, for $r \ll s \ll 1$,

$$
c_{1} \eta\left|x_{n}\right|+s^{2} \geq \mathrm{C} x_{n}^{2}+\mathrm{C}\left|x_{n}\right|^{3}+\theta(r) \quad \text { for }\left|x_{n}\right| \leq s,
$$

we deduce that

$$
v(x)=r^{-3} u(r x, t)<\phi_{z^{\prime}, c}(x) \quad \text { on } \partial \mathrm{B}_{s}(z)
$$

for all $c>0$.
Now, assume there exists $c_{*}>0$ be such that $\phi_{z^{\prime}, c_{*}}$ touches $v$ from above at some point $x_{0} \in \overline{\mathbf{B}_{s}(z)}$. Since $v<\phi_{z^{\prime}, c_{*}}$ on $\partial \mathbf{B}_{s}(z)$, the contact point is inside $\mathrm{B}_{s}(z)$. Also, since $\Delta v=r^{-1}$ in $\{u(r \cdot, t)>0\}$ while $\Delta \phi_{z^{\prime}, c_{*}} \leq r^{-1}$, it follows by the maximum principle that $x_{\circ} \notin\{u(r \cdot, t)>0\}$. Thus,

$$
0=r^{-3} u\left(r x_{\circ}, t\right)=v\left(x_{\circ}\right)=\phi_{z^{\prime}, c_{*}}\left(x_{\circ}\right) \geq c_{*}>0
$$

a contradiction. This proves that $v \leq \phi_{z^{\prime}, c}$ for all $c>0$, and letting $c \rightarrow 0$ we obtain

$$
0 \leq r^{-3} u(r z, t)=v(z) \leq \phi_{z^{\prime}, 0}(z)=0 .
$$

Since $z^{\prime} \in \mathrm{B}_{2}^{\prime}$ is arbitrary, this proves (9.6), and the lemma follows.

We finally prove:
Theorem 9.5. - Let $u: \mathrm{B}_{1} \times(-1,1) \rightarrow[0, \infty)$ satisfy (6.1) and (1.2). Then:
(a) In dimension $n=2$ we have $\operatorname{dim}_{\mathcal{H}}\left(\pi_{2}(\boldsymbol{\Sigma})\right) \leq 1 / 4$.
(b) In dimension $n=3$ we have $\operatorname{dim}_{\mathcal{H}}\left(\pi_{2}(\boldsymbol{\Sigma})\right) \leq 1 / 2$.
(c) In dimensions $n \geq 4$, for $\mathcal{H}^{1}$-a.e. $t \in(-1,1)$ we have

$$
\mathcal{H}^{n-4}\left(\boldsymbol{\Sigma} \cap \pi_{2}^{-1}(\{t\})\right)=0
$$

In particular, for $n \leq 4$, the singular set is empty for a.e. $t$.

Proof. - First of all, as in the proof of Theorem 8.7, we have the following:

- $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{m}^{a}\right)\right) \leq n-3$ for $0 \leq m \leq n-2$;
- $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{m} \backslash \boldsymbol{\Sigma}_{m}^{a}\right)\right) \leq n-2$ for $0 \leq m \leq n-2$;
- $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 3}\right)\right) \leq n-3$;
- $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}\right)$ is contained in a countable union of $(n-2)$-dimensional Lipschitz graphs;
- $\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{3 r d} \backslash \boldsymbol{\Sigma}_{n-1}^{>3}\right)=\emptyset$;
- $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}^{*}\right)\right) \leq n-2 ;$
- $\operatorname{dim}_{\mathcal{H}}\left(\pi_{1}\left(\boldsymbol{\Sigma}^{*}\right)\right) \leq n-1$.

Furthermore, thanks to Lemmas 9.3 and 9.4, we have:

- In $\boldsymbol{\Sigma}_{m}^{a}$ for $0 \leq m \leq n-2$, we can use Corollary 7.8 with $\beta=n-3$ and $k=2$;
- In $\boldsymbol{\Sigma}_{m} \backslash \boldsymbol{\Sigma}_{m}^{a}$ we can use Corollary 7.8 with $\beta=n-2$ and $k=3$;
- In $\boldsymbol{\Sigma}_{n-1} \backslash \boldsymbol{\Sigma}_{n-1}^{\geq 3}$ we can use Corollary 7.8 with $\beta=n-3$ and $k=1+\alpha_{o}$;
- In $\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 n d}$ (up to taking a countable union, and up to reversing time) we can use Proposition 7.7 with $\beta=n-2$ and $k=2$;
- In $\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}^{*}$ we can use Proposition 7.7 with $\beta=n-2$ and $k=2$;
- In $\boldsymbol{\Sigma}^{*}$ we can use Corollary 7.8 with $\beta=n-1$ and $k=4$.

Hence, combining these information, we deduce that:

- $\operatorname{dim}_{\mathcal{H}}\left(\boldsymbol{\Sigma}_{m}^{a} \cap \pi_{2}^{-1}(\{t\})\right) \leq n-5$ for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$;
- $\operatorname{dim}_{\mathcal{H}}\left(\left(\boldsymbol{\Sigma}_{m} \backslash \boldsymbol{\Sigma}_{m}^{a}\right) \cap \pi_{2}^{-1}(\{t\})\right) \leq n-5$ for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$;
- $\operatorname{dim}_{\mathcal{H}}\left(\boldsymbol{\Sigma}_{n-1}^{<3} \cap \pi_{2}^{-1}(\{t\})\right) \leq n-4-\alpha$ 。for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$;
- $\mathcal{H}^{n-4}\left(\left(\boldsymbol{\Sigma}_{n-1}^{\geq 3} \backslash \boldsymbol{\Sigma}_{n-1}^{3 r d}\right) \cap \pi_{2}^{-1}(\{t\})\right)=0$ for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$;
- $\mathcal{H}^{n-4}\left(\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}^{*}\right) \cap \pi_{2}^{-1}(\{t\})\right)=0$ for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$;
- $\operatorname{dim}_{\mathcal{H}}\left(\boldsymbol{\Sigma}^{*} \cap \pi_{2}^{-1}(\{t\})\right) \leq n-5$ for $\mathcal{H}^{1}$-a.e. $t \in \mathbf{R}$.

Thus, part (c) is proved. Parts (a) and (b) follow exactly in the same way, but using instead Proposition 7.7(a) and Corollary 7.8(a).

Remark 9.6. - Thanks to Remark 8.5, one could actually slightly improve the estimate for the set $\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}^{*}$ and show that $\operatorname{dim}_{\mathcal{H}}\left(\left(\boldsymbol{\Sigma}_{n-1}^{>3} \backslash \boldsymbol{\Sigma}^{*}\right) \cap \pi_{2}^{-1}(\{t\})\right) \leq n-4-\frac{1}{2}$. However, all the other estimates are sharp (at least with respect to the techniques introduced in this paper), and in particular we believe that it is very unlikely that one could prove a stronger result with these techniques.

As a consequence of the previous estimates, we finally obtain our main results:
Proof of Theorems 1.1 and 1.2. - The results follow immediately from Theorem 9.5.

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## Appendix A: Some results on the Signorini problem

For the convenience of the reader, in this appendix we gather some classical results on the Signorini problem (5.1) that we use several times throughout the paper

Lemma A.1. - The only $1 D$ solutions to (5.1) that vanish at the origin are given by $q\left(x_{n}\right)=$ $-c\left|x_{n}\right|+b x_{n}$, for some $c \geq 0$ and $b \in \mathbf{R}$.

Proof. - Since $q=q\left(x_{n}\right)$, it follows from (5.1) that $q$ must be affine in $\mathbf{R}^{n} \backslash\{0\}$, hence $q\left(x_{n}\right)=a-c\left|x_{n}\right|+b x_{n}$ for some $a, b, c \in \mathbf{R}$. The condition $\Delta q \leq 0$ implies that $c \geq 0$. Also, since $q(0)=0$ we deduce that $a=0$, as desired.

Lemma A.2. - Let $\lambda>0$, and let $i$ denote the imaginary unit. The only $2 D \lambda$-homogeneous solutions of (5.1) (i.e., $q=q\left(x_{n}, x_{n-1}\right)$ and $q(r x)=r^{\lambda} q(x)$ for every $r>0$ ) are given by

$$
\left\{\begin{array}{l}
q\left(x_{n}, x_{n-1}\right)=c i^{1-\lambda} \operatorname{Re}\left(\left|x_{n}\right|+i x_{n-1}\right)^{\lambda}+b \operatorname{Re}\left(x_{n}+i x_{n-1}\right)^{\lambda}, \\
\quad \text { if } \lambda \in\{1,3,5, \ldots\} \\
q\left(x_{n}, x_{n-1}\right)=c i^{\lambda} \operatorname{Re}\left(x_{n}+i x_{n-1}\right)^{\lambda}+b \operatorname{Im}\left(x_{n}+i x_{n-1}\right)^{\lambda}, \\
\quad \text { if } \lambda \in\{2,4,6, \ldots\}, \\
q\left(x_{n}, x_{n-1}\right)=c \operatorname{Re}\left(x_{n}+i x_{n-1}\right)^{\lambda}, \\
\quad \text { if } \lambda \in\left\{\frac{3}{2}, \frac{7}{2}, \frac{11}{2}, \ldots\right\},
\end{array}\right.
$$

where $c \geq 0$ and $b \in \mathbf{R}$. In particular the set of possible homogeneities is $\{1,2,3,4,5, \ldots\} \cup$ $\left\{\frac{3}{2}, \frac{7}{2}, \frac{11}{2}, \ldots\right\}$.

Proof. - A proof of this result can be found, for instance, in [FS18, Proposition A.1]; see also [GP09, Remark 1.2.7].

The following result is proved in [ACS08, Lemma 1].
Lemma A.3. - Let $q$ be a solution of (5.1) and assume that $q(0)=0$. Let $\phi(\cdot, q)$ be as in (2.4). Then $r \mapsto \phi(r, q)$ is monotone nondecreasing. Moreover, if $\mathrm{I} \ni r \mapsto \phi(r, q) \equiv \lambda>0$ for some open interval $\mathrm{I} \subset \mathbf{R}^{+}$, then $q$ is $\lambda$-homogeneous.

We conclude this section with a uniqueness result.
Lemma A.4. - Let $q_{i}, i=1,2$ be two solutions of (5.1) satisfying $q_{1} \geq q_{2}$ in $\mathbf{B}_{1}$ and $q_{i}(0)=0$. Assume that $\phi\left(0^{+}, q_{2}\right)>1$, or that $q_{2} \equiv 0$. Then $q_{1} \equiv q_{2}$.

Proof. - We use coordinates $\left(x^{\prime}, x_{n}\right) \in \mathbf{R}^{n-1} \times \mathbf{R}$. Assume by contradiction that $q_{1} \not \equiv q_{2}$. Then, applying Hopf's Lemma at the origin, we deduce that $\partial_{x_{n}}\left(q_{1}-q_{2}\right) \times$ $\left(0,0^{+}\right)>0$. Also, our assumption on $q_{2}$ implies that $\nabla q_{2}(0,0)=0$, thus $\partial_{x_{n}} q_{1}\left(0,0^{+}\right)>0$. On the other hand, the distributional Laplacian of $q_{1}$ on $\left\{x_{n}=0\right\}$ is given by $2 \partial_{x_{n}} q_{1}\left(x^{\prime}, 0^{+}\right)$. Since $\Delta q_{1} \leq 0$, this gives the desired contradiction.

## Appendix B: Odd frequency points in the Signorini problem

The aim of this section is to show how the arguments developed in this paper (see in particular Section 5) can be applied in the context of the Signorini problem to prove both uniqueness and nondegeneracy of blow-ups at all points of odd frequency for solutions of the Signorini problem
(B.1) $\left\{\begin{array}{lll}\Delta u \leq 0 & \text { and } \quad u \Delta u=0 & \text { in } \mathrm{B}_{1} \\ \Delta u=0 & & \text { in } \mathrm{B}_{1} \backslash\left\{x_{n}=0\right\} \\ u \geq 0 & \text { on } \mathrm{B}_{1} \cap\left\{x_{n}=0\right\},\end{array}\right.$
see Theorem B. 7 below. Since this was an open problem in this topic which we expect to be of interest to a wide audience, we prefer to give a complete and self-contained proof (rather than referring to parts of this paper) so that this appendix can become of reference for future results on the Signorini problem.

Note that this appendix extends the results of [GP09] (which were dealing only with even frequencies) to the sets $\Gamma_{2_{m+1}}(u), m \in \mathbf{N}$ (see [GP09] for an explanation of this notation).

Since the odd part of a solution of (B.1) is harmonic and vanishes on $\left\{x_{n}=0\right\}$, to understand the structure of the solution and the free boundary it suffices to study even solutions, that is, $u\left(x^{\prime}, x_{n}\right)=u\left(x^{\prime},-x_{n}\right)$. For this reason, also when studying global homogeneous solutions, we can restrict ourselves to even functions.

We begin by recalling Lemma 5.1: If $q$ is a $\lambda$-homogeneous even solution of (B.1) and $\lambda=2 m+1$ is an odd integer, then $q \equiv 0$ on $\left\{x_{n}=0\right\}$ (this result was not known before). As a consequence of this fact and the Liouville Theorem for harmonic functions vanishing on a hyperplane, ${ }^{9} q$ must be a harmonic polynomial on each side sides of $\left\{x_{n}=0\right\}$.

Then, since $\left.q\right|_{\left\{x_{n}=0\right\}}=0, q$ is even, and $q$ is harmonic in $\mathbf{R}^{n} \backslash\left\{x_{n}=0\right\}$, we deduce that

$$
q\left(x^{\prime}, x_{n}\right)=-\left|x_{n}\right|\left(q_{0}\left(x^{\prime}\right)+x_{n}^{2} q_{1}\left(x^{\prime}, x_{n}\right)\right), \quad\left(x^{\prime}, x_{n}\right) \in \mathbf{R}^{n-1} \times \mathbf{R},
$$

where $q_{0}$ and $q_{1}$ are polynomials. Furthermore, since $\Delta q \leq 0$, the polynomial $q_{0}\left(x^{\prime}\right)$ is nonnegative.

In the sequel it will be useful to define the "trace operator" T as

## (B.2) $\quad q \mapsto \mathrm{~T}[q]:=q_{0}$.

Since $q_{0} \equiv 0$ implies that $q \equiv 0$ (as a consequence the harmonicity of $q$ outside of $\left\{x_{n}=0\right\}$ ), one easily deduces that T is injective.

We will need a monotonicity formula that is the analogue of Lemma 5.3.
Lemma B.1. - Let $u: \mathrm{B}_{1} \rightarrow \mathbf{R}$ be an even solution of (B.1) with $\phi\left(0^{+}, u\right)=\lambda$, where $\lambda$ is an odd integer, and define $u_{r}(x):=u(r x)$. Also, let $q$ be any $\lambda$-homogeneous even solution of (B.1). Then, for any $\varrho \in(0,1)$,

$$
\frac{d}{d r} \int_{\partial \mathbf{B}_{e}} u_{r} q=\frac{\lambda}{r} \int_{\partial \mathbf{B}_{e}} u_{r} q-\frac{\varrho}{r} \int_{\mathrm{B}_{e}} u_{r} \Delta q .
$$

In particular

$$
\frac{d}{d r}\left(\frac{1}{r^{\lambda}} \int_{\partial \mathrm{B}_{1}} u_{r} q\right) \geq 0 .
$$

Proof. - We have

$$
\begin{aligned}
\frac{d}{d r} \int_{\partial \mathrm{B}_{e}} u_{r} q & =\int_{\partial \mathrm{B}_{e}} \frac{x}{r} \cdot \nabla u_{r} q=\frac{\varrho}{r} \int_{\partial \mathrm{B}_{e}} \partial_{\nu} u_{r} q=\frac{\varrho}{r} \int_{\mathrm{B}_{e}} \operatorname{div}\left(\nabla u_{r} q\right) \\
& =\frac{\varrho}{r}\left(\int_{\mathrm{B}_{e}} \nabla w_{r} \nabla q+\int_{\mathrm{B}_{e}} \Delta w_{r} q\right) \\
& =\frac{\varrho}{r}\left(\int_{\partial \mathrm{B}_{e}} u_{r} \partial_{\nu} q-\int_{\partial \mathrm{B}_{e}} u_{r} \Delta q+\int_{\mathrm{B}_{e}} \Delta u_{r} q\right) .
\end{aligned}
$$

[^9]Since $q$ is $\lambda$-homogeneous, we find that $\varrho \int_{\partial \mathrm{B}_{e}} u_{r} \partial_{\nu} q=\lambda \int_{\partial \mathrm{B}_{e}} u_{r} q$. Also, since $q$ vanishes on $\left\{x_{n}=0\right\}$ (by Lemma 5.1) and $\Delta u$ is a measure supported on $\left\{x_{n}=0\right\}$, we have $\int_{\mathrm{B}_{e}} \Delta u_{r} q=0$. This proves the first statement.

Finally, taking $\varrho=1$ and using that $-u_{r} \Delta q \geq 0$ in $\mathbf{R}^{n}$ (since $\Delta q \leq 0$ is supported on $\left\{x_{n}=0\right\}$, and $u_{r} \geq 0$ there) we obtain

$$
\frac{d}{d r}\left(\frac{1}{r^{\lambda}} \int_{\partial \mathrm{B}_{1}} u_{r} q\right)=-\frac{1}{r^{\lambda+1}} \int_{\partial \mathrm{B}_{1}} u_{r} \Delta q \geq 0 .
$$

As a consequence of the previous result, we deduce the following:
Proposition B.2. - Let $u: \mathrm{B}_{1} \rightarrow \mathbf{R}$ be an even solution of $(\mathrm{B} .1)$ with $\phi\left(0^{+}, u\right)=\lambda$, where $\lambda=2 m+1$ is an odd integer. Then the limit

$$
\tilde{q}:=\lim _{r \downarrow 0} \frac{u(r \cdot)}{r^{\lambda}}
$$

exists and it is a $\lambda$-homogeneous even solution of (B.1).

$$
\begin{aligned}
& \text { Proof. - Let } \\
& \qquad q^{(i)}=\lim _{\substack{(i) \\
r_{k}^{(i)} \downarrow 0}} \frac{1}{\left(r_{k}^{(i)}\right)^{\lambda}} u_{r_{k}^{(i)}}, \quad i=1,2,
\end{aligned}
$$

be two accumulation points along different sequences $r_{k}^{(i)}$. Then, given a $\lambda$-homogeneous solution of Signorini $q$, we can apply Lemma B. 1 to deduce that $r \mapsto \frac{1}{r^{\lambda}} \int_{\partial \mathrm{B}_{1}} u_{r} q$ has a unique limit as $r \rightarrow 0$. In particular this implies that

$$
\int_{\partial \mathbf{B}_{1}} q^{(1)} q=\int_{\partial \mathbf{B}_{1}} q^{(2)} q,
$$

therefore, choosing $q=q^{(1)}-q^{(2)}$, we deduce that $q^{(1)} \equiv q^{(2)}$.
The next step consists in showing the following nondegeneracy property: if $\phi\left(0^{+}, u\right)=\lambda$, then the limit $\tilde{q}$ obtained in Proposition B. 2 cannot be identically zero. This is the most delicate part of this appendix, and the proof of this fact requires a new compactness lemma and an interesting ODE type formula obtained below.

Lemma B.3.-Let $u: \mathrm{B}_{1} \rightarrow \mathbf{R}$ be an even solution of (B.1) satisfying $\phi\left(0^{+}, u\right)=\lambda$ with $\lambda$ odd, set $u_{r}(x):=u(r x)$, and let $\tilde{u}_{r}:=u_{r} /\left\|u_{r}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}$. Given $\eta>0$ there exists $\delta=\delta(n, \lambda, \eta)$ such that, if for some $r \in(0,1 / 2)$ and for some $\lambda$-homogeneous even solution $q$ of (B.1) we have

$$
\left\|\tilde{u}_{r}-q\right\|_{L^{\infty}\left(\mathbf{B}_{2}\right)} \leq \delta,
$$

then

$$
\tilde{u}_{r}=0 \quad \text { on }\left\{x_{n}=0\right\} \cap\left(\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}\right) \cap\{\mathrm{T}[q] \geq \eta\},
$$

where $\mathrm{T}[q]$ is defined as in (B.2).
Proof. - Fix $z=\left(z^{\prime}, 0\right) \in\left(\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}\right)$ such that $\mathrm{T}[q]\left(z^{\prime}\right) \geq \eta$, and given $c>0$ we define

$$
\phi_{c}(x):=-(n-1)\left|x_{n}\right|^{2}+\left|x^{\prime}\right|^{2}+c .
$$

Let $\varrho>0$ be sufficiently small (depending only on $n$ and $\eta$ ) and take $\delta=\varrho^{3}$. Then, for $|x|=\varrho$ we have
(B. 3 )

$$
\begin{aligned}
& u_{r}(z+x) \leq q(z+x)+\delta \\
& \quad=-\left|x_{n}\right| q_{0}\left(z^{\prime}\right)+\mathrm{O}\left(\varrho^{3}\right)+\delta \leq-\eta\left|x_{n}\right|+\mathrm{O}\left(\varrho^{3}\right)+\delta \\
& \quad \leq-n\left|x_{n}\right|^{2}+|x|^{2} \leq \phi_{c}(x) \quad \forall c \geq 0
\end{aligned}
$$

Since $\phi_{c}>q(z+\cdot)$ inside $\mathbf{B}_{\varrho}$ for $c$ large, we can decrease $c$ until a contact point occur inside $\overline{\bar{B}_{\varrho}}$. Since $\phi_{0}(0)=0 \leq q(z)$ (since $z \in\left\{x_{n}=0\right\}$ ), we see that such a contact point must occur for some value $c_{*} \geq 0$.

If $c_{*}=0$ then we have $u_{r}(z) \leq \phi_{0}(0)=0$, as wanted. Hence, it suffices to show that $c_{*}>0$ is impossible.

Assume by contradiction that there exists $c_{*}>0$ such that $\phi_{c_{*}} \geq q(z+\cdot)$ in $\overline{\mathbf{B}_{Q}}$, and $\phi_{c_{*}}\left(x_{\circ}\right)=q\left(z+x_{\circ}\right)$ for some $x_{\circ} \in \overline{\mathbf{B}_{\varrho}}$. By (B.3) we see that $\phi_{c_{*}}$ and $u_{r}(z+\cdot)$ must touch at an interior point, that is $x_{\circ} \notin \partial \mathbf{B}_{\varrho}$. Also, since $\phi_{c_{*}}$ is harmonic, it cannot touch $u_{r}(z+\cdot)$ at some point where it is harmonic too. Thus, $x_{\circ}$ must belong to $\left\{x_{n}=0\right\} \cap\left\{u_{r}(z+\cdot)=0\right\}$. This gives $0=u_{r}\left(z+x_{\circ}\right)=\phi_{c_{*}}\left(x_{\mathrm{o}}\right)=\left|x_{\mathrm{o}}\right|^{2}+c_{*}>0$, a contradiction.

Another fundamental tool is the following ODE-type formula.
Lemma B.4. - Let $u: \mathrm{B}_{1} \rightarrow \mathbf{R}$ be an even solution of (B.1) satisfying $\phi\left(0^{+}, u\right)=\lambda$, with $\lambda$ odd. Set $u_{r}(x):=u(r x), h(r):=\left\|u_{r}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}$, and $\tilde{u}_{r}:=u_{r} / h(r)$. Let $q$ be an even $\lambda$-homogeneous solution of (B.1), and define
(B.4)

$$
\psi(r ; q):=\int_{\partial \mathrm{B}_{1}} \tilde{u}_{r} q-2 \int_{\partial \mathbf{B}_{1 / 2}} \tilde{u}_{r} q .
$$

Then

$$
\begin{aligned}
& \frac{d}{d r} \psi(r ; q)=-\theta(r) \psi(r ; q)-\frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta q \\
& \text { where } \theta(r):=\left(\frac{h^{\prime}(r)}{h(r)}+\frac{\lambda}{r}\right)=\left(\log \left(h(r) / r^{\lambda}\right)\right)^{\prime}
\end{aligned}
$$

Proof. - As in the proof of Lemma B.1, we obtain

$$
\frac{d}{d r} \int_{\partial \mathrm{B}_{e}} u_{r} q=\frac{\lambda}{r} \int_{\partial \mathrm{B}_{e}} u_{r} q-\frac{\varrho}{r} \int_{\mathrm{B}_{e}} u_{r} \Delta q .
$$

Thus, since $\tilde{u}_{r}=u_{r} / h(r)$, we deduce that

$$
\frac{d}{d r} \int_{\partial \mathrm{B}_{e}} \tilde{u}_{r} q=\left(-\frac{h^{\prime}(r)}{h(r)}+\frac{\lambda}{r}\right) \int_{\partial \mathrm{B}_{e}} \tilde{u}_{r} q-\frac{\varrho}{r} \int_{\mathbf{B}_{e}} \tilde{u}_{r} \Delta q,
$$

and the lemma follows combining the identities for $\varrho=1$ and $\varrho=1 / 2$.
In the sequel, for $\lambda=2 m+1, m \in \mathbf{N}$, we denote

$$
\mathcal{Q}_{\lambda}:=\{\text { even } \lambda \text {-homogeneous solutions of (B.1) }\} .
$$

Also, given $f \in \mathrm{~L}_{\text {loc }}^{1}\left(\mathbf{R}^{n}\right)$, we define the radial symmetrization in the first $(n-1)$ variables as
(B.5)

$$
\widehat{f}\left(x^{\prime}, x_{n}\right):=\int_{\operatorname{SO}(n-1)} f\left(\mathrm{M} x^{\prime}, x_{n}\right) d \mathrm{M}, \quad x=\left(x^{\prime}, x_{n}\right) \in \mathbf{R}^{n-1} \times \mathbf{R},
$$

where the previous average is with respect to the Haar measure of $\mathrm{SO}(n-1)$.
Lemma B.5. - Given $\lambda \geq 3$ odd, there exists a unique $\mathrm{Q} \in \mathcal{Q}_{\lambda}$ satisfying

$$
\begin{equation*}
\mathrm{Q}=\widehat{\mathrm{Q}} \text { and }\|\mathrm{Q}\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1 \tag{B.6}
\end{equation*}
$$

Moreover, for any other $q \in \mathcal{Q}_{\lambda}$ we have

$$
\int_{\partial \mathrm{B}_{1}} q \mathrm{Q} \geq c_{n, \lambda}\|q\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}>0
$$

where $c_{n, \lambda}$ is some positive constant depending only on $n$ and $\lambda$.
Proof. - We begin by noticing that $\mathrm{Q}=\widehat{\mathrm{Q}}$ belongs to $\mathcal{Q}_{\lambda}$ if and only if
(B.7) $\quad \mathrm{Q}(x)=\sum_{k=0}^{\frac{\lambda-1}{2}} a_{k}\left|x^{\prime}\right|^{\lambda-1-2 k}\left|x_{n}\right|^{1+2 k}, \quad a_{0} \leq 0, \quad \Delta\left(\sum_{k=0}^{\frac{\lambda-1}{2}} a_{k}\left|x^{\prime}\right|^{\lambda-1-2 k} x_{n}^{1+2 k}\right)=0$.

Setting $r^{\prime}:=\left|x^{\prime}\right|$ and noticing that $\Delta \mathrm{Q}=\partial_{r^{\prime} r^{\prime}} \mathrm{Q}+\frac{n-2}{r^{\prime}} \partial_{\mu^{\prime}} \mathrm{Q}+\partial_{x_{n} x_{n}} \mathrm{Q}$, we can rewrite (B.7) as

$$
\begin{aligned}
& \sum_{j=0}^{\frac{\lambda-3}{2}}\left(a_{j}(\lambda-1-2 j)(\lambda-2-2 j+n-2)\right. \\
& \left.\quad+a_{j+1}(3+2 j)(2+2 j)\right)\left(r^{\prime}\right)^{\lambda-3-2 j} x_{n}^{1+2 j}=0 .
\end{aligned}
$$

Therefore, (B.7) is satisfied if and only if

$$
\begin{aligned}
& a_{j}(\lambda-1-2 j)(\lambda-2-2 j+n-2)+a_{j+1}(3+2 j)(2+2 j)=0 \\
& \quad \text { for all } j=0,1, \ldots \frac{\lambda-3}{2} .
\end{aligned}
$$

This means that all the coefficients are uniquely determined (by induction over $j$ ) once $a_{0} \leq 0$ is fixed, and $a_{0}<0$ is uniquely determined imposing that $\|\mathrm{Q}\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1$. This concludes the first part of the proof.

To prove the second part, note that if $q \in \mathcal{Q}_{\lambda}$ then $\widehat{q} \in \mathcal{Q}_{\lambda}$ (see (B.5)). We now recall that, to define the trace operator $T$, we used the expansion

$$
q\left(x^{\prime}, x_{n}\right)=-\left|x_{n}\right|\left(q_{0}\left(x^{\prime}\right)+x_{n}^{2} q_{1}\left(x^{\prime}, x_{n}\right)\right), \quad \text { so that } \mathrm{T}[q]=q_{0} .
$$

Since $\mathrm{T}[q] \equiv 0$ implies $q \equiv 0$, it follows by compactness that

$$
\|\mathrm{T}[q]\|_{\mathrm{L}^{2}\left(\partial \mathbf{B}_{1}\right)} \geq \tilde{c}_{n, \lambda}\|q\|_{\mathrm{L}^{2}\left(\partial \mathbf{B}_{1}\right)} \quad \forall q \in \mathcal{Q}_{\lambda}, \text { for some } c_{n, \lambda}>0 .
$$

Also

$$
\widehat{q}\left(x^{\prime}, x_{n}\right)=-\left|x_{n}\right|\left(\widehat{q_{0}}\left(x^{\prime}\right)+x_{n}^{2} \widehat{q_{1}}\left(x^{\prime}, x_{n}\right)\right), \quad \text { that is } \mathrm{T}[\widehat{q}]=\widehat{\mathrm{T}[q]} .
$$

Thus, given $q \in \mathcal{Q}_{\lambda}$, since $\widehat{q} \in \mathcal{Q}_{\lambda}$ depends only on the variables $r^{\prime}=\left|x^{\prime}\right|$ and $x_{n}$, it follows from the first part of the proposition that $\widehat{q}$ must by a positive multiple of Q that is, $\widehat{q}=t \mathrm{Q}$, where $t \geq c_{n, \lambda}\|q\|_{L^{2}\left(\partial \mathrm{~B}_{1}\right)}>0$. Hence, since $\widehat{\mathrm{Q}}=\mathrm{Q}$ and using the invariance of the Haar measure $d \mathrm{M}$ on $\mathrm{SO}(n-1)$ under the transformation $\mathrm{M} \mapsto \mathrm{M}^{-1}$, we get

$$
\begin{aligned}
\int_{\partial \mathbf{B}_{1}} q \mathbf{Q} & =\int_{\partial \mathbf{B}_{1}} q \widehat{\mathbf{Q}}=\int_{\partial \mathbf{B}_{1}} d x \int_{\operatorname{SO}(n-1)} q\left(x^{\prime}, x_{n}\right) \mathrm{Q}\left(\mathrm{M} x^{\prime}, x_{n}\right) d \mathbf{M} \\
& =\int_{\partial \mathbf{B}_{1}} d x{\underset{\operatorname{SO}(n-1)}{ } q\left(\mathrm{M}^{-1} x^{\prime}, x_{n}\right) \mathrm{Q}\left(x^{\prime}, x_{n}\right) d \mathbf{M}}=\int_{\partial \mathbf{B}_{1}} d x{\underset{\operatorname{SO}(n-1)}{ } q\left(\mathrm{M} x^{\prime}, x_{n}\right) \mathrm{Q}\left(x^{\prime}, x_{n}\right) d \mathbf{M}}=\int_{\partial \mathbf{B}_{1}} \widehat{q} \mathbf{Q}=t \int_{\partial \mathbf{B}_{1}} \mathrm{Q}^{2} \geq c_{n, \lambda}\|q\|_{\mathrm{L}^{2}\left(\partial \mathbf{B}_{1}\right)}
\end{aligned}
$$

In the following proposition we will use the notation $\mathrm{X} \asymp \mathrm{Y}$ for $\mathrm{X} \leq \mathrm{C}(n, \lambda) \mathrm{Y}$ and $\mathrm{Y} \leq \mathrm{C}(n, \lambda) \mathrm{X}$.

Proposition B.6. - Let $u: \mathrm{B}_{1} \rightarrow \mathbf{R}$ be an even solution of $(\mathrm{B} .1)$ with $\phi\left(0^{+}, u\right)=\lambda$, where $\lambda$ is an odd integer. Suppose that $\|u\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1$, and set $u_{r}(x):=u(r x)$. Then

$$
0<c r^{\lambda} \leq\left\|u_{r}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \leq r^{\lambda} \quad \forall r \in(0,1]
$$

where $c$ depends only on $n$ and $\lambda$.
Proof. - The inequality $\left\|u_{r}\right\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)} \leq r^{\lambda}$ follows from the fact that $r^{-2 \lambda} \mathrm{H}(r, u)$ is monotone nondecreasing since $\phi(r, u) \geq \lambda$ (see [ACS08, Lemma 2]). We need to show the bound from below (the nondegeneracy).

Define

## (B.8)

$$
\Psi(r):=\max \left\{\psi(r ; q): q \in \mathcal{Q}_{\lambda} \text { and }\|q\|_{\mathrm{L}^{2}\left(\partial \mathrm{~B}_{1}\right)}=1\right\}
$$

where $\psi$ is given by (B.4), and let $q_{r}^{*}$ be the function at which the above maximum is attained (note $\mathcal{Q}_{\lambda}$ is a closed convex subset of a finite dimensional vector space). Also, let Q be as in Lemma B.5, and define $\Phi(r):=\psi(r, \mathrm{D})$. Then, as a consequence of Lemma B.4, we have

$$
\frac{d}{d r} \Psi(r)=\theta(r) \Psi(r)-\frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta q_{r}^{*} \quad \text { for a.e. } r>0
$$

and
(B.9) $\quad \frac{d}{d r} \Phi(r)=\theta(r) \Phi(r)-\frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta \mathrm{Q} \quad \forall r>0$.

We now claim that

$$
\Psi(r) \asymp \Phi(r) \asymp \frac{\Psi(r)}{\Phi(r)} \asymp 1 \quad \text { as } r \downarrow 0 .
$$

Indeed, the accumulation points of $\tilde{u}_{r}$ (as $r \downarrow 0$ and in the $\mathrm{C}_{\mathrm{loc}}^{0}\left(\mathbf{R}^{n}\right)$ topology) belong to the unit ball of $\mathcal{Q}_{\lambda}$ (see [ACS08]) and therefore $\tilde{u}_{r}-q_{r}=o(1)$ for some $q_{r} \in \mathcal{Q}$. Hence, by definition of $\Psi$,

$$
\begin{aligned}
\Psi(r) \geq \psi\left(r ; q_{r}\right) & =\int_{\partial \mathrm{B}_{1}} \tilde{u}_{r} q_{r}-2 \int_{\partial \mathrm{B}_{1 / 2}} \tilde{u}_{r} q_{r}=\int_{\partial \mathrm{B}_{1}} q_{r}^{2}-2 \int_{\partial \mathrm{B}_{1 / 2}} q_{r}^{2}+o(1) \\
& =\left(1-2^{-n-1-2 \lambda}\right) \int_{\partial \mathbf{B}_{1}} q_{r}^{2}+o(1) \geq \frac{1}{2}>0 .
\end{aligned}
$$

Note that the above computation shows also that $\psi(r, q)=\left(1-2^{-n-1-2 \lambda}\right) \int_{\partial \mathrm{B}_{1}} q_{r} q+o(1)$, from which it follows that $q_{r}^{*}=q_{r}+o(1)$ as $r \downarrow 0$ (recall that $q_{r}^{*}$ is a maximizer in (B.8)).

Similarly, using Lemma B.5,

$$
\begin{aligned}
\Phi(r) & =\int_{\partial \mathbf{B}_{1}} \tilde{u}_{r} \mathbf{Q}-2 \int_{\partial \mathbf{B}_{1 / 2}} \tilde{u}_{r} \mathbf{Q}=\int_{\partial \mathbf{B}_{1}} q_{r} \mathbf{Q}-2 \int_{\partial \mathbf{B}_{1 / 2}} q_{r} \mathbf{Q}+o(1) \\
& \geq c_{n, \lambda}\left(1-2^{-n-1-2 \lambda}\right)+o(1) \geq \frac{c_{n, \lambda}}{2}>0,
\end{aligned}
$$

where $c_{n, \lambda}$ is the constant from Lemma B.5. Finally, it is clear that $\Psi(r)$ and $\Phi(r)$ are bounded by above, so the claim is proved.

Using the expressions for $\frac{d}{d r} \Psi$ and $\frac{d}{d r} \Phi$, we find

$$
\frac{d}{d r}\left(\frac{\Psi(r)}{\Phi(r)}\right)=-\frac{1}{r} \frac{\Psi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta q_{r}^{*}-\Phi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{w}_{r} \Delta \mathrm{Q}}{\Phi(r)^{2}}
$$

We claim that, given $\varepsilon>0$, for $r$ sufficiently small,
(B.10)

$$
\left|\int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta q_{r}^{*}\right| \leq \varepsilon\left|\int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta \mathrm{Q}\right| .
$$

Indeed, introducing the notation $\mathrm{B}_{r}^{\prime}:=\mathrm{B}_{r} \cap\left\{x_{n}=0\right\}$ and using Lemma B.3, given $\eta>0$ and choosing $r>0$ is sufficiently small so that $\left\|\tilde{u}_{r}-q_{r}^{*}\right\|_{L^{\infty}\left(\mathbf{B}_{2}\right)} \leq \delta(n, \eta)$ (recall that $q_{r}^{*}=$ $q_{r}+o(1)$ as $\left.r \downarrow 0\right)$, we have

$$
\begin{aligned}
0 \leq-\int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} u_{r} \Delta q_{r}^{*} & =\int_{\mathrm{B}_{1}^{\prime} \backslash \mathbf{B}_{1 / 2}^{\prime}} u\left(r x^{\prime}, 0\right) \mathrm{T}\left[q_{r}^{*}\right] d x^{\prime} \\
& =\eta \int_{\left(\mathrm{B}_{1}^{\prime} \backslash \mathrm{B}_{1 / 2}^{\prime}\right) \cap\left\{T\left[q_{r}^{*}\right] \leq \eta\right\}} u_{r} d x^{\prime} \leq \eta \int_{\mathrm{B}_{1}^{\prime} \backslash \mathbf{B}_{1 / 2}^{\prime}} u_{r} d x^{\prime}
\end{aligned}
$$

(recall that $u_{r} \geq 0$ on $\left\{x_{n}=0\right\}$ ), while

$$
-\int_{\mathbf{B}_{1} \backslash \mathbf{B}_{1 / 2}} u_{r} \Delta \mathrm{Q}=2\left|a_{0}\right| \int_{\mathrm{B}_{1}^{\prime} \backslash \mathrm{B}_{1 / 2}^{\prime}} u\left(r x^{\prime}, 0\right)\left|x^{\prime}\right|^{\lambda-1} d x^{\prime} \geq c_{n, \lambda}^{\prime} \int_{\mathrm{B}_{1}^{\prime} \backslash \mathbf{B}_{1 / 2}^{\prime}} u_{r} d x^{\prime},
$$

for some constant $c_{n, \lambda}^{\prime}>0$. Hence, dividing by $h(r)$, we obtain

$$
0 \leq-\int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta q_{r} \leq \mathrm{C}_{n, \lambda} \eta \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta \mathrm{Q},
$$

and (B.10) follows.

Then, thanks to (B.10), we deduce that

$$
\begin{aligned}
\frac{d}{d r}\left(\frac{\Psi(r)}{\Phi(r)}\right) & =-\frac{1}{r} \frac{\Psi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{l / 2}} \tilde{u}_{r} \Delta q_{r}^{*}-\Phi(r) \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta \mathrm{Q}}{\Phi(r)^{2}} \\
& \asymp \frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta \mathrm{Q}
\end{aligned}
$$

for $r \leq r_{0}$ small enough.
Integrating the above ODE over $\left[\hat{r}, r_{0}\right]$, since the integral of $\frac{d}{d r}\left(\frac{\Psi(r)}{\Phi(r)}\right)$ over $\left[\hat{r}, r_{0}\right]$ is uniformly bounded as $\hat{r} \rightarrow 0$, we deduce that the negative term $\frac{1}{r} \int_{\mathrm{B}_{1} \backslash \mathrm{~B}_{1 / 2}} \tilde{u}_{r} \Delta \mathrm{Q}$ is integrable over [0, $r_{0}$ ]. Hence, since $\Phi(r) \asymp 1$ and $\theta(r)=\frac{d}{d r} \log \left(h(r) / r^{\lambda}\right)$, it follows from (B.9) that

$$
\frac{d}{d r} \log \Phi(r)=\frac{d}{d r} \log \left(h(r) / r^{\lambda}\right)+g(r), \quad \text { with } g \in \mathrm{~L}^{1}\left(\left[0, r_{0}\right]\right) \text {. }
$$

Integrating over $\left[\hat{r}, r_{0}\right]$ and using again that $\Phi(r) \asymp 1$, we deduce that $\log \left(h(\hat{r}) / \hat{r}^{\lambda}\right)$ is uniformly bounded as $\hat{r} \rightarrow 0$, therefore $h(r) \asymp r^{\lambda}$, as desired.

As a consequence of Propositions B. 2 and B.6, we get the uniqueness and nondegeneracy of blow-ups:

Theorem B.7.-Let $u: \mathrm{B}_{1} \rightarrow \mathbf{R}$ be an even solution of (B.1) with $\phi\left(0^{+}, u\right)=\lambda$, where $\lambda=2 m+1$ is an odd integer. Then the limit

$$
\tilde{q}:=\lim _{r \downarrow 0} \frac{u(r \cdot)}{r^{\lambda}}
$$

exists, is non-zero, and it is a $\lambda$-homogeneous even solution of (B.1).
Thanks to this result, by classical arguments (see Proposition 8.3 and Lemma 6.6) one easily obtains the following rectifiability result, that was already proved with completely different methods in [FS18]:

Corollary B.8. - Let $u: \mathbf{B}_{1} \rightarrow \mathbf{R}$ be an even solution of (B.1). Then, for any odd integer $\lambda \geq 3$, the set of free boundary points of frequency $\lambda$ is $(n-2)$-rectifiable.

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[^1]:    ${ }^{1}$ More precisely, Lemma 4.3 gives the monotonocity of a "truncated frequency function", where the size of $w_{3}(r \cdot)$ is corrected by adding a term $r^{\gamma}$ with $\gamma \in(3,4)$. In order to relate the (almost) monotonicity of this truncated frequency function to the one of $\left\|\nabla \tilde{w}_{3, r}\right\|_{\mathrm{L}^{2}\left(\mathrm{~B}_{1}\right)}$, one needs to know that the size of $w_{3}(r \cdot)$ dominates $r^{\gamma}$ for some $\gamma \in(3,4)$.

[^2]:    ${ }^{2}$ In the past, other kind of truncations have been introduced (see in particular [CSS08]), but they do not work in our case due to the fact that D is not equal to (a multiple of) the derivative of H , as it is instead the case for the Signorini problem.

[^3]:    ${ }^{3}$ The formula for the ansatz can found by inspection, by trying to find the coefficients that guarantee the validity of Lemma 4.6.

[^4]:    ${ }^{4}$ More precisely, $\Sigma_{n-1}^{\geq 3} \backslash \Sigma_{n-1}^{3 n d}$ is the set in which any second blow-up (for $u-p_{2}$ ) is 3-homogeneous and nonharmonic, while $\Sigma_{n-1}^{3 n d} \backslash \Sigma_{n-1}^{>3}$ is the set in which the third blow-up (for $u-p_{2}-p_{3}$ ) is 3-homogeneous and non-harmonic.

[^5]:    ${ }^{5}$ The proof of this implication is standard. We want to prove that $\int \nabla \mathrm{Q} \cdot \nabla \xi=0$ for all $\xi \in \mathrm{C}_{c}^{1}\left(\mathbf{R}^{n}\right)$. But since $\left\{p_{2}=0\right\}$ has zero harmonic capacity, any given $\xi$ can be approximated in $\mathrm{W}^{1,2}$ norm by functions $\xi_{k}$ which vanish on $\left\{p_{2}=0\right\}$, and thus for which $\int \nabla \mathrm{Q} \cdot \nabla \xi_{k}=-\int \Delta \mathrm{Q} \xi_{k}=0$. The desired conclusion follows by taking the limit as $k \rightarrow \infty$.

[^6]:    ${ }^{6}$ Note again that $\hat{\mathbf{Q}}^{(1)}$ has a sign, $\hat{\mathbf{Q}}^{(2)}$ is (the translation of) a $\lambda^{2 n d}$-homogeneous solution of Signorini with $\lambda^{2 n d}>2$, and $\hat{Q}^{(3)}$ is a odd quadratic harmonic polynomial, and thus they are linearly independent.

[^7]:    ${ }^{7}$ In many textbooks, the definition of $\mathcal{H}_{\delta}^{\beta}$ includes a normalization constant chosen so that the Hausdorff measure of dimension $k$ coincides with the standard $k$-dimensional volume on smooth sets. However such normalization constant is irrelevant for our purposes, so we neglect it.

[^8]:    ${ }^{8}$ Note that, as a consequence of [Caf98], points in $\Sigma_{0}$ are always isolated and $u$ is strictly positive in a neighborhood of them.

[^9]:    ${ }^{9}$ More precisely, if we consider the odd reflection of $\left.u\right|_{\left\{x_{n}>0\right\}}$, then we obtain a global $\lambda$-homogeneous harmonic functions in the whole space. By Liouville Theorem, this functions mush be a $\lambda$-homogeneous harmonic polynomial.

